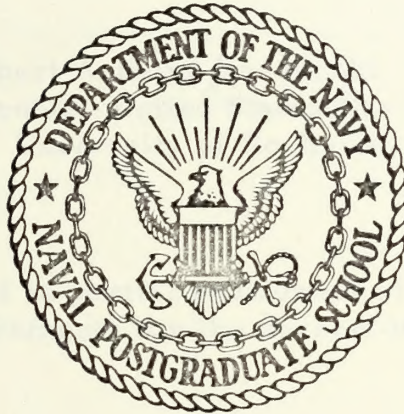


APPLICATION OF HOLOGRAPHIC INTERFEROMETRY TO
DENSITY FIELD DETERMINATION IN TRANSONIC
CORNER FLOW

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THESIS

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by

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ABSTRACT

The successful application of holographic interferometry to the study of density fields around opaque bodies in wind tunnel experiments has been reported in the literature. The present report extends this technique to the study of the three-dimensional asymmetric flow fields encountered near the wing-fuselage junction of an aerodynamic model in the transonic flow regime. Finite fringe interferometry has been used to obtain fringe information about a partially transparent wing-body structure. A FORTRAN computer program was utilized to invert the fringe information and produce a plot of the density field around the model. The resulting asymmetric density field and shock wave structure are shown to be an accurate representation of the phenomena encountered in aerodynamic corner flow.

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I. INTRODUCTION

The field of flow measurement has been revolutionized in recent years with the perfection of holography and holographic interferometry techniques. High power Q-switched and dye-switched lasers and sophisticated double-pulsing trigger mechanisms provide exposure times on the order of twenty nanoseconds, thereby "freezing" the flow during the hologram production process. The precision optical quality components and measurement techniques of Mach-Zehnder interferometry have given way to the much less restrictive requirements of holographic interferometry which provide high quality interferograms in three dimensions.

Techniques for the application of holography to interferometry have been reported by Heflinger, et al. [1], and by Brooks [2]. In the determination of the density field around a free jet in the supersonic regime, Matulka [3, 4] expressed the fringe data in a series of orthogonal polynomials and transformed them to polynomials representing density using an inversion technique reported in [5, 6]. The method was extended by Jagota [7, 8] to the determination of the three-dimensional density field around a ten-degree half angle cone in a supersonic wind tunnel. The ability to produce readable holograms in wind tunnel studies using transparent phase objects was verified by Heyer [9]. In the present report the aforementioned techniques

have been combined to study the three-dimensional density field near the wing-fuselage junction of an aerodynamic model in transonic flow. The experiment was conducted at the Naval Ship Research and Development Center, Carderock, Maryland in an eighteen inch transonic blow-down wind tunnel at a Mach number of 0.937, using a semi-transparent model of original design.

Since reasonably small variations in density were anticipated, a finite fringe technique was used in obtaining the interferograms. The horizontal finite fringe field was produced by a vertical translation of a diffusing glass in the scene beam a distance of 0.003 inches between the two exposures of the holographic plate. Fringe data obtained from the interferograms were reduced to density information using a modified form of the inversion computer program used in [7]. A self-testing procedure incorporated in the program verified the resulting density data as an accurate representation of the actual flow around the model.

II. EXPERIMENTAL APPARATUS

A. THE WIND TUNNEL

The investigation was conducted in the Naval Ship Research and Development Center blowdown supersonic wind tunnel. Transonic flow conditions were produced through incorporation of slotted upper and lower tunnel walls. The tunnel is a nominal eighteen inch blow-down-to-vacuum facility with a test section fourteen inches by eighteen inches in cross-section and twenty-nine inches in length with the slotted surfaces installed. Optical quality windows twenty-two inches in diameter in the side walls provided complete viewing of the flow in the test section as the model was rolled through 180 degrees for hologram production. A functional schematic of the wind tunnel is shown in Figure 1.

B. THE MACH NUMBER AND PRESSURE MEASUREMENT PROCEDURE

The Mach number at the test section is determined as a function of total and static pressure measurements and is maintained by carefully controlled butterfly valve settings. Total pressure is determined by recording atmospheric pressure prior to tunnel operation and accounting for the pressure loss between the plenum and the test section during operation. Static pressure is measured directly at a central wall port in the test section. Data recordings

were made on a Beckman Instruments Company 210 Digital Recorder, shown in Figure 2, and were read out on line via a Franklin Strip Tape Printer.

C. THE HOLOGRAPHIC ARRANGEMENT

The holographic arrangement is illustrated in Figure 3 and shown in photographs included as Figures 4, 5, 6, and 7. Two large wooden tables were constructed and linked together with two-by-four beams under the tunnel to form the experimental platform. Thick rubber pads were attached to the table legs to dampen possible floor vibrations. The bulk of the platform provided sufficient stability and vibration damping for the experiment. The monochromatic light source used was a KORAD K-1 pulsed ruby laser operating at a wavelength of 6943 Angstroms, together with a Pockels cell Q-switching device. The resultant effective exposure time was approximately twenty nanoseconds, eliminating the problems due to possible model vibration during hologram exposure. To maintain the laser head and output etalon at a constant temperature of 27.0 degrees Centigrade, a LAUDA constant temperature circulator Model K2R was used.

The reference beam was directed under the wind tunnel by four front surface mirrors, and the beam size was manipulated by means of lenses (Figure 3). The scene beam was routed through the test section to intersect the reference beam on the holographic plate at an angle of approximately fifty degrees. A diffuse glass, mounted on

a precision X-Y translation table in the scene beam, was used to produce light field holograms. Alignment of the Q-switched laser and system optics was accomplished using a continuous wave, low-power helium-neon laser. Reference grids were mounted accurately on the outer surfaces of the tunnel windows using a surveyor's transit. Details of the model mounting and reference grids are shown in Figure 7. To enable hologram production during daylight hours, the entire tunnel room was blacked out using drop curtains and light shields.

D. THE WIND TUNNEL MODEL

The aerodynamic model used is shown in Figures 8, 9, and 10. The metal portions of the model were stainless steel. The greater part of the modified double wedge platform wing was constructed of optical lucite, as was the portion of the fuselage at the wing root. Detailed model dimensions are shown in Figure 11. The choice of aerodynamic design provided good flow characteristics and a relatively stable lambda-type shock wave on the wing; the largely transparent construction facilitated hologram production through 180 degrees of view.

The model was rotated about its sting mount in the wind tunnel from the zero degree position, wings level, to the 180 degree position, wings level inverted. Alignment for the desired rotation angle was accomplished by manually aligning prescribed lines on the sting mount collar with a scribed mark on the sting support.

III. ANALYTICAL EVALUATION OF THE DENSITY FIELD

A. THE BASIC EQUATION OF INTERFEROMETRY

Interferograms are created when two originally coherent light beams are superimposed and projected on a viewing screen. The two rays will reinforce or annul each other, depending on their relative phase difference at the screen. This phase difference is directly a function of the optical pathlengths traversed by the two waves.

Consider a coherent beam which is split and then recombined on a viewing screen. A difference in optical pathlengths of the two component beams may be achieved by causing the beams to traverse through different media prior to recombination, with their physical pathlengths maintained equal. Each component beam will travel at a speed c_0/n where c_0 is the speed of light in a vacuum and n is the index of refraction of the medium traversed. The difference in optical pathlength is then given by

$$L = L (n_2 - n_1) = c_0 \Delta t \quad (1)$$

where Δt is the time difference of travel in the two media. If the optical pathlength is changed by an amount $N\lambda$, where λ is the wavelength of the light source and N is an integer, then the order of interference changes by an amount N . In other words, a shift of N fringes occurs in the interference pattern. The fringe shift may be expressed as

$$g = L/\lambda \quad (2)$$

where

g = fringe shift

λ = light source wavelength

L = change in optical pathlength

Substituting equation (1) into equation (2) yields

$$g = \frac{L}{\lambda} (n_2 - n_1) \quad (3)$$

The index of refraction for a given medium is a function of density. In the case of gases, since the speed of light is very nearly the same as in a vacuum, the index of refraction is well represented by the first two terms of a Taylor series expansion [10]:

$$n = 1 + \beta \frac{\rho}{\rho_s} \quad (4)$$

where β = dimensionless constant related to the Gladstone-Dale constant by $K = \beta/\rho_s$

ρ_s = reference density at 0° C, 760 mm. Hg.

The value of β for air at $\lambda = 5893$ Angstroms (deep red light) is 0.000292; variation with wavelength is very small.

For a fixed difference in the index of refraction between the two component beams the fringe shift relation becomes:

$$g = \beta \frac{L}{\lambda} \left(\frac{\rho_2 - \rho_\infty}{\rho_s} \right) \quad (5)$$

For variable density in the test section, the net change in optical pathlength is the integrated effect along the beam path, or

$$g = \frac{\beta}{\lambda \rho_s} \int_0^L (\rho - \rho_\infty) ds = Q \int_0^L f(x, y, z_c) ds \quad (6)$$

where:
$$Q = \frac{\beta \rho_{\infty}}{\lambda \rho_s} \quad (7)$$

$$f(x, y, z_c) = \frac{\rho(x, y, z_c)}{\rho_{\infty}} - 1 \quad (8)$$

z_c = plane of constant z

ds = incremental distance along beam path

Equation (6) is the basic integral equation for the unknown density.

With known fringe shift values from an interferogram, the equation is inverted to obtain the density along a beam path.

B. THE INTEGRAL INVERSION

The integral inversion procedure utilized in this investigation was first reported by C. D. Maldonado, et al [5, 6]. It was used subsequently by R. D. Matulka [3] and R. C. Jagota [7] to determine the density variation in an asymmetric free jet and about a cone at angle of attack in supersonic flow, respectively. The procedure involves the representation of the function $f(x, y, z_c)$ of Equation (6) in a complete set of orthogonal functions, with the expansion coefficients evaluated by use of the orthogonality condition between the functions f and g of Equation (6). The set of functions used is orthogonal over the entire (x, y) plane for every z_c and remains an orthogonal set under a rotation of the coordinate system. The coordinate system used for the inversion is shown in Figure 12. It consists of (a) a set of fixed coordinates x, y and (b) a set of moving coordinates

x', y' in which the fringe number function is defined and which rotates with respect to x, y as the viewing angle through the test section is varied.

In operator form, Equation (6) can be represented as

$$g(\xi, y', z_c) = T f(x, y, z_c) \quad (9)$$

and f is evaluated by inversely transforming the equation to obtain

$$f(x, y, z_c) = T^{-1} g(\xi, y', z_c) \quad (10)$$

This inversion is achieved by utilizing a pair of orthogonal polynomials $U_{m+2k}^{+m}(\alpha x, \alpha y)$ and $H_{m+2k}(\alpha y')$ which are related by the transform relationship

$$T[U_{2k}(\alpha x, \alpha y) e^{-\alpha^2 x'^2}] = \frac{e^{\pm i m \xi}}{[k!(m+k)!]^{1/2}} \cdot \frac{1}{2^{m+2k}} \cdot H_{m+2k}(\alpha y') \quad (11)$$

where $H_{m+2k}(\alpha y')$ are Hermite polynomials of order $m+2k$. The

unknown function $f(x, y, z_c)$ is expanded in a set of functions U_{m+2k}^{+m} as

$$\begin{aligned} f(x, y, z_c) = & \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \epsilon_m \left\{ C_{m+2k}^{+m}(\alpha) U_{m+2k}^{+m}(\alpha x, \alpha y) \right. \\ & \left. + C_{m+2k}^{-m}(\alpha) U_{m+2k}^{-m}(\alpha x, \alpha y) \right\} e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \end{aligned} \quad (12)$$

where $\epsilon_m = 1/2$ for $m = 0$, $\epsilon_m = 1$ for $m = 1, 2, 3, \dots$, and

C_{m+2k}^{+m} are the unknown coefficients of the expansion. α is an arbitrary scale factor which may be considered the reciprocal of a non-dimensionalizing coefficient.

Utilizing the transform relation of Equation (11), Equation (6)

becomes

$$g(\xi, y', z_c) = \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \varepsilon_m [k!(m+k)! 2^{2(m+2k)}]^{1/2} \times [C_{m+2k}^{+m}(\alpha) e^{im\xi} + C_{m+2k}^{-m}(\alpha) e^{-im\xi}] H_{m+2k}(\alpha) e^{-\alpha^2 y'^2} \quad (13)$$

Equation (13) is subject to the orthogonality condition

$$\int_{-\pi}^{\pi} e^{\pm im\xi} e^{\mp in\xi} d\xi \int_{-\infty}^{+\infty} H_{m+2k}(\alpha y') H_{n+2l}(\alpha y') e^{-\alpha^2 y'^2} dy' = \frac{2\pi^{3/2}}{\alpha} [(m+2k)!(n+2l)! 2^{m+2k} 2^{n+2l} \delta_{mn} \delta_{(m+2k)(n+2l)}] \quad (14)$$

where δ is the Kronecker delta function. The solution of Equation

(14) applied to Equation (13) yields the series coefficients

$$C_{m+2k}^{\pm m}(\alpha) = \frac{\alpha}{2\pi^{3/2}} \left[\frac{k!(m+k)!}{(m+2k)!} \right] \int_{-\pi}^{\pi} \int_{-\infty}^{+\infty} g(\xi, y', z_c) H_{m+2k}(\alpha y') e^{\mp im\xi} dy' d\xi \quad (15)$$

With the substitution of the coefficients of Equation (15), Equation (7)

becomes

$$f(x, y, z_c) = \frac{\alpha}{\pi^{3/2}} \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \varepsilon_m \frac{[k!(m+k)!]^{1/2}}{(m+2k)!} e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \times \text{Re} \left[\int_{-\pi}^{\pi} \int_{-\infty}^{+\infty} g(\xi, y', z_c) e^{-im\xi} H_{m+2k}(\alpha y') dy' d\xi \times U_{m+2k}^m(\alpha x, \alpha y) \right] \quad (16)$$

The functions $U_{m+2k}^{\pm m}$ are defined as

$$U_{m+2k}^{\pm m}(\alpha x, \alpha y) = (-1)^k \alpha \left[\frac{k!(\alpha^2 x^2 + \alpha^2 y^2)^m}{\pi(m+k)!} \right]^{1/2} e^{\pm im\phi} L_k^m(\alpha^2 x^2 + \alpha^2 y^2) \quad (17)$$

where $\phi = \tan^{-1}(y/x) - (\pi/2)$ and L_k^m are the associated Laguerre polynomials

$$L_k^m(\alpha^2 x^2 + \alpha^2 y^2) = \sum_{s=0}^{\infty} \frac{(m+k)!}{(k-s)!(m+s)! s!} \left[(-1)(\alpha^2 x^2 + \alpha^2 y^2) \right]^s \quad (18)$$

Insertion of Equation (17) into Equation (16) yields

$$f(x, y, z_c) = \left(\frac{\alpha}{\pi} \right)^2 \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \epsilon_m \frac{(-1)^k k!}{(m+2k)!} (\alpha^2 x^2 + \alpha^2 y^2)^{m/2} L_k^m(\alpha^2 x^2 + \alpha^2 y^2) \quad (19)$$

$$\times \left[B_{m+2k}^m(\alpha) \cos(m\phi) + D_{m+2k}^m(\alpha) \sin(m\phi) \right] e^{-(\alpha^2 x^2 + \alpha^2 y^2)}$$

where

$$B_{m+2k}^m(\alpha) = \int_{-\pi}^{\pi} \int_{-\infty}^{\infty} g(\xi, y', z_c) \cos(m\xi) H_{m+2k}(\alpha y') dy' d\xi \quad (20)$$

$$D_{m+2k}^m(\alpha) = \int_{-\pi}^{\pi} \int_{-\infty}^{\infty} g(\xi, y', z_c) \sin(m\xi) H_{m+2k}(\alpha y') dy' d\xi \quad (21)$$

Equations (19), (20), and (21) are the basic equations used to obtain the density distribution from the experimentally determined fringe variations in a completely asymmetric flow field.

C. THE NUMERICAL PROCEDURE

Because the function $g(\xi, y', z_c)$ is an experimentally determined quantity the unknown coefficients $B_{m+2k}^m(\alpha)$ and $D_{m+2k}^m(\alpha)$ in the series representation of $f(x, y, z_c)$ in Equation (19) cannot be calculated analytically. It is therefore necessary to evaluate the double integrals of Equations (20) and (21) numerically. This is accomplished by noting in Figure 12 and Equation (8) that there is an area outside which the density is invariant, namely outside the test

section where the known density is ρ_∞ . Since the function $f(x, y, z_c) = 0$ outside this circular domain, the limits of integration of $+\infty$ and $-\infty$ in Equations (20) and (21) can be replaced by finite values. The fringe distribution is then approximated by small increments over the test domain, resulting in the representation of the B and D coefficients as double series:

$$B_{m+2k}^m(\alpha) = \sum_{i=1}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta \xi_j, x_i + \Delta x_i) \int_{\xi_j}^{\xi_{j+1}} \cos(m\xi) d\xi \int_{x_i}^{x_{i+1}} H_{m+2k}(\alpha x) dx \quad (22)$$

and

$$D_{m+2k}^m(\alpha) = \sum_{i=1}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta \xi_j, x_i + \Delta x_i) \int_{\xi_j}^{\xi_{j+1}} \sin(m\xi) d\xi \int_{x_i}^{x_{i+1}} H_{m+2k}(\alpha x) dx \quad (23)$$

Using the derivative formula for Hermite polynomials, Equations (22)

and (23) can be manipulated to yield workable series expressions:

$$B_{m+2k}^m(\alpha) = \left[\frac{1}{2\alpha m} \cdot \frac{1}{(m+2k+1)} \right] \sum_{i=0}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta \xi_j, x_i + \Delta x_i) \times \left[\sin(m\xi_{j+1}) - \sin(m\xi_j) \right] \left[H_{m+2k+1}(\alpha x_{i+1}) - H_{m+2k+1}(\alpha x_i) \right] \quad (24)$$

$$D_{m+2k}^m(\alpha) = - \left[\frac{1}{2\alpha m} \cdot \frac{1}{(m+2k+1)} \right] \sum_{i=0}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta \xi_j, x_i + \Delta x_i) \times \left[\cos(m\xi_{j+1}) - \cos(m\xi_j) \right] \left[H_{m+2k+1}(\alpha x_{i+1}) - H_{m+2k+1}(\alpha x_i) \right] \quad (25)$$

Since it is impossible to sum over an infinite number of terms,

Equation (19) is necessarily expressed as the sum of a finite series:

$$f(x, y, z_c) = \left(\frac{\alpha}{\pi}\right)^2 \sum_{k=0}^K \sum_{m=0}^M \varepsilon_m (-1)^k \left[\frac{k!}{(m+2k)!} \right] (\alpha^2 x^2 + \alpha^2 y^2) \quad (26)$$

$$\times \left[\sum_k^m (\alpha^2 x^2 + \alpha^2 y^2) \left[B_{m+2k}^m(\alpha) \cos(m\phi) + D_{m+2k}^m(\alpha) \sin(m\phi) \right] e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \right]$$

It has been demonstrated that judicious selection of the parameters

$\Delta \xi$, Δx , K , M , and α yields density distributions with very good accuracy [3, 6].

IV. EXPERIMENTAL PROCEDURE

A. LABORATORY TECHNIQUES

In order to visualize the general flow patterns and localize shock or expansion waves about the model a series of standard Schlieren photographs were taken at varying flow Mach numbers. Pictures were produced for roll angles of 0° , 45° , and 90° at Mach numbers from 0.925 to 1.10. A representative series of Schlieren photographs is shown in Figures 13, 14, and 15. Analysis of the Schlieren photographs dictated a flow Mach number of 0.937 for the experimental study; this Mach number yielded uniform upstream flow conditions and located the lambda-type shock wave ideally near the center of the lucite section of the model wing.

The coherence length of the pulsed ruby laser was approximately ten centimeters for the output power utilized. This reduced the normally critical requirement for pathlength equality in the scene and reference beams that must be fulfilled in the classical Mach-Zehnder interferometric approach. Consequently, a length of string proved to be a sufficiently accurate measuring device to maintain the two beam pathlengths within the coherence length of the laser, a requirement for interferogram production. To compensate for the fact that the scene beam traversed approximately five inches of glass tunnel walls and lucite grids which the reference beam did not, the scene beam



was adjusted to be some 2.5 inches shorter than the reference beam. Reference beam pathlength was maintained at approximately 138 inches throughout the experiment.

Holograms produced using the basic holographic setup shown in Figure 3 exhibited clear, well-defined fringe patterns in nearly every instance. In deciding on the final arrangement, several techniques were tested to improve upon fringe pattern definition. Horizontal, vertical and diagonal translations of the diffuser plate in the scene beam were considered, varying from 0.001 inches to 0.005 inches. A vertical translation of 0.003 inches yielded clear horizontal fringes that were quite easily analyzed. The transverse mode selector aperture was varied from 1.0 mm. to 3.0 mm. in increments of 0.5 mm; best lighting of the model resulted with use of a 2.5 mm. aperture. The temperature of the cooling water circulated through the laser head and etalon was varied from 26.0°C . to 28.0°C . in increments of 0.2°C ., with 27.0°C . providing the best fringe definition. Finally, a variety of beam splitters and lenses were tested prior to final selection of the best available optics arrangement for the experiment. A 2:1 reference to scene beam strength ratio was found to yield very good holograms.

Two double exposure holograms were taken for each model viewing angle. The first, labeled a double-static exposure, consisted of two exposures in a no-flow condition with a 0.003 inch vertical translation of the diffuser plate between exposures. The fringe

patterns in this hologram provided a measure of the effect of tunnel wall glass, grid plexiglass and model lucite on the subsequent double exposure. The second, or static-dynamic, exposure consisted of a no-flow exposure, a 0.003 inch diffuser translation, and finally an exposure at flow Mach number 0.937. The fringe deviations recorded in the region behind the lucite portion of the model by the double-static hologram were measured and subtracted from the fringe shifts measured in the corresponding static-dynamic hologram.

Holograms were produced on Agfa-Gaevert 8E75 holographic plates, 4 inches by 5 inches in size. As recommended by Collier, et al. [11] the development process included:

1. Five minutes in Kodak D-19 developer
2. Thirty seconds in a flowing water bath
3. Five minutes in standard rapid fixer
4. Thirty seconds in a flowing water bath
5. One and one-half minutes in Kodak Hypo Clearing agent
6. Five minutes in a flowing water bath
7. Five minutes in methanol bath
8. One minute in a flowing water bath
9. Drying

B. PHOTOGRAPHIC TECHNIQUES

Normal reconstructions of the original scene were made by illuminating the holograms with a seven milliwatt continuous wave

helium-neon laser beam at a wavelength of 6328 Angstroms. There was some slight distortion in the reconstructed scene because of the difference in wavelengths of the original scene beam and the reconstruction beam; however, the effect was almost totally negated by shrinkage of the holographic plate emulsion during the development process.

A common technique of image reconstruction was employed, utilizing a conjugate reference beam to reilluminate the exposed hologram, as shown in Figure 16. The resulting scene was recorded on photographic film. Individual points on the photograph are produced by a series of non-parallel rays originating from various source points on the diffuser plate in the scene beam. Using a reilluminating beam of small diameter has the effect of a small aperture at the focal point of the imaging lens, filtering out all but a set of nearly parallel rays, as shown in Figure 17. The real images produced in this manner have a large depth of field, permitting simultaneous projection on the film of front and rear grids, the model and the fringe patterns. The imaging lens was focused as near to the plane of the fringes as possible, producing photographs at various planes of constant z_c .

C. DATA REDUCTION

Photographic interferograms were obtained using the arrangement shown in Figure 16, with the camera viewing screen in the position of the real image. The line of sight in the plane desired was achieved by translating and elevating the hologram until common points on the

front and rear grids were aligned. The Graphic View camera, with the wide angle lens aperture set fully open at f7.8, was adjusted to yield the best focus on the fringe plane. Exposure times of from 1/5 to 3/4 seconds were used to produce workable interferogram photographs on Polaroid Type 55 P/N film.

Fringe shift analysis was accomplished on 8 inch by 10 inch enlargements of the 4 inch by 5 inch film used to record the images. The enlargements were placed face down on a light table, and the fringes, model contours and shock wave were traced on the back surface at the desired cross-sectional plane. Fringe shift values were recorded by measuring the distance between (1) the intersection of the hypothetically undeviated fringe and the cross-sectional plane and (2) the intersection of the deviated fringe with the same plane. Fringe shifts so obtained were corroborated by placing the negative in a photographic enlarger and tracing the lines of interest directly onto graph paper.

The known model fuselage diameter of 1.1 inches was compared with that measured in each individual photograph to yield magnification factors relating projected dimensions to actual dimensions. These factors were then used as corrections to the fringe shift measurements. A grid reference point located on the cross-sectional plane of interest served as the datum for all fringe shift measurements. A base point was located at the intersection of the cross-sectional plane of interest and the body longitudinal axis. Fringe shift measurements were

corrected using this base point as the new datum so that the inversion circle was properly centered on the body axis. The radius of the inversion circle was selected to be nearly equal to the semi-span of the wing. Fringe shift measurements were converted to fringe numbers using the average free stream spacing, while fringe locations were nondimensionalized using the inversion circle diameter. From the data so obtained, the radial variation of fringe number was plotted for each viewing angle. Fringe numbers at 201 equidistant points, as required for input into Mode 3 of the computer program, HOLOFER, were recorded from the resulting curves. Further details concerning this inversion computer program are outlined in Appendix B. A typical reduction of an interferogram to obtain the radial variation of fringe number at a particular cross-sectional plane is detailed in Appendix A.

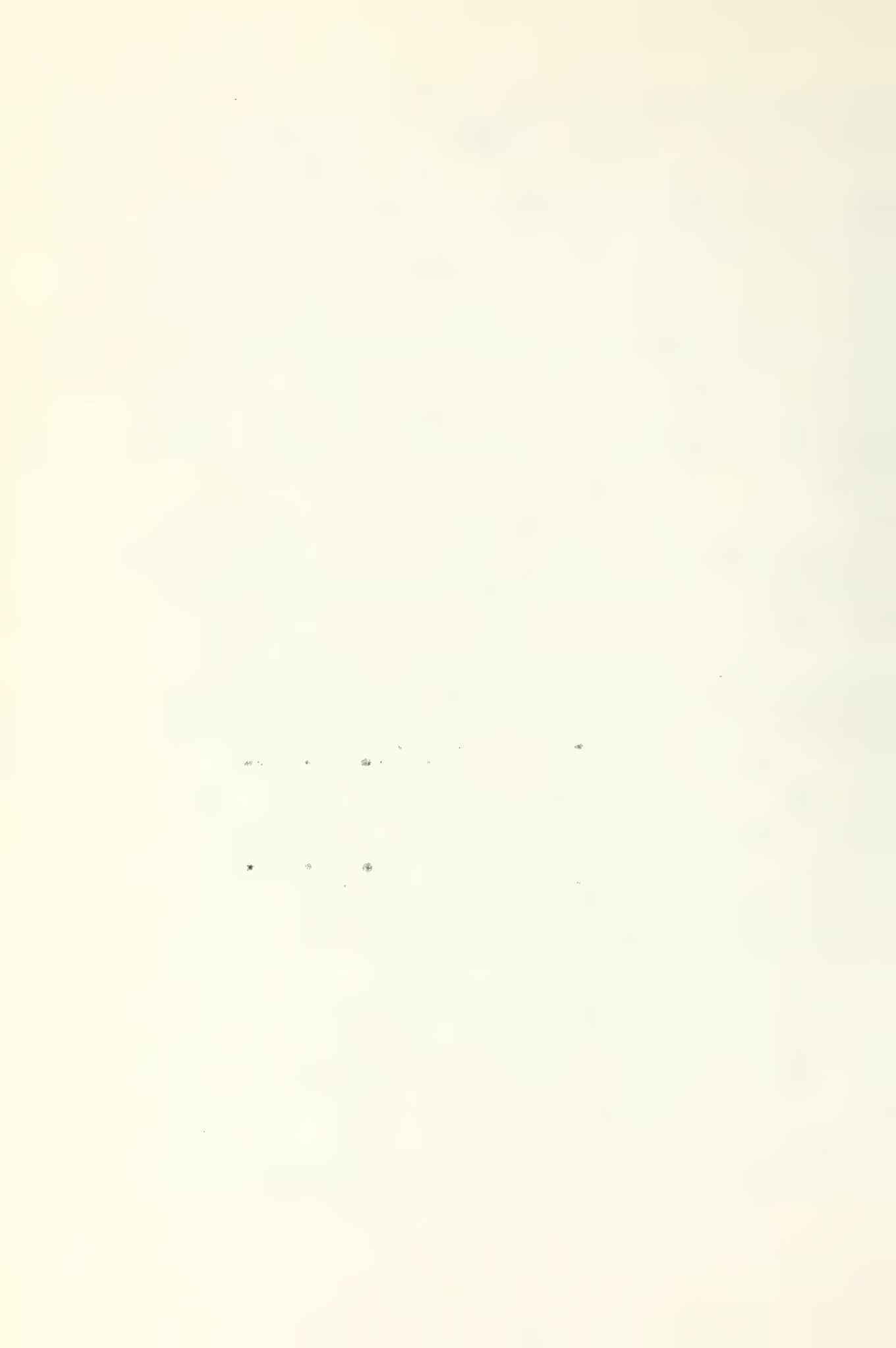
V. EXPERIMENTAL RESULTS AND DISCUSSION

A pair of double exposure holograms was taken of the model at $11\frac{1}{4}$ degree intervals through a 180 degree field of view. Experimental data from the wind tunnel runs are recorded in Table I. Initial resulting density patterns indicated relatively smooth contours across adjacent intervals; the interval was therefore doubled to $22\frac{1}{2}$ degrees to simplify and speed the analysis. Fringe data were first inserted into the inversion computer program along nine lines of sight in the 180 degree field of view. A numerical comparison of views from 0 degrees to 90 degrees and from 90 degrees to 180 degrees verified to within 0.20 percent the assumption of a single plane of symmetry in the experiment. The fringe data input was then reduced to five lines of sight in a 90 degree field of view, as shown in Figure 18. The resulting output was an inverted density field along nine radial lines spanning a 180 degree field of view, with a mirror image on the opposite side of the plane of symmetry, as shown in Figure 19.

The static-dynamic photographic interferogram for the 0 degree view, along with its corresponding double-static interferogram, is shown in Figure 20. The diffraction effects caused by the presence of the lucite portions of the model are clearly visible in the double-static exposure, where the free stream fringe lines are bent and displaced toward the model axis. This displacement was measured

and subtracted from subsequent measurements made on the static-dynamic exposure, as outlined in Section IV.A. Photographic interferograms of the remainder of the static-dynamic exposures are shown in Figures 21 through 24. Clearly visible and reduceable in nearly all views were (1) the region of uniform subsonic flow, commonly called the free stream condition, (2) the transition from local subsonic to local supersonic flow, and (3) the lambda-type shock wave on the model wing. These characteristics are shown in schematic representation in Figure 25.

Contour plots of the density function, as expressed in Equation (8) of Section III.A., for successive z-planes of analysis are shown in Figures 26 and 27. The cross-sectional plane of analysis for the plot of Figure 26 was located at 186.75 mm. from the model nose along the longitudinal axis. For Figure 27, the plane of interest was 195.25 mm. from the model nose. It is apparent from both contour plots that the model went to a very small angle of attack under the loading forces produced during tunnel operation; this is evidenced by the compression of the contour lines above the model and the corresponding expansion of the contours below the model. Measurements made from photographic interferograms confirm this angle of attack to be, at most, 0.05 degrees. The closed contours above and adjacent to the wing surface in both figures may very well be the result of a vortex originating at the intersection of the wing leading edge and the



fuselage on either side of the model and traveling aft and outward over the wing surface.

A comprehensive quantitative analysis of the shock wave structure was not undertaken, with the exception of estimating the strength of the shock by comparison of fringe line separation immediately ahead of and aft of the shock wave. Fringe line separation measurements on either side of the shock wave were converted first to density information and thence to pressure information, disregarding compressibility effects. An approximate strength value of 0.207 was computed using the accepted definition of $(p_2 - p_1)/p_1$. This corresponds to a local Mach number of 1.08 in the supersonic region just ahead of the shock wave. Qualitative shock wave analysis resulted in the construction of a three-dimensional structural representation as shown in Figure 28, using input information from several interferogram viewing angles. While location of the leading and trailing edges of the lambda-type shock wave was very accurate, interior structure was largely indefinable due to "smearing" and blurring of the fringes transiting the shock wave itself.

As a preliminary step to possible future studies in this field, photographic interferograms were made from holograms produced with the aerodynamic model set at small angles of attack. Orientations included angles of attack of five and ten degrees, with roll angles varying from zero to ninety degrees. Although the holograms themselves were of very good quality, the photographic reproductions

were relatively poor due to the fact that an inferior photographic arrangement had to be used. They were therefore omitted from this report. It was inferred from the holograms, however, that a complete study at angle of attack using the basic procedures followed in the present study would be both totally feasible and rewardingly fruitful.

The original character of the experimental data prevented comparison with published results. Qualitative studies of transonic phenomena are widely available, and the general characteristics of the resulting density field and shock wave structure serve to bear out the schematics based on theoretical and mathematical models. Moreover, the self-testing mode of the inversion computer program, HOLOFER, verified the consistency and reproductibility of the resulting density distributions to within 2.0 percent through proper choice of the input parameters, primarily the slope-matching parameter α . The errors encountered in the final results are due primarily to errors in the fringe data input to the inversion program. The intrinsic presence of laser speckle, the extended pathlengths of the scene and reference beams and the unavoidable beam scattering and diffraction within the lucite model sections, created difficulty in obtaining precisely the slope of the fringe lines behind the lucite sections. Fringe spacing measurements in the free stream flow were conservatively judged accurate to within 0.5 mm. This assumption was quite reasonable since all measurements were effected with a scale graduated at half millimeter intervals. This figure of 0.5 mm.,

combined with the mean free stream fringe spacing of 5.197 mm. for all interferograms, indicated measurement accuracy to within one tenth (0.1) of a fringe. The mean systematic error of the free stream spacing in each view was computed to be a maximum of 3.9 percent. Associated with this systematic error was a random error of 2.1 percent in the measurement of fringe shifts in each view to a conservative accuracy of 0.5 mm. The resulting error for each viewing angle was therefore a maximum of 6.0 percent, found by merely adding the two types of error for each view. The minimum error limit was found by considering the error resulting from the reproduction of the same interferogram view five separate times. Statistically, with 6.0 percent error in each view, the composite error for the repeated view is 2.6 percent. As five different views, or lines of sight, were used for the data between zero and ninety degrees, the final total error in the analysis was therefore in the interval between 2.6 percent and 6.0 percent. To insure contour clarity and guard against overlapping, the maximum error figure of 6.0 percent was used in construction of the plots shown in Figures 26 and 27. In general, the rather large fringe shifts led to a very low mean fringe sensitivity value of 0.1259. This coefficient indicated a resulting density function (Equation (8)) inaccuracy of less than 1.5 percent for a fringe shift measurement misreading of 0.5 mm.

Physical limitations of beam diameter and hologram plate area dictated the choice of an inversion circle diameter somewhat smaller

than the full data circle normally used in the finite fringe procedure. This, in effect, introduced an inconsistency in reference density into the analysis since the density on the selected circle and immediately outside it was not the calculated ρ_{∞} ; the density function was therefore not zero outside the actual inversion region. To alleviate this inconsistency, a new, updated reference density was computed for each cross-sectional plane of analysis by averaging the density values on the selected circle from the first inversion process. The actual reference densities used were $\rho = 1.777 \text{ mg/cc}$ for the 186.75 mm. plane and $\rho = 1.642 \text{ mg/cc}$ for the 195.25 mm. plane. This procedure was justified since all density values between the selected circle and the full data circle were constant to within approximately fifteen percent. The updated reference densities were then used to produce the final output density field. The net effect was a scaled, uniform shift toward density function values slightly lower than those computed on the basis of the original reference density.

VI. CONCLUSIONS

The finite fringe procedure for the production of holographic interferograms has been applied successfully to the determination of the three-dimensional density distribution of the flow near the wing-fuselage junction of a partially transparent aerodynamic model in the transonic regime. Density contours accurate to within six percent enabled a thorough analysis of the flow field to be conducted, highlighting flow characteristics and the presence of the shock wave. Subsequent studies of similar models at angle of attack have been shown to be entirely feasible. Procedures used in the experiment also exhibit promise for the direct analysis of duct and inlet flows as well as comprehensive study of shock wave structure.

The inversion computer program, HOLOFER, was found to be adequate in handling a general asymmetric flow field analysis. However, it was considered quite cumbersome and difficult to modify for various experimental situations. Subsequent analysts will find the procedures advocated by Sweeney and Vest [12] for the recording and analyzing of interferograms of considerable interest. In addition, the efforts of Van Houton [13], who utilized the method proposed by Junginger and van Haeringen [14], may prove valuable in reducing computer time significantly.

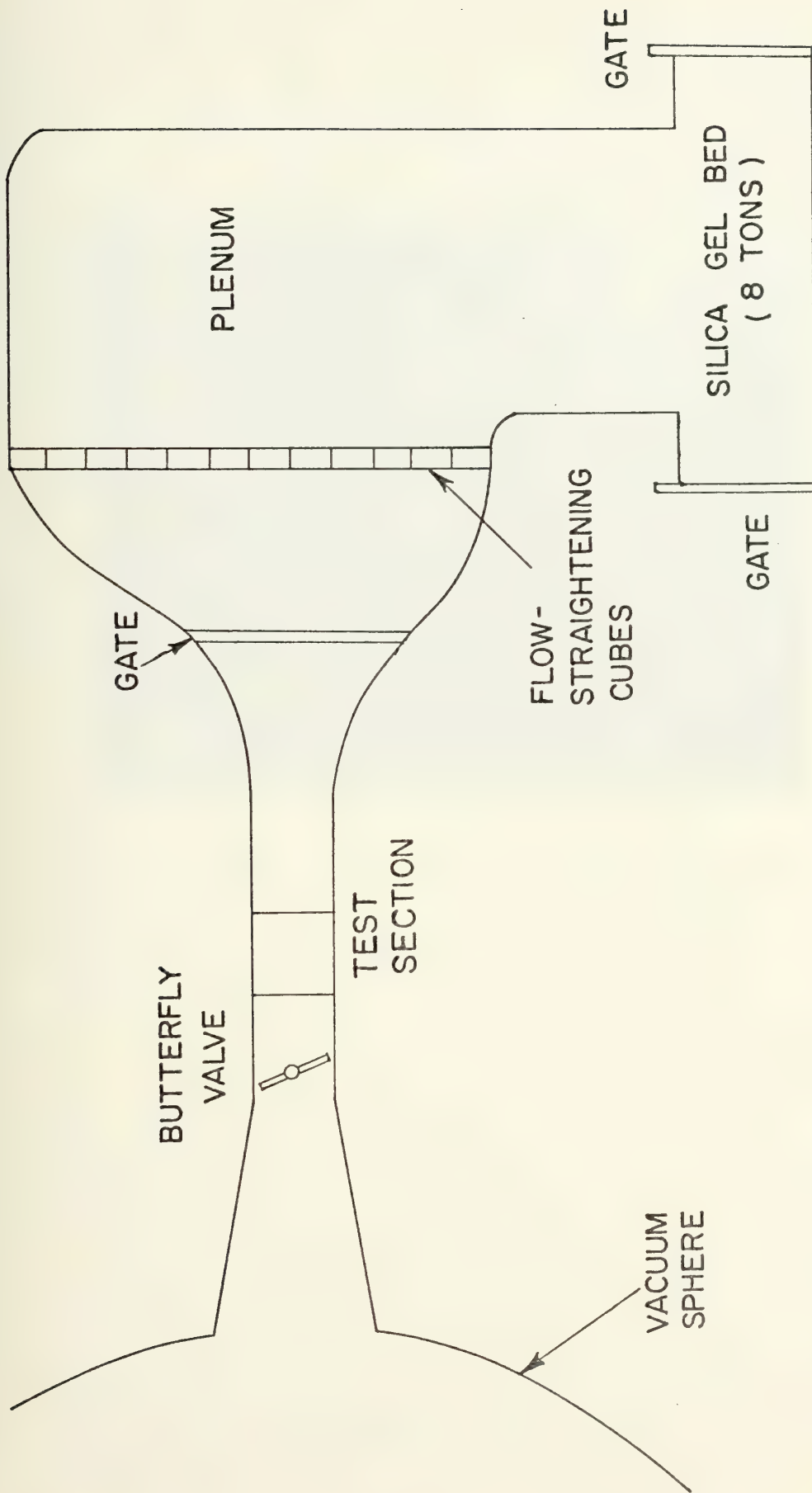


FIGURE 1. FUNCTIONAL SCHEMATIC OF NSRDC WIND TUNNEL

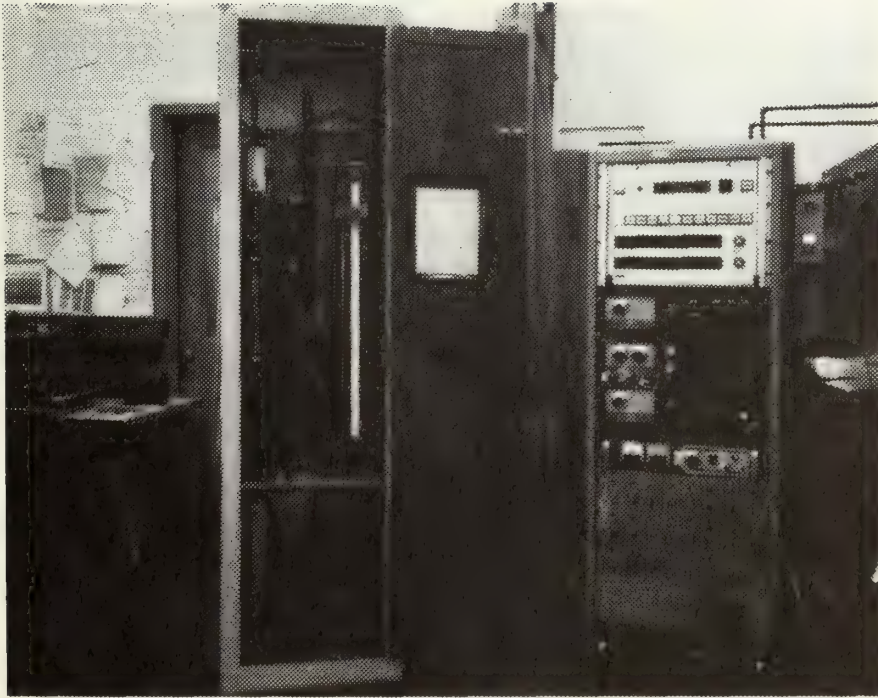


FIGURE 2. BECKMAN 210 DATA RECORDING SYSTEM

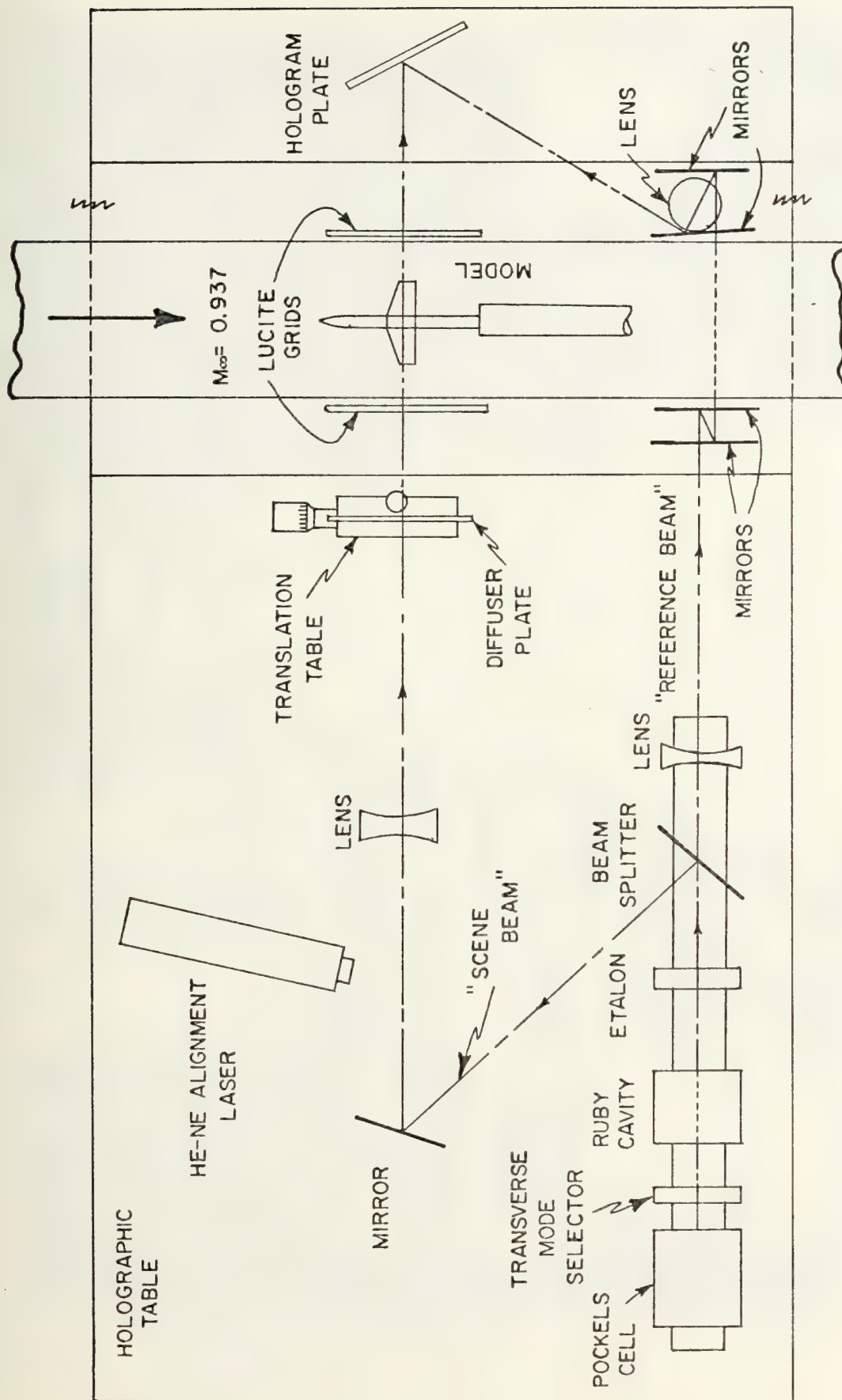


FIGURE 3. SCHEMATIC DRAWING OF THE HOLOGRAPHIC ARRANGEMENT

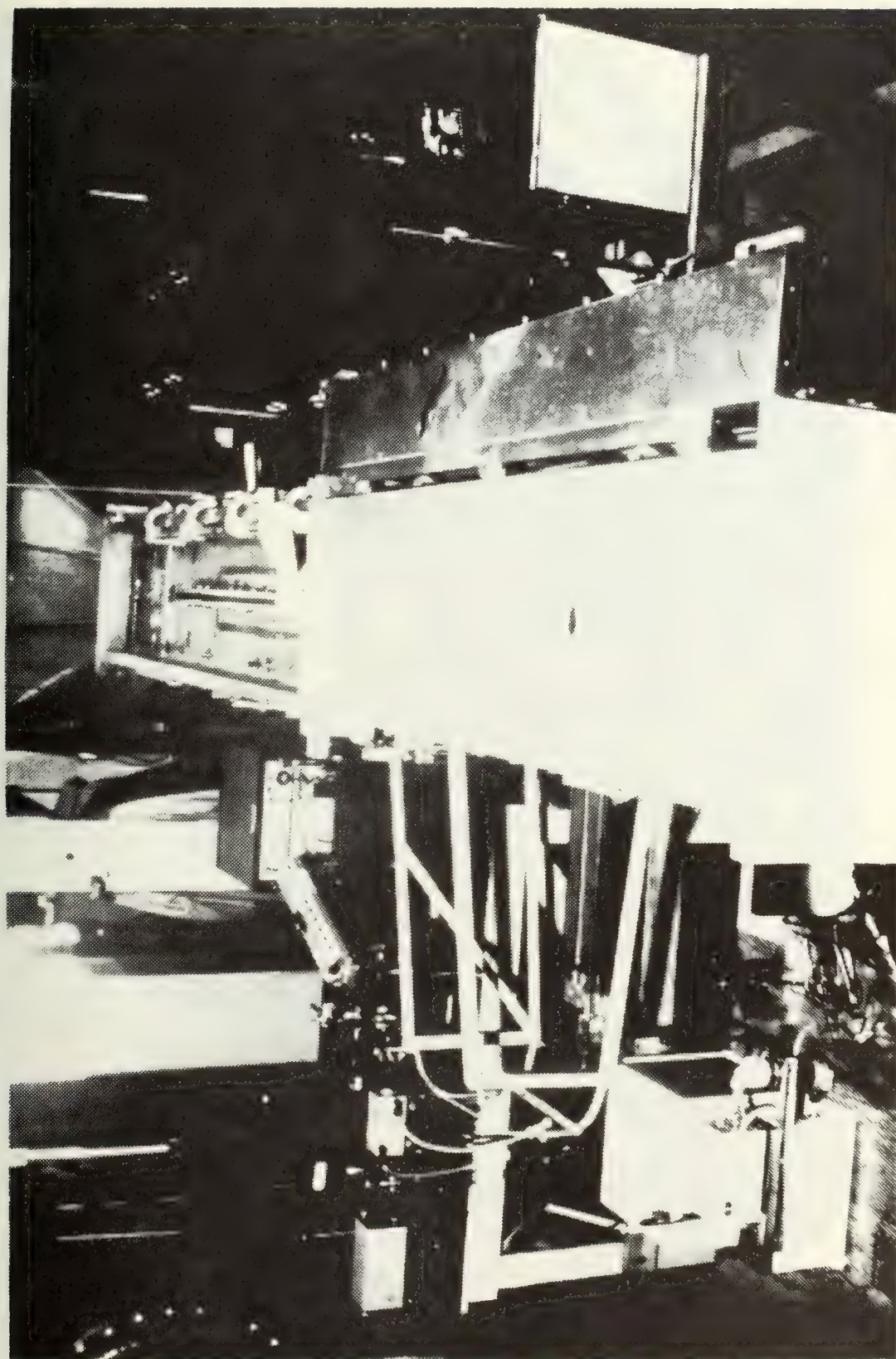


FIGURE 4. OVERHEAD VIEW OF TUNNEL AND ENTIRE HOLOGRAPHIC SYSTEM

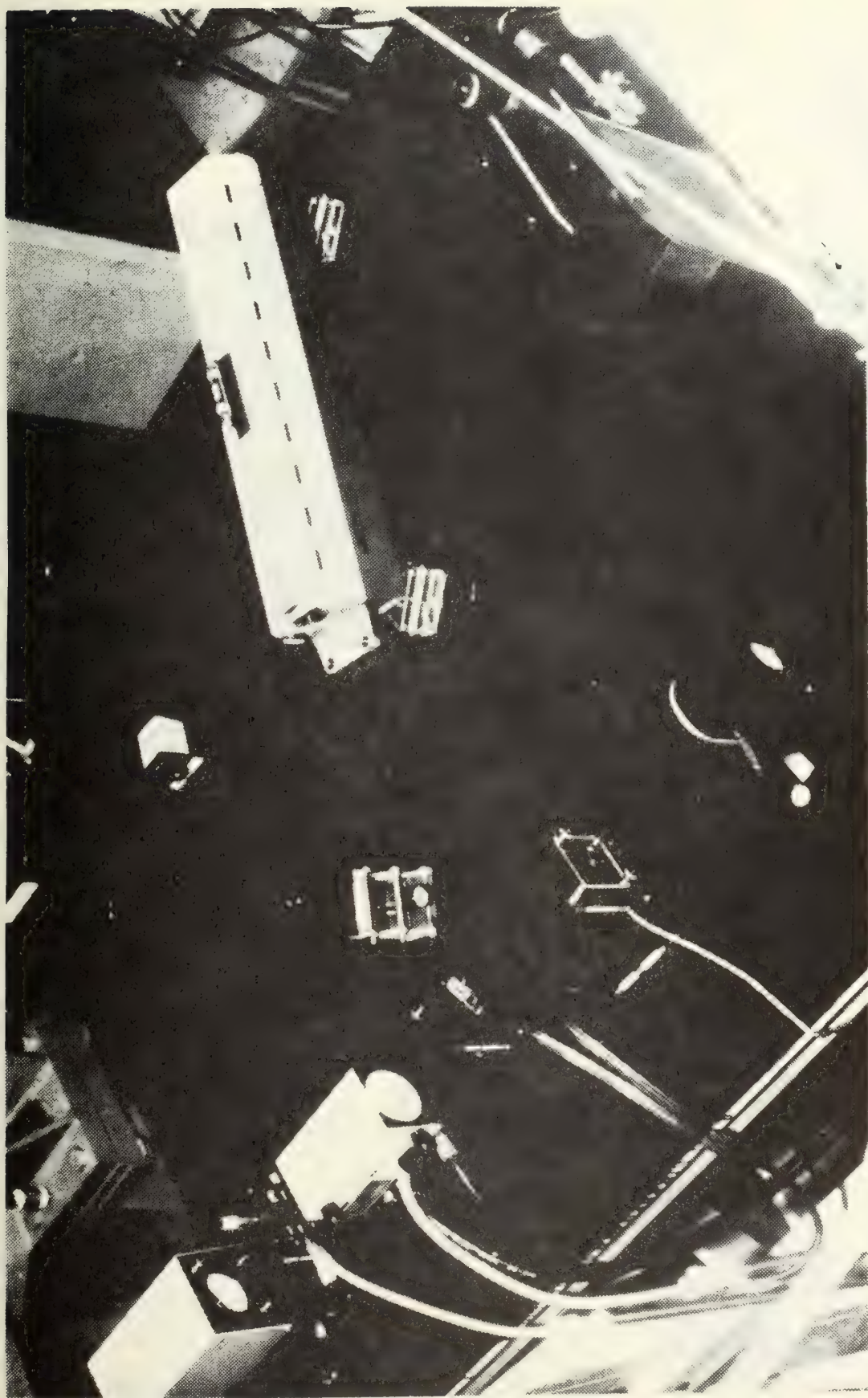


FIGURE 5. OBLIQUE VIEW OF HOLOGRAPHIC TABLE

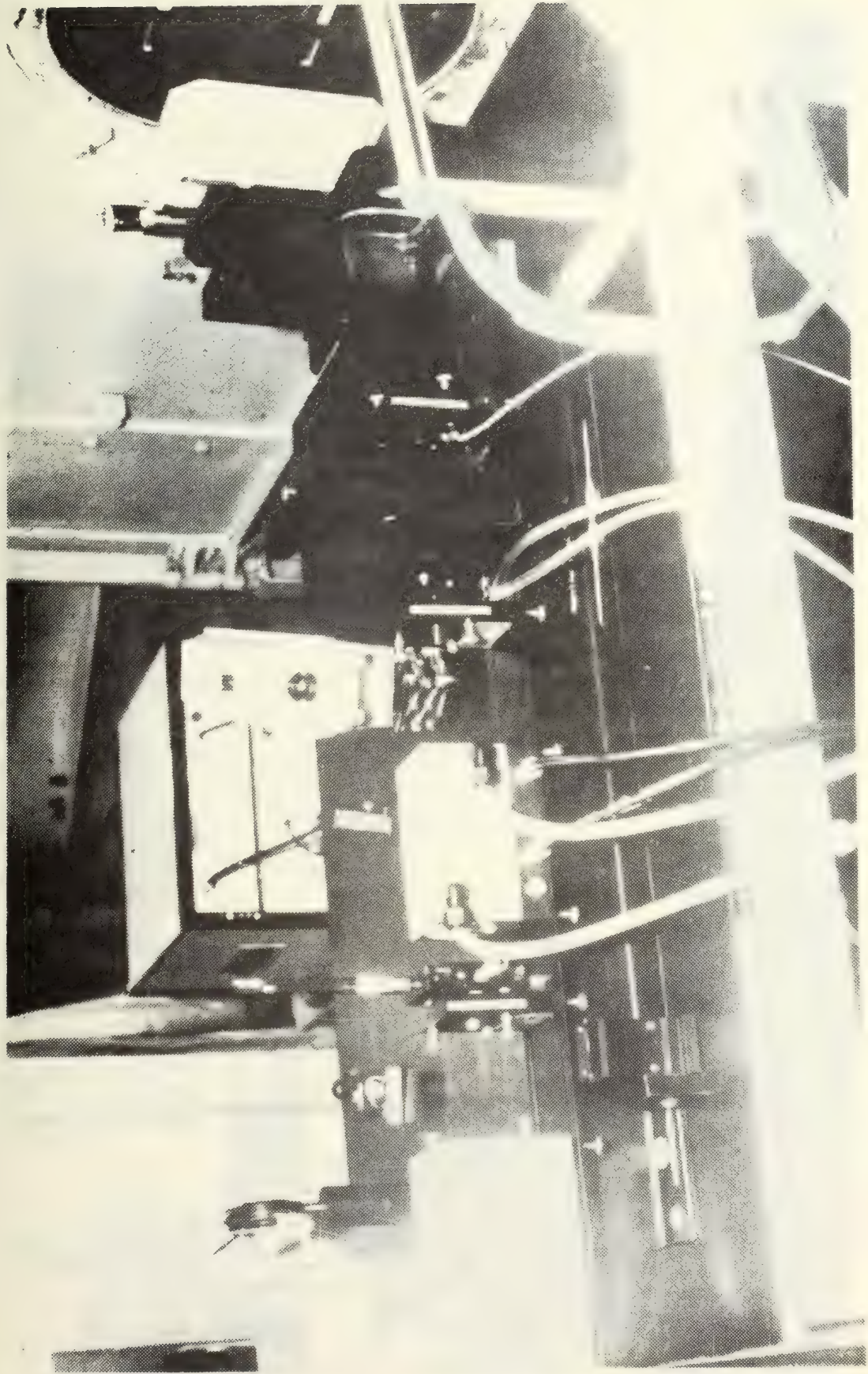


FIGURE 6. OBLIQUE VIEW OF HOLOGRAPHIC TABLE

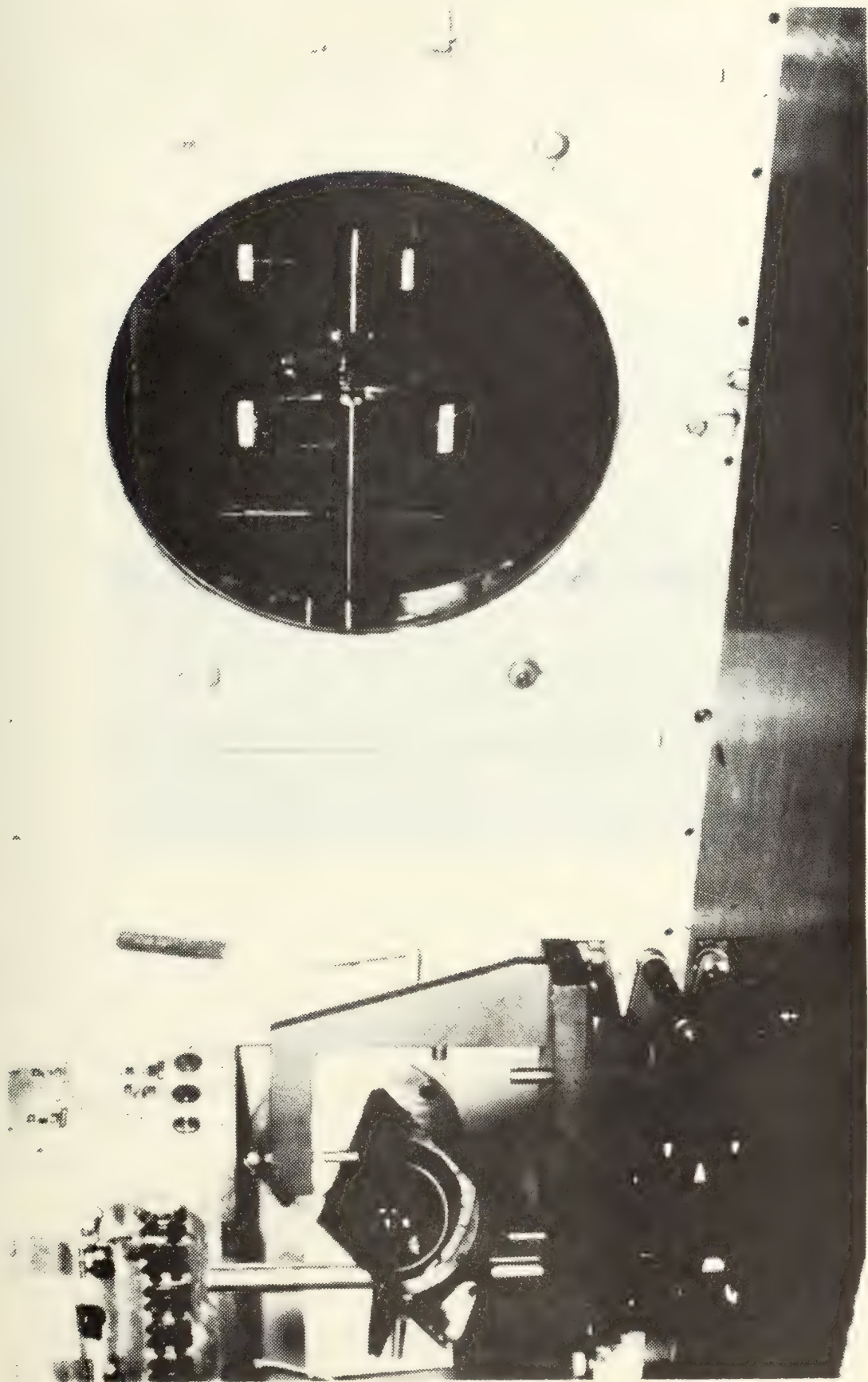


FIGURE 7. DETAIL OF MODEL MOUNTING AND REFERENCE GRIDS

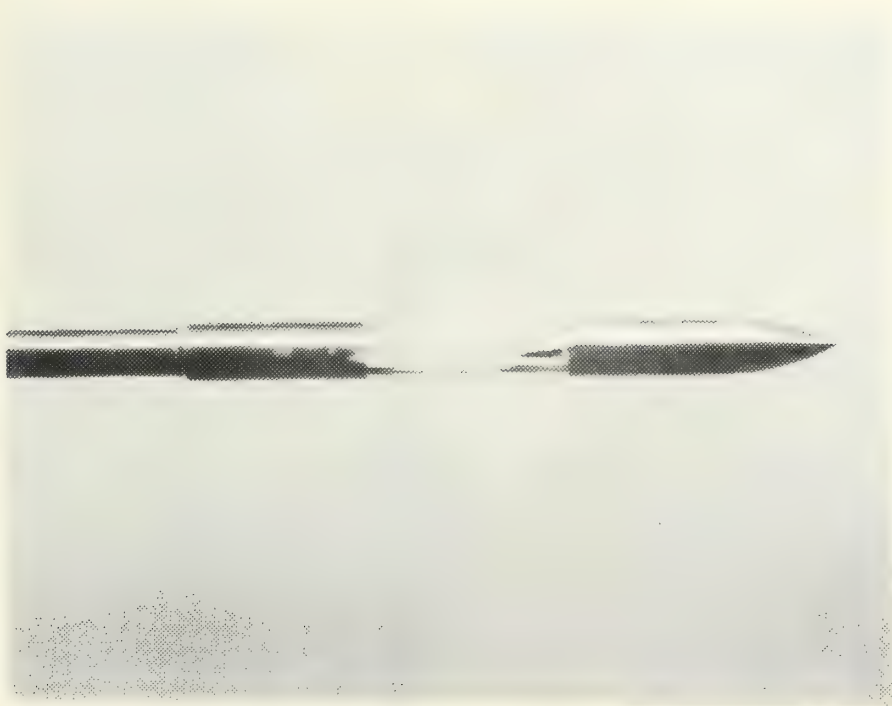


FIGURE 8. AERODYNAMIC TEST MODEL: 0 DEG. ROLL ANGLE

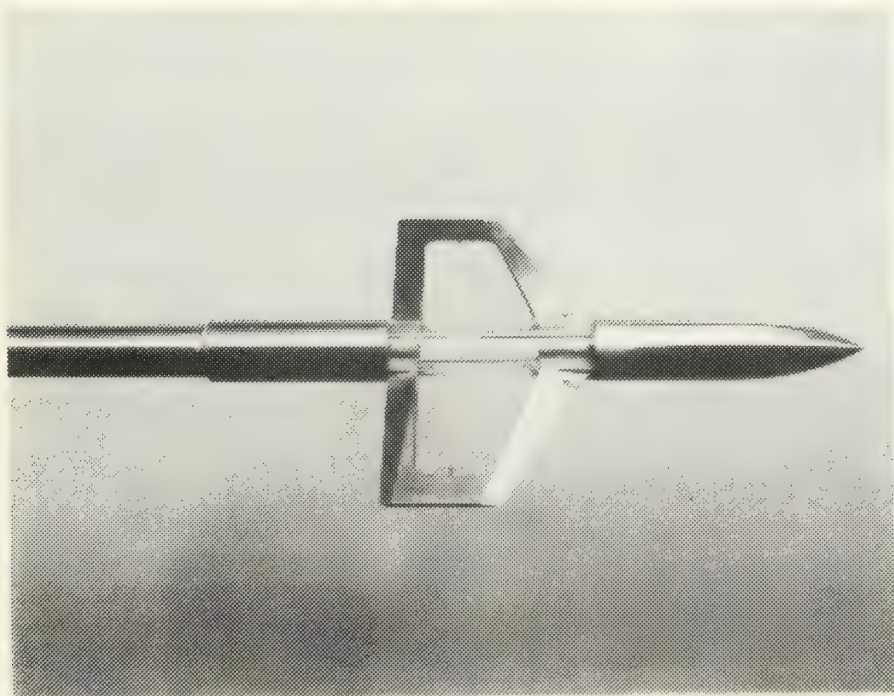


FIGURE 9. AERODYNAMIC TEST MODEL; 45 DEG. ROLL ANGLE

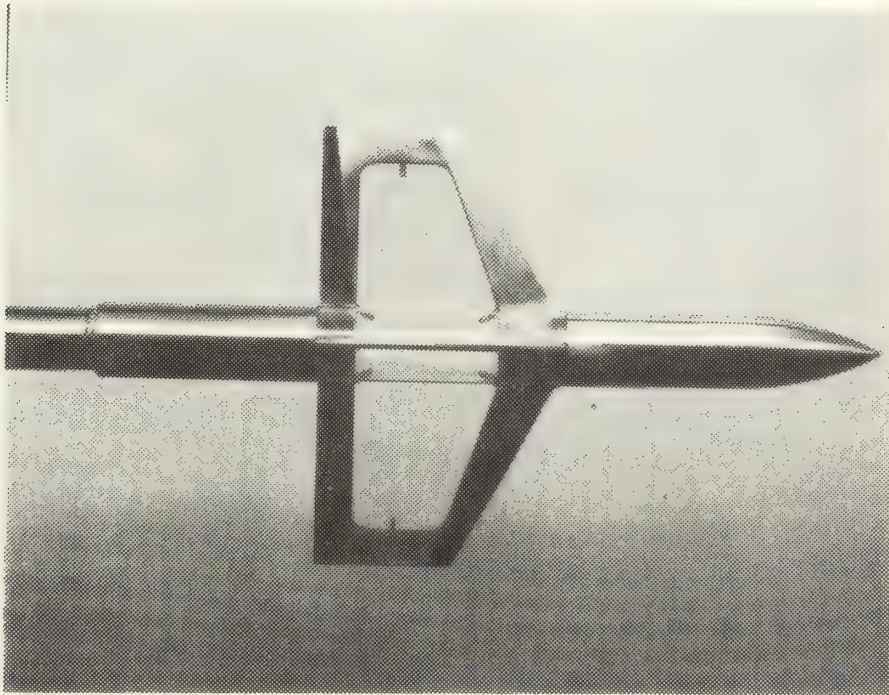


FIGURE 10. AERODYNAMIC TEST MODEL; 90 DEG. ROLL ANGLE

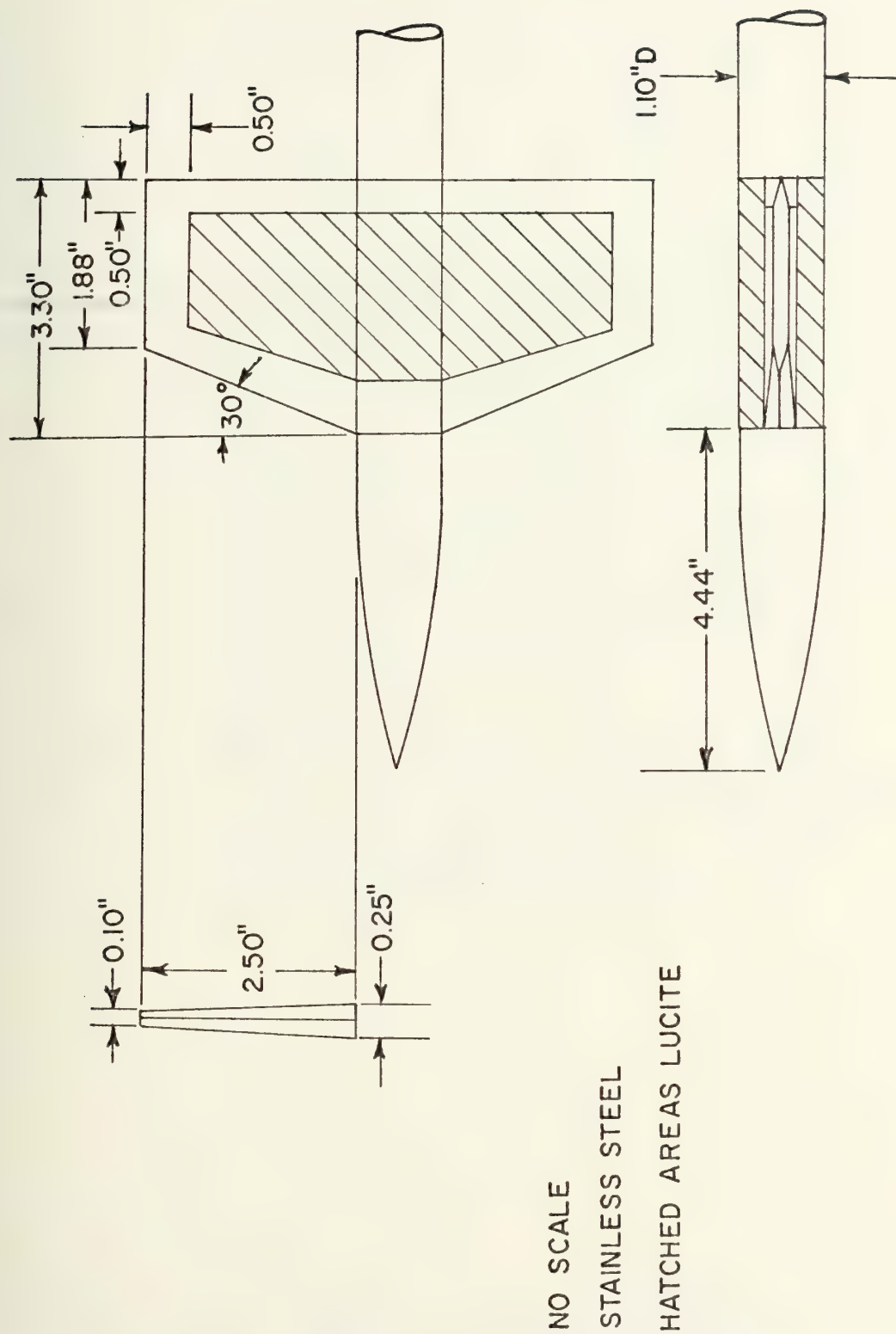


FIGURE 11. DETAILS OF THE AERODYNAMIC TEST MODEL

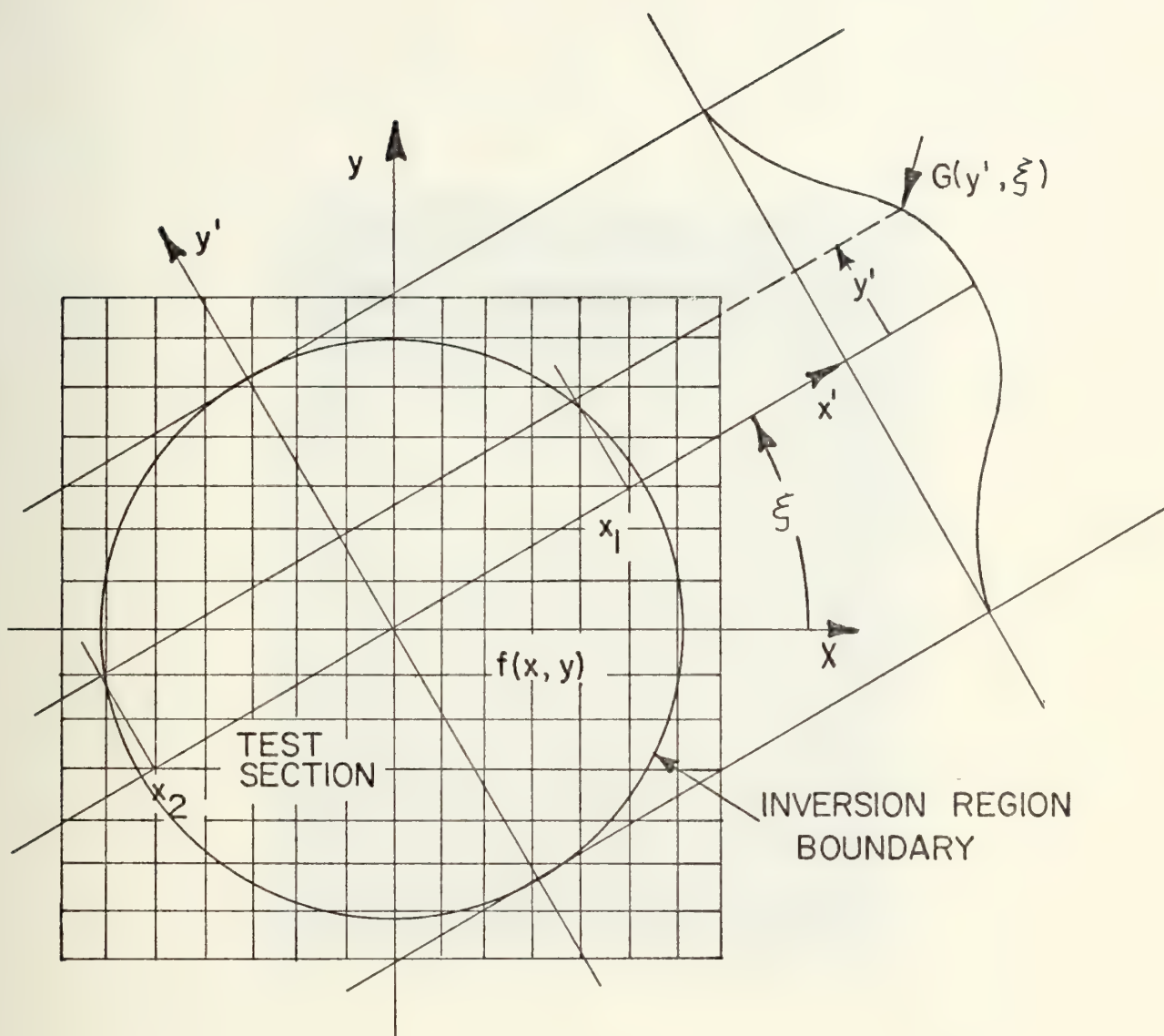


FIGURE 12. CO-ORDINATE SYSTEM USED FOR INVERSION OF FRINGE NUMBER TO DENSITY

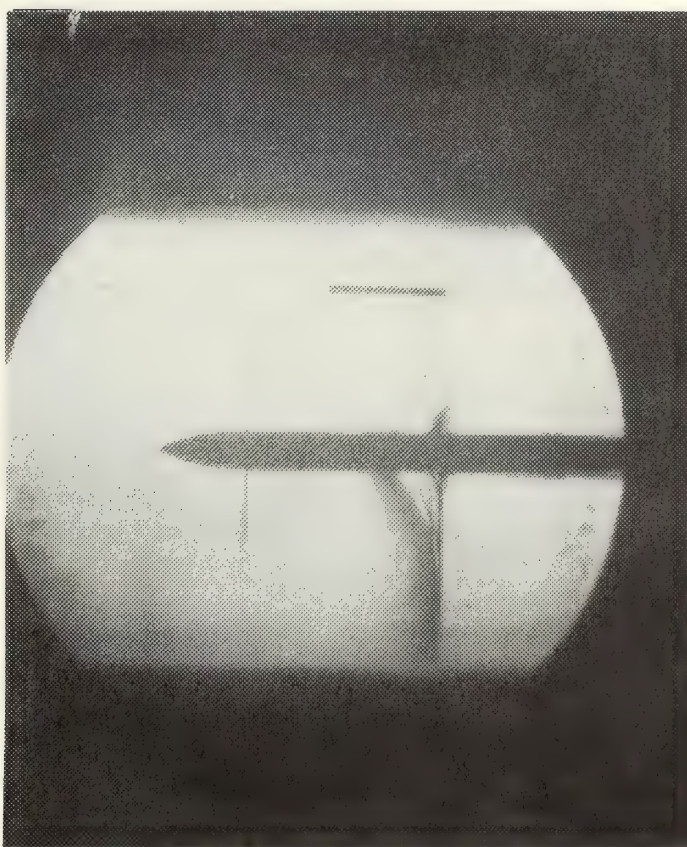


FIGURE 13. SCHLIEREN PHOTOGRAPH; 0 DEG. ROLL ANGLE,
0.967 MACH NUMBER

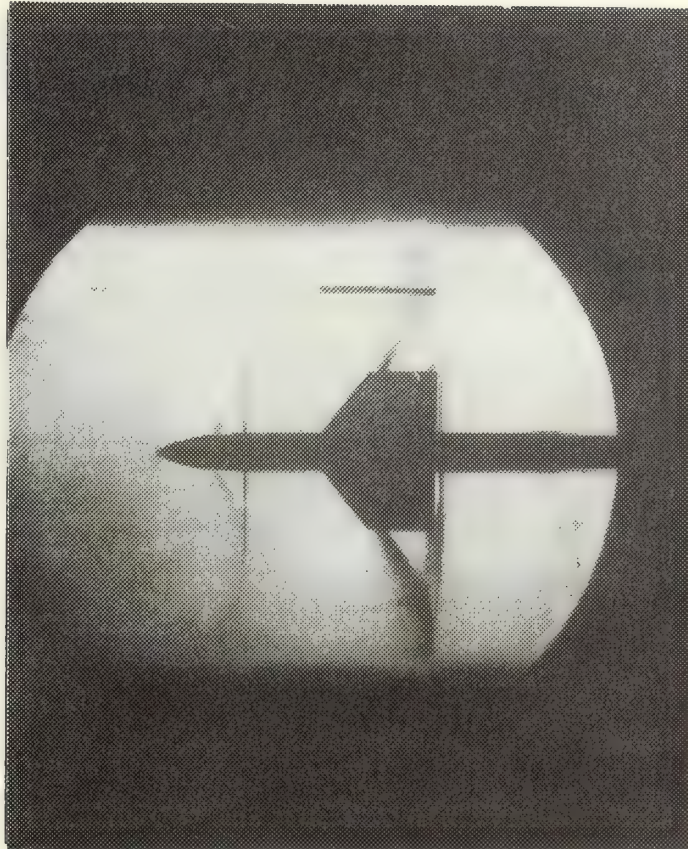


FIGURE 14. SCHLIEREN PHOTOGRAPH; 45 DEG. ROLL ANGLE.
0.967 MACH NUMBER

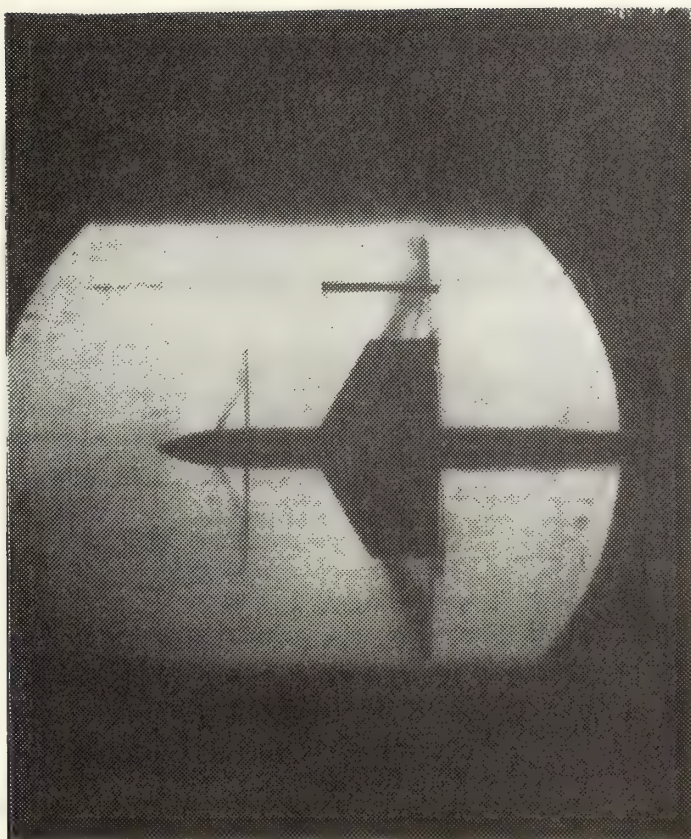


FIGURE 15. SCHLIEREN PHOTOGRAPH; 90 DEG. ROLL ANGLE,
0.967 MACH NUMBER

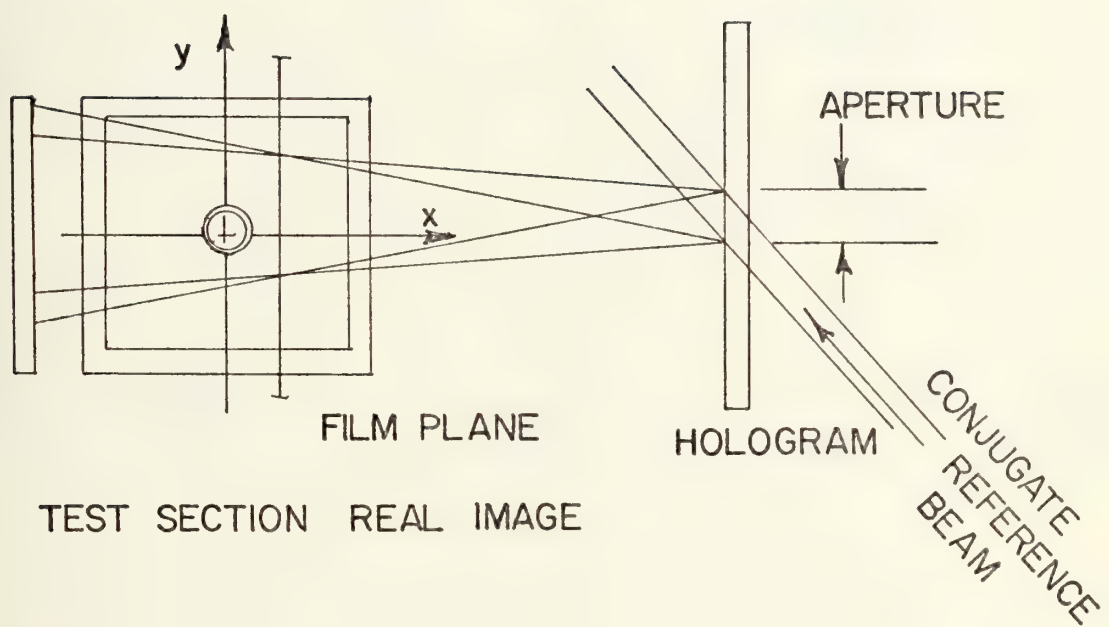


FIGURE 16. LENSLESS PHOTOGRAPHIC TECHNIQUE USING A CONJUGATE REFERENCE BEAM OF SMALL DIAMETER

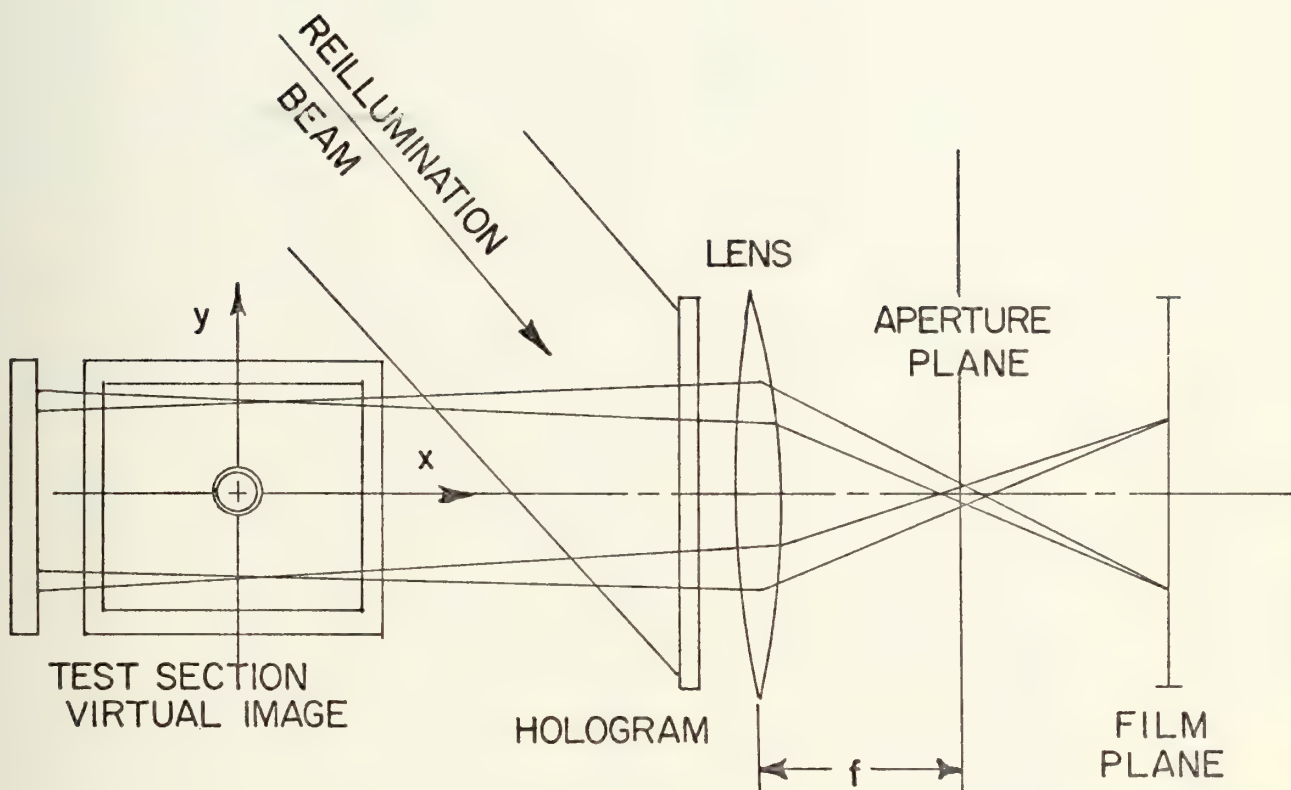


FIGURE 17. SPATIAL FILTERING TECHNIQUE FOR SELECTING PHOTOGRAPH OF CONSTANT ANGLE LINES OF LIGHT

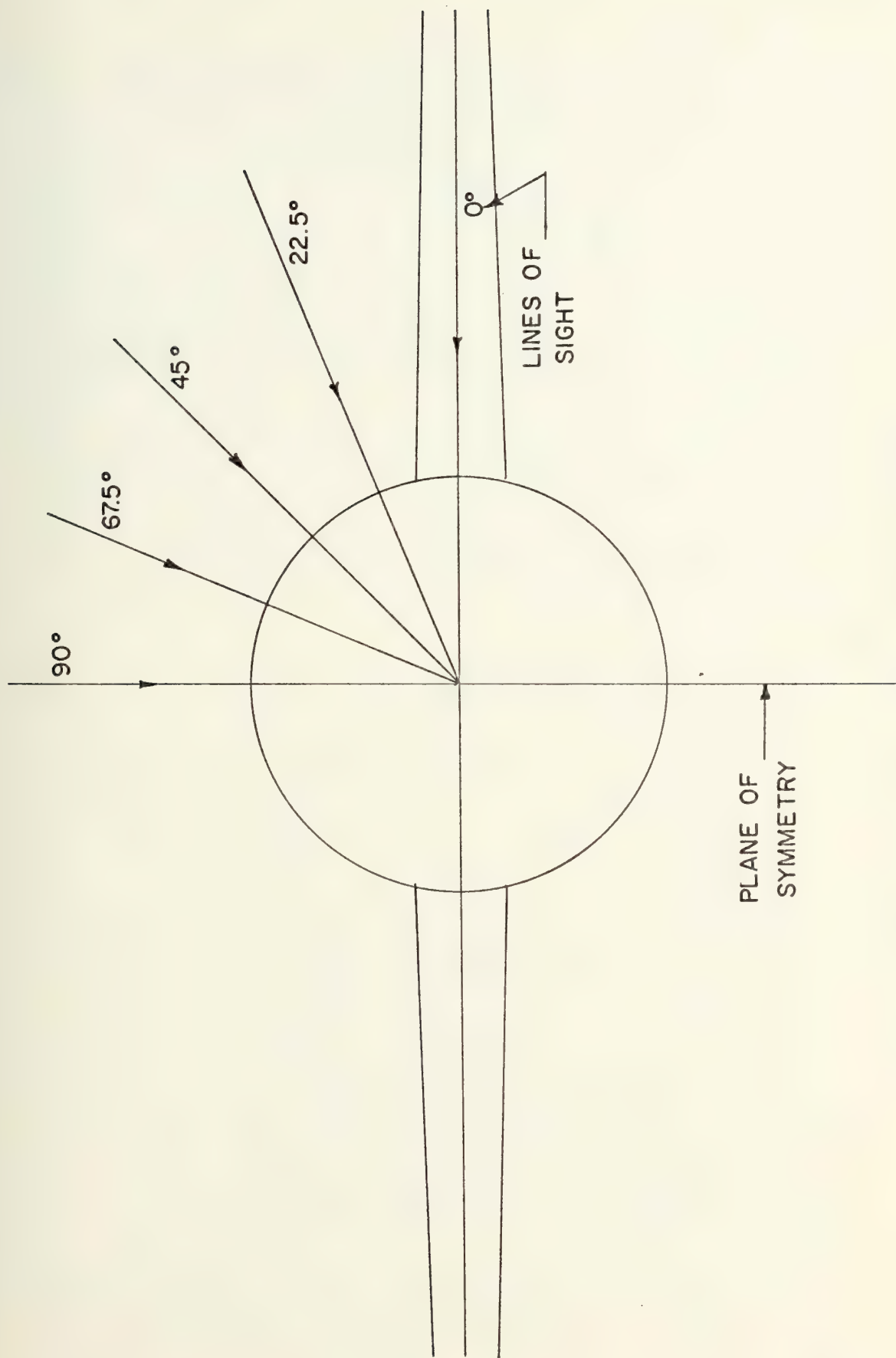


FIGURE 18. FRINGE DATA INPUT INFORMATION

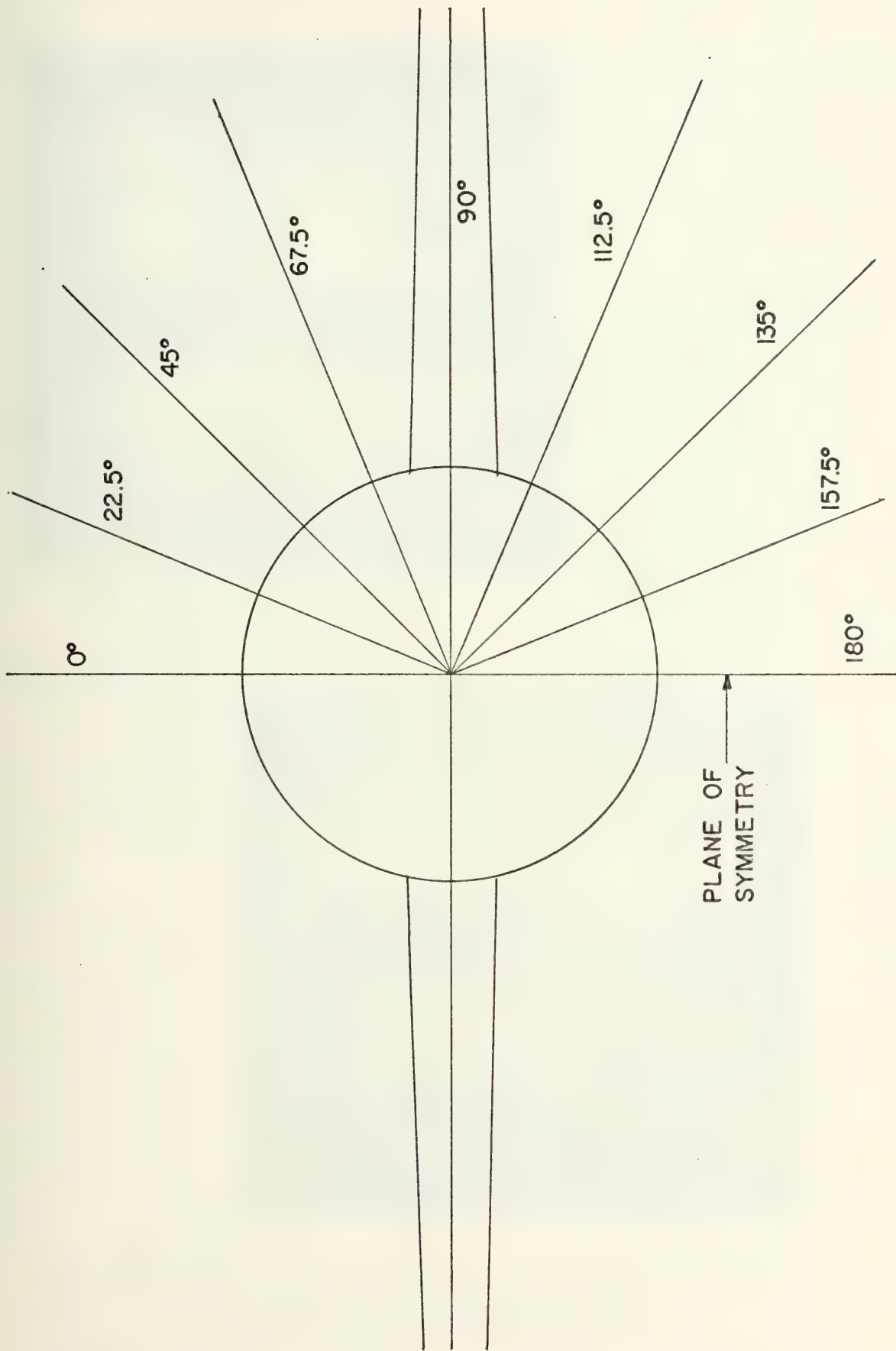
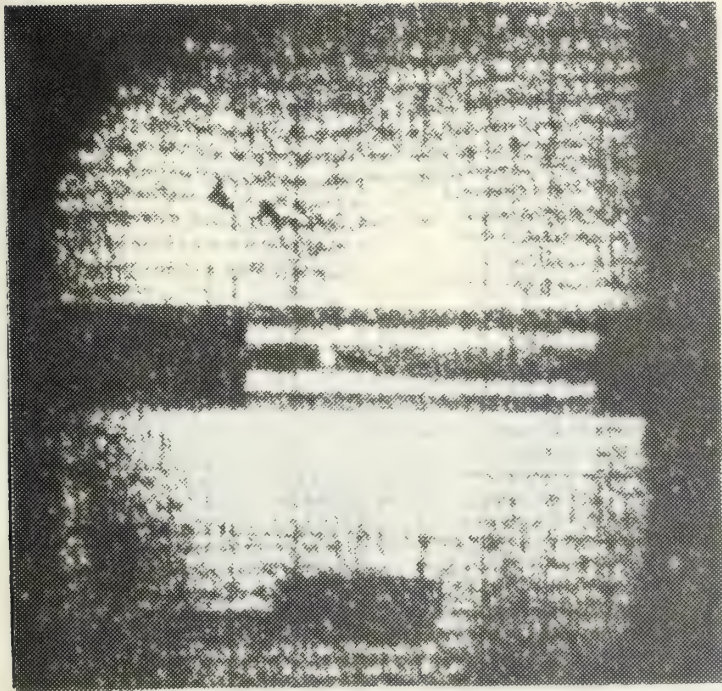


FIGURE 19. DENSITY DATA OUTPUT INFORMATION



DOUBLE-STATIC
INTERFEROGRAM

STATIC-DYNAMIC
INTERFEROGRAM

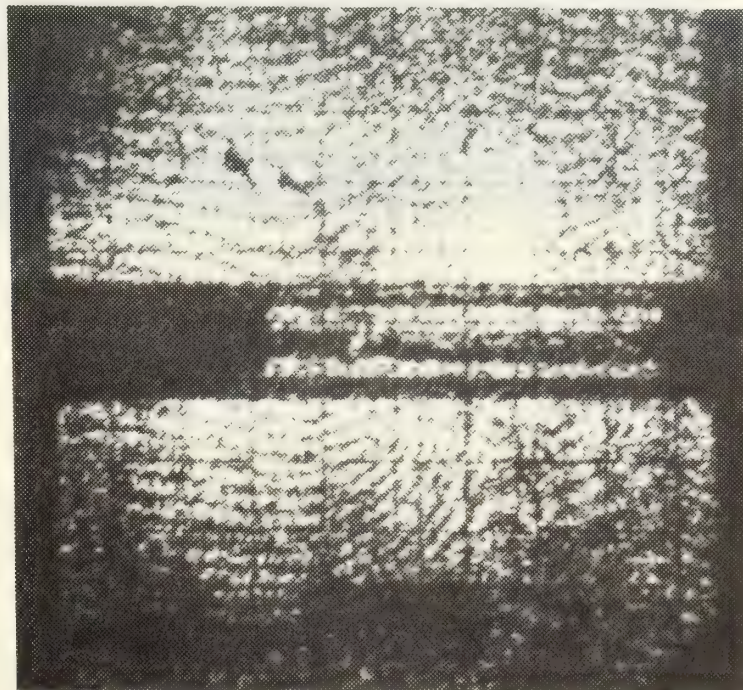


FIGURE 20. PHOTOGRAPHIC INTERFEROGRAMS
FOR 0 DEG. VIEWING ANGLE

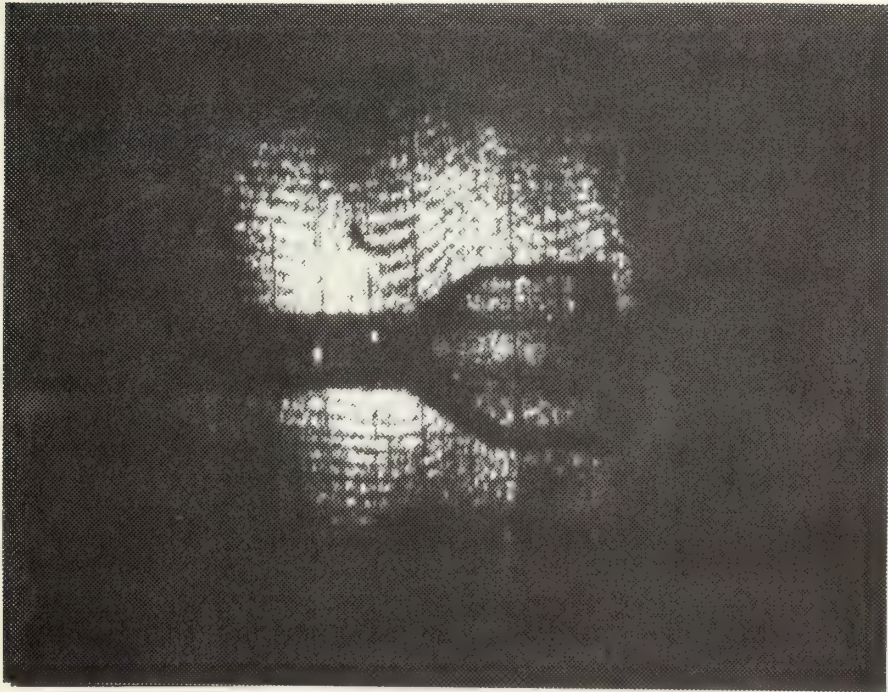


FIGURE 21. STATIC-DYNAMIC INTERFEROGRAM
FOR $22\frac{1}{2}$ DEG. VIEWING ANGLE



FIGURE 22. STATIC-DYNAMIC INTERFEROGRAM
FOR 45 DEG. VIEWING ANGLE

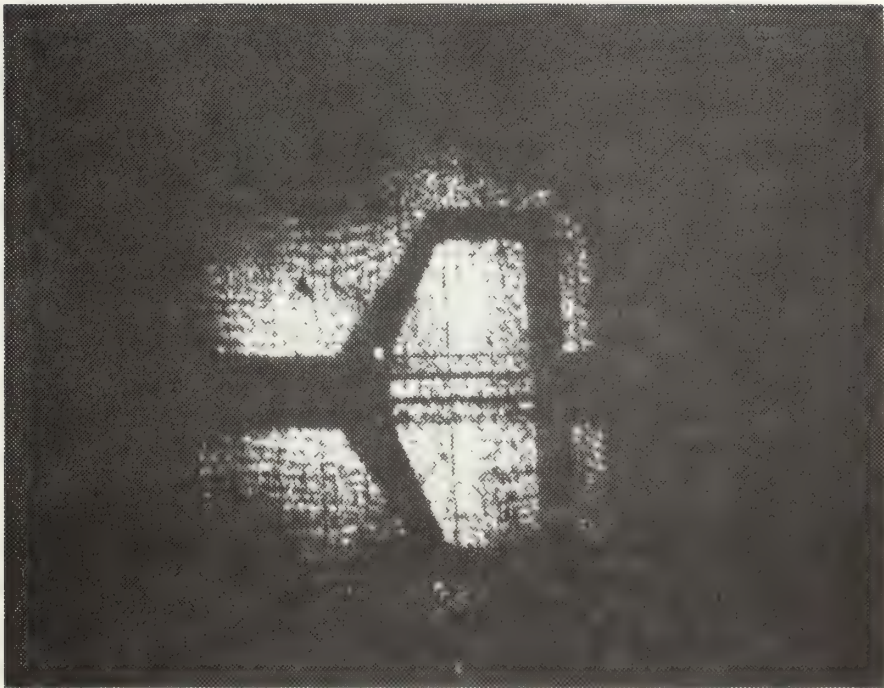


FIGURE 23. STATIC-DYNAMIC INTERFEROGRAM
FOR $67\frac{1}{2}$ DEG. VIEWING ANGLE

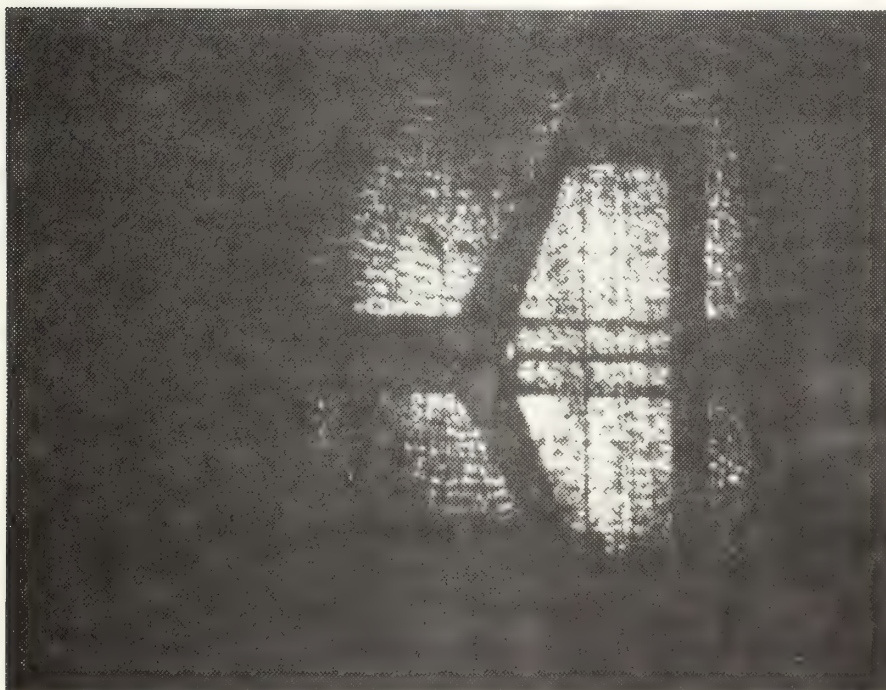


FIGURE 24. STATIC-DYNAMIC INTERFEROGRAM
FOR 90 DEG. VIEWING ANGLE

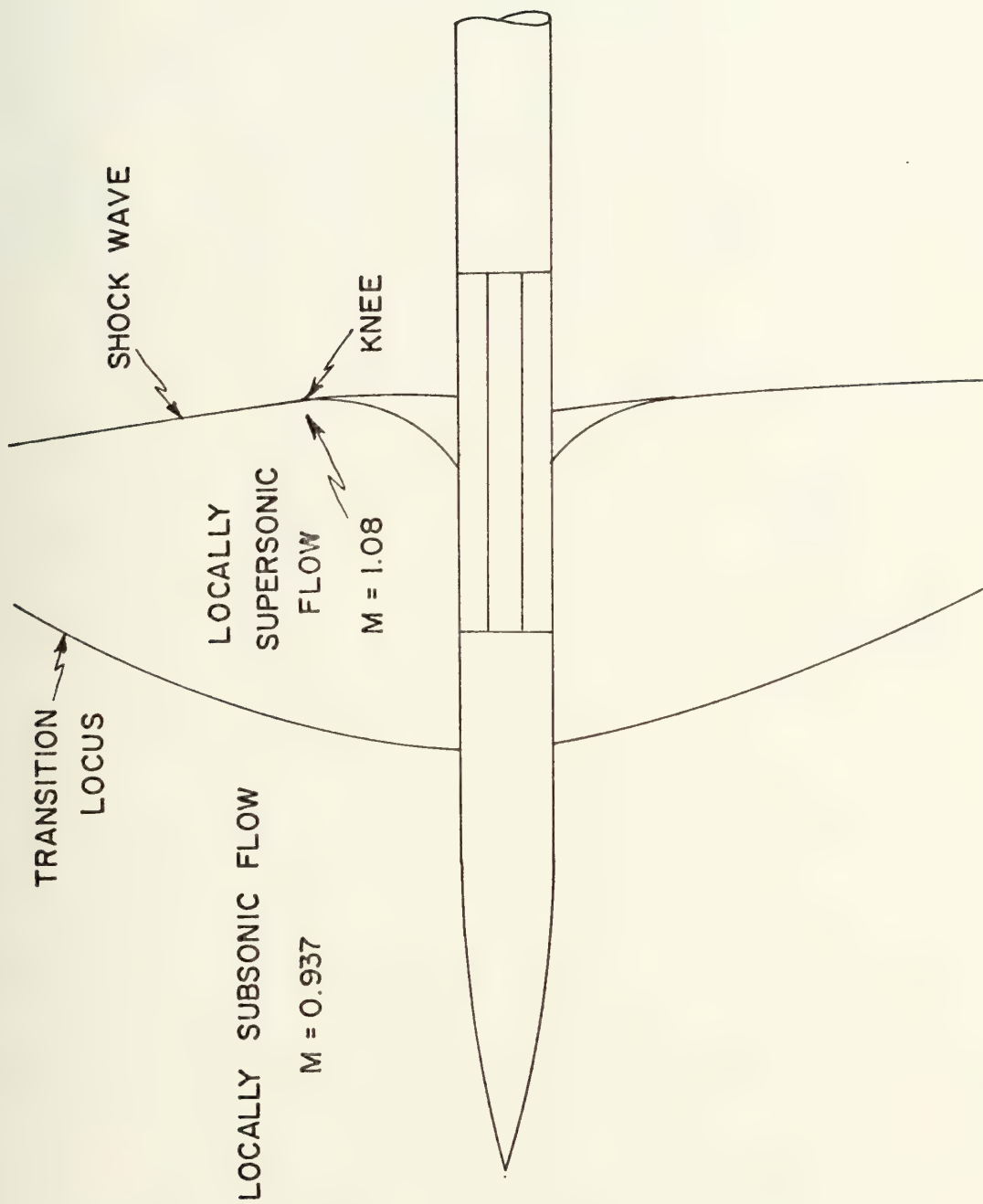


FIGURE 25.
CHARACTERISTIC TRANSONIC FLOW REGIONS; FROM SEVERAL INTERFEROGRAMS

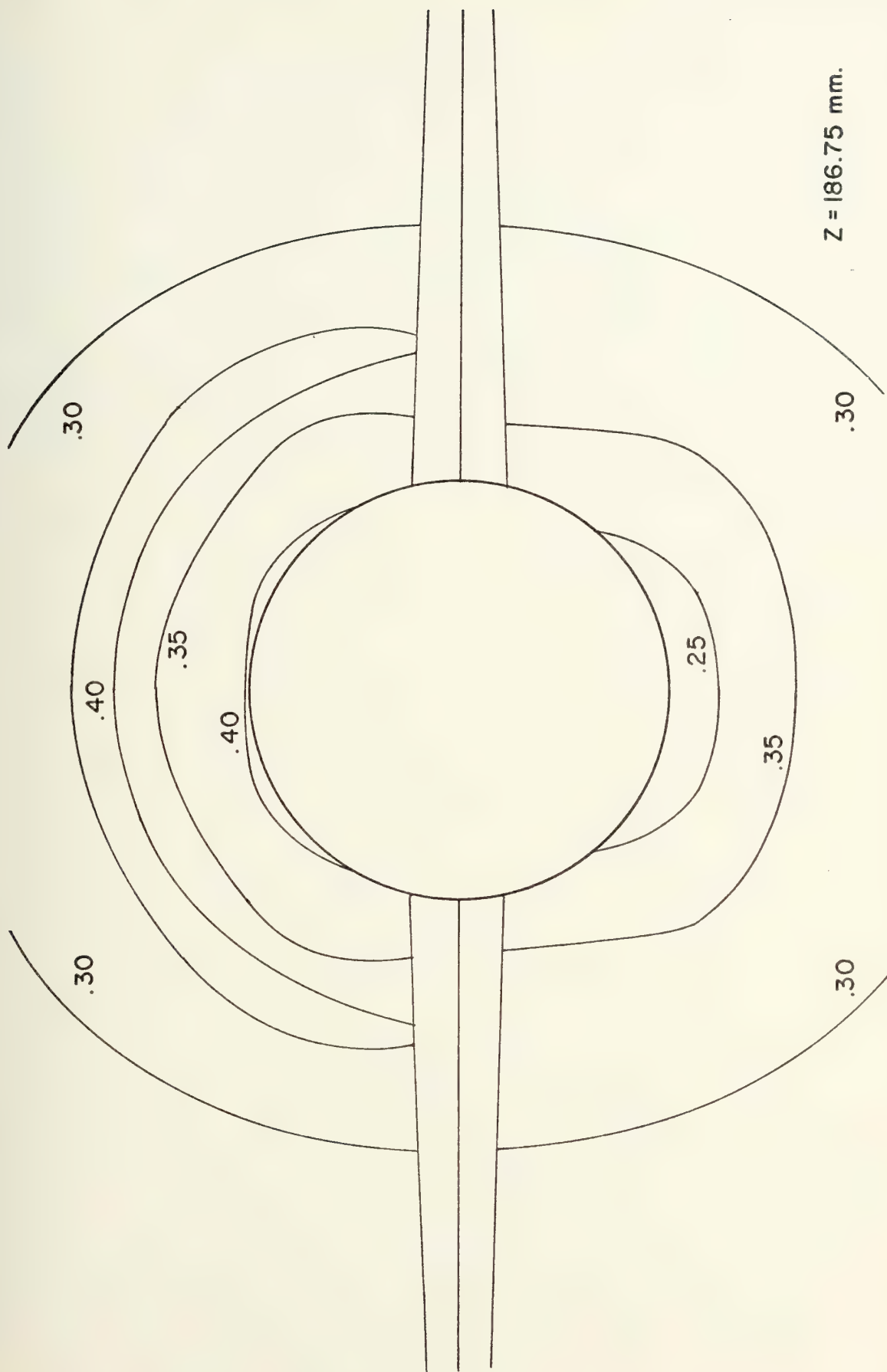


FIGURE 26. CONTOUR PLOT OF DENSITY FUNCTION, $(\rho / \rho_{\infty}) - 1$, FOR GIVEN CROSS-SECTIONAL PLANE

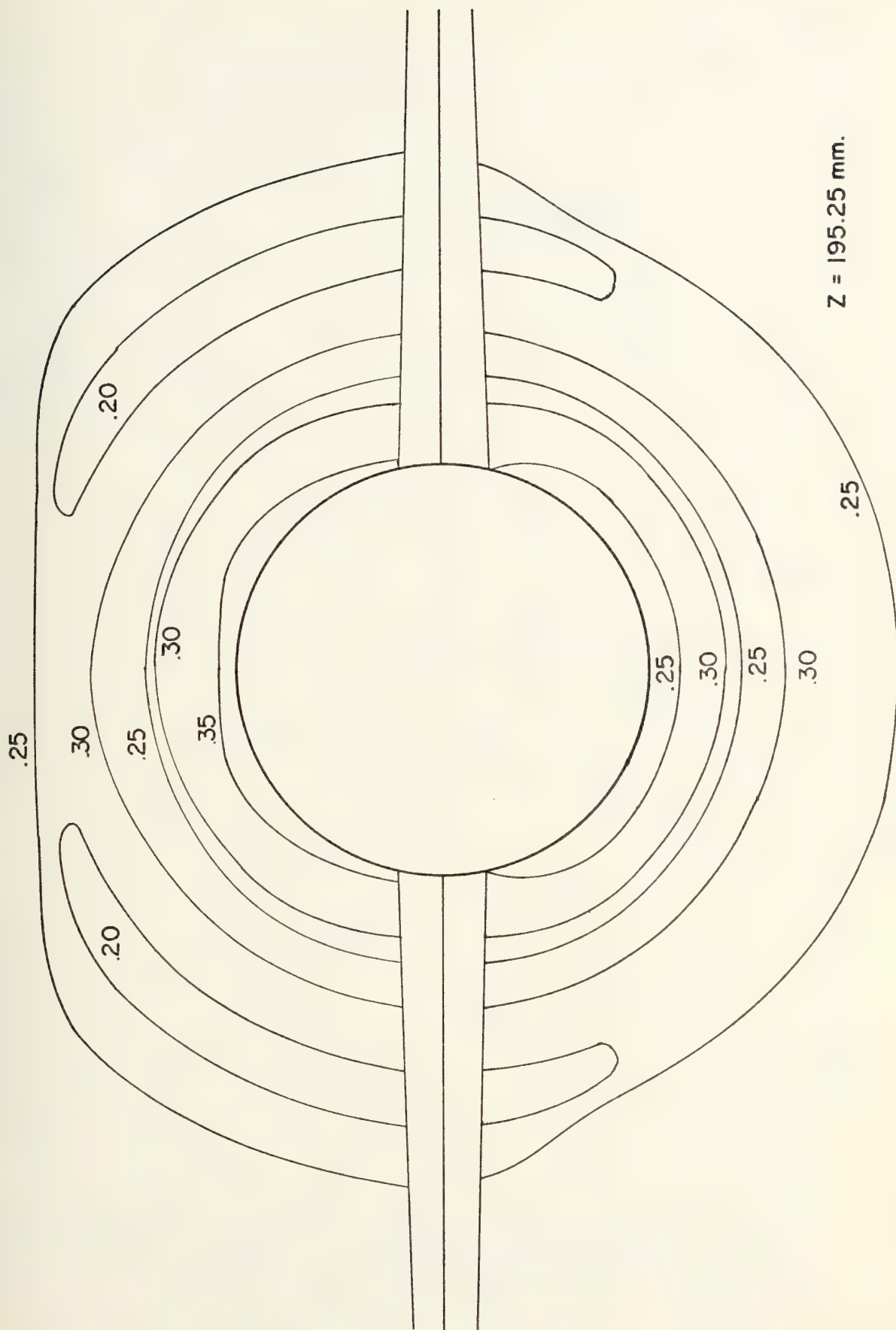


FIGURE 27.
CONTOUR PLOT OF DENSITY FUNCTION, $(\rho / \rho_{\infty}) - 1$, FOR GIVEN CROSS-SECTIONAL PLANE

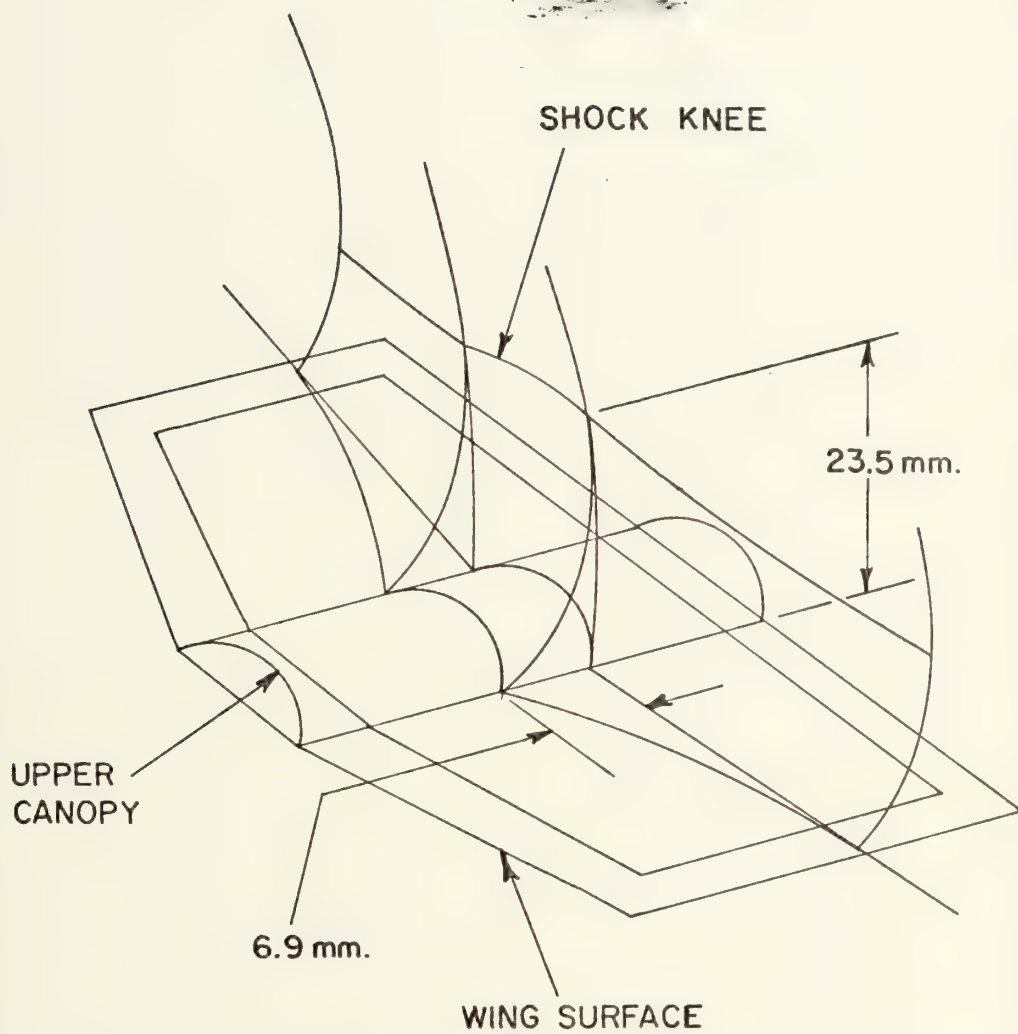


FIGURE 28. THREE DIMENSIONAL SCHEMATIC OF SHOCK WAVE STRUCTURE; CONSTRUCTED USING SEVERAL INTERFEROGRAM VIEWS

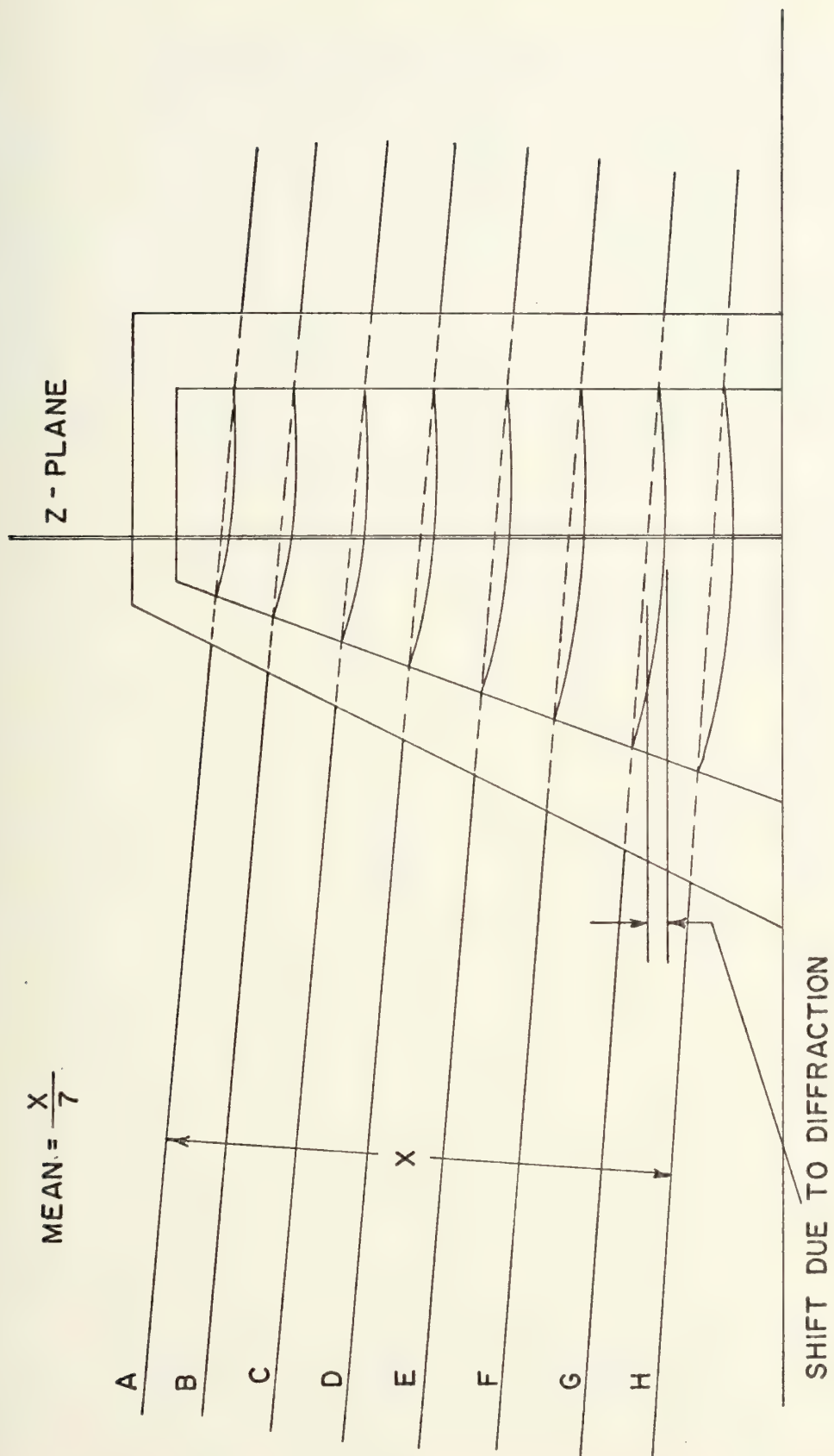


FIGURE 29. DOUBLE-STATIC HOLOGRAM REDUCTION PROCESS; $Z = 186.75$ mm.

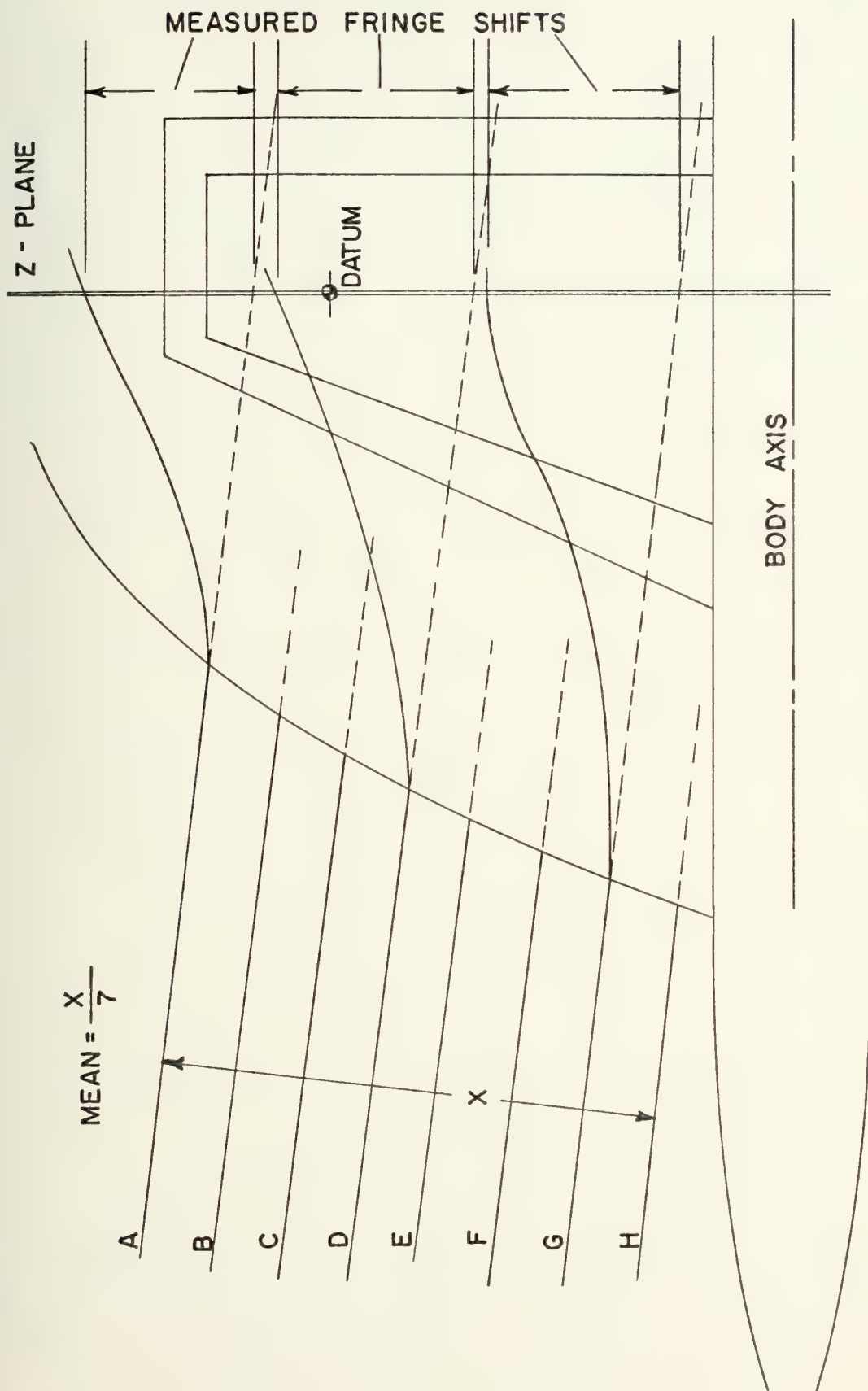


FIGURE 30. STATIC-DYNAMIC HOLOGRAM REDUCTION PROCESS; $Z = 186.75$ mm.



FIGURE 31. RADIAL VARIATION OF FRINGE NUMBER;
ZERO DEGREE VIEW, $Z = 186.75$ mm.

RUN NUMBER	HOLOGRAM NUMBER	DOUBLE EXPOSURE	ROLL ANGLE	P _A (psi)	P _T (psi)	T _T (deg. F)	MACH NUMBER
1	9	S/S	0.00	14.781	14.181	64.6	.9371
2	10	S/D	0.00				
3	23	S/S	11.25				
4	24	S/D	11.25	14.820	14.228	64.2	.9398
5	19	S/S	22.50				
6	20	S/D	22.50	14.784	14.187	64.6	.9361
7	15	S/S	33.75				
8	16	S/D	33.75	14.800	14.206	64.2	.9392
9	11	S/S	45.00				
10	12	S/D	45.00	14.785	14.185	64.6	.9391
11	25	S/S	56.25				
12	26	S/D	56.25	14.830	14.237	64.3	.9361
13	21	S/S	67.50				
14	22	S/D	67.50	14.787	14.195	64.3	.9365
15	17	S/S	78.75				
16	18	S/D	78.75	14.870	14.285	64.8	.9342
17	13	S/S	90.00				
18	14	S/D	90.00	14.787	14.195	64.6	.9402
19	35	S/S	101.25				
20	36	S/D	101.25	14.781	14.274	61.5	.9392
21	27	S/S	112.50				
22	28	S/D	112.50	14.875	14.281	62.5	.9371
23	37	S/S	123.75				
24	38	S/D	123.75	14.878	14.284	61.0	.9369
25	29	S/S	135.00				
26	30	S/D	135.00	14.875	14.278	62.0	.9372
27	39	S/S	146.25				
28	40	S/D	146.25	14.881	14.286	60.3	.9352
29	31	S/S	157.50				
30	32	S/D	157.50	14.875	14.276	61.8	.9350
31	41	S/S	168.75				
32	42	S/D	168.75	14.881	14.287	61.0	.9371
33	33	S/S	180.00				
34	34	S/D	180.00	14.875	14.282	62.0	.9357

TABLE I. EXPERIMENTAL DATA FOR WIND TUNNEL RUNS AT NSRDC

FRINGE I.D.	MEASURED SHIFT	CORRECTED SHIFT	FRINGE NUMBER	DISTANCE FROM DATUM	DEMAG. DISTANCE	DISTANCE FROM AXIS	NONDIMEN. LOCATION
A	25.0	21.475	3.53	-27.0	-25.1	-63.8	-.798
B	29.6	26.075	4.29	-23.2	-21.6	-59.8	-.748
C	30.4	26.875	4.42	-18.5	-17.2	-55.4	-.693
D	31.6	28.075	4.62	-13.0	-12.1	-50.3	-.623
E	35.5	31.975	5.23	- 9.5	- 8.8	-47.0	-.588
F	36.0	32.475	5.34	- 3.6	- 3.4	-41.6	-.520
G	34.0	30.475	5.23	+32.8	+30.5	- 7.7	-.096
H	36.4	32.875	5.64	+37.0	+34.4	- 3.8	-.048
I	40.0	36.475	6.26	+40.0	+37.2	- 1.0	-.013
J	41.5	37.975	6.51	+44.4	+41.3	+ 3.1	+0.039
K	43.5	39.975	6.86	+48.5	+45.2	+ 7.0	+0.088
L	45.3	41.775	7.17	+52.7	+49.1	+10.9	+0.136
M	47.4	43.875	7.53	+57.0	+53.1	+14.9	+0.186
N	48.6	45.075	7.73	+60.5	+56.3	+18.1	+0.226
O	46.8	43.275	7.42	+67.0	+62.4	+24.2	+0.303
P	38.5	34.975	5.99	+80.3	+74.8	+36.6	+0.458
Q	34.0	30.475	5.23	+89.6	+83.4	+45.2	+0.565

Mean free stream spacing = 5.96 mm.

Diffraction Correction = 3.525 mm.

Magnification factor = 1.074

Body axis location = +38.2 mm

All measurements in millimeters

- = above, + = below

TABLE II.

RECORDED DATA FOR ZERO DEGREE VIEW AT Z = 186.75 mm. FROM BODY NOSE

APPENDIX A

REDUCTION OF AN INTERFEROGRAM TO OBTAIN FRINGE SHIFT DATA

The reduction of data for a cross-sectional plane of interest involved a complete analysis of both doubly exposed holograms, and their corresponding interferograms, for each viewing angle.

For each view, or line of sight, the double-static exposure was first placed face down on a light table. The average fringe spacing was recorded on the back by measuring the perpendicular distance between two sufficiently separated fringe lines in the free stream and dividing by the number of spacings in between. The method is portrayed in Figure 29. Of primary importance was the measurement of the uniform shift of fringe lines due to the beam diffractive quality of the lucite sections of the model; this was taken to be the average distance between hypothetically unaltered fringe lines and the corresponding shifted fringes at their intersection with the z -plane.

The static-dynamic exposure was then placed face down on the light table and the fringe and model contours traced on the back, as shown in Figure 30. Again, the average free stream fringe line spacing was measured and checked against the value from the double-static exposure. A reference point, or datum, was selected at the intersection of a well-defined horizontal grid line and the z -plane of interest. Fringe shifts were computed in the following manner:

1. The distance from datum to a hypothetically undeviated fringe line at its intersection with the z-plane was measured in millimeters.
2. The distance from datum to the actual deviated, or shifted, fringe at its intersection with the z-plane was measured in millimeters. Measured shifts should be made perpendicular to the fringe direction. The present procedure is convenient and only introduces an error of about 1% in the overall level of the density field.
3. The raw fringe shift distance was then the distance in 1. above minus the distance in 2.
4. To correct for diffraction error, the uniform shift measured in the double-static exposure was then subtracted from the distance in 3.

Fringe numbers were calculated by dividing the shift for each fringe by the average free stream spacing for the exposure. Magnification factors were computed for each exposure by comparing photographically recorded model diameter with actual model diameter.

Since the datum location varied slightly from exposure to exposure, it was necessary to reference all measurements to a central point for each plane of analysis. This point was taken to be the intersection of the longitudinal axis of the model fuselage and the cross-sectional plane of interest. The intersections of shifted fringe lines with the z-plane were referenced to this fixed point and demagnified. The resulting distances were then nondimensionalized using the selected inversion circle radius. Table II contains typical data recorded for the zero degree view at the 186.75 mm. plane of interest.

A plot of fringe numbers versus nondimensionalized fringe location was produced for each viewing angle using data obtained from the interferograms. Fringe numbers at 201 equidistant points, as required for input into Mode 3 of the inversion computer program, were recorded from the resulting curves. The curve plotted using the data from Table II is shown in Figure 31.

APPENDIX B

APPLICATION OF COMPUTER PROGRAM "HOLOFER"

The computer program used in this study is designed to invert an array of fringe numbers across a field to obtain the associated density field using a special adaptation of the inversion technique first proposed by C. D. Maldonado [5,6]. Three different modes of operation are available to the operator, as described below:

(a) Mode 1

This mode provides the basic self-testing capability of the computer program. It can either generate its own input density field using Subroutine FUNCT or read in a density field through Subroutine FREAD. The fringe number array corresponding to the input density data is first generated; this information is then used as the input for the inversion to obtain the original density distribution once again. This self-testing procedure was utilized in the present investigation to determine the best value of the scale factor, α , required to assure accurate density in the region of inversion.

(b) Mode 2

This mode generates the fringe array at regular intervals across the test field from irregularly spaced fringe values read in through Subroutine SHEET. Simulated fringe arrays may be generated

by one of the functions specified in Subroutine FUNCT if NCODE = 1 is specified. Mode 2 has not been utilized in the present investigation.

(c) Mode 3

Mode 3 operation reads in a regularly spaced array of input density values directly utilizing Subroutine READ, which is first called by Subroutine GARRAY. The parameters in the first two cards preceding the input fringe data serve to identify the size, location and symmetry of the fringe field. The following parameters were used in calculating the asymmetric density field in the present experiment:

<u>PARAMETER</u>	<u>INPUT</u>
NOF	Run Number
IMAX	201
JMAX	20
ISYM	1
JSYM	1
IMS	201
JMS	5
Z	0.0, 1.0
XO	0.0
YO	0.0
PHISYM	0.0

Amplifying details and applications of the inversion computer program are found in References [3, 7, 9]. A listing of the program is included in this appendix for reference.

NPTS=AR(15)
 NLINS=AR(16)
 SD=AR(18)
 PHIZ=AR(19)
 DELPHI=AR(20)
 YPZERO=AR(21)
 YPRNG=AR(22)
 XPZERO=AR(23)
 XPRNG=AR(24)
 NOF=AR(25)
 NAF=AR(26)
 IPT=AR(27)
 KPT=AR(28)
 LPT=AR(29)
 BND=AR(30)
 A=AR(31)
 B=AR(32)
 C=AR(33)
 D=AR(34)
 E=AR(35)
 P=AR(36)
 S=AR(37)
 T=AR(38)
 U=AR(39)
 V=AR(40)
 W=AR(41)
 Q=AR(42)

WRITE(6,90)
 WRITE(6,98)
 WRITE(6,91)
 WRITE(6,92)
 WRITE(6,93)
 WRITE(6,94)
 WRITE(6,95)
 WRITE(6,96)
 WRITE(6,97)
 WRITE(6,98)
 FORMAT(//)

CAL00480
 CAL00490
 CAL00500
 CAL00510
 CAL00520
 CAL00530
 CAL00540
 CAL00550
 CAL00560
 CAL00570
 CAL00580
 CAL00590
 CAL00600
 CAL00610
 CAL00620
 CAL00630
 CAL00640
 CAL00650
 CAL00660
 CAL00670
 CAL00680
 CAL00690
 CAL00700
 CAL00710
 CAL00720
 CAL00730
 CAL00740

90
 91
 92
 93
 94
 95

```

    1, MEXTRA * KEXTRA ** SET00310
    1, BETA * EPS * RHO-INF * LAMBDA ** SET00340
    1, STD DEV * F11.6, 2X, F8.7, F9.1, F11.6) * POINTS * LINES * DIAGNOS * SET
    1, XPRNGE * YPZERO * YPRNGE * XPZERO **
    1, MAP BND * ADD FUN * GARRAY * GRAPH * LIN PRT ** SET00400
  
```



```

96  FORMAT (/5X,*, A *, B *, C *, D *, E *, SET00420
97  FORMAT (/5X,*, S *, T *, U *, V *, W *, SET00440
98  FORMAT (/3X,75A1)
    IF (DGN.GE.4) WRITE (6,89) (AR(I),I=1,42)
    NNN=2
    IF (MODE.LT.0) NNN=1
    IF (MODE.GT.5) NNN=3
    IF (MODE.GT.5) MODE=MODE-10
    NGP=0
    IF (KLIMIT.LT.KEXTRA) KEXTRA=KLIMIT
    IF (MLIMIT.LT.MEXTRA) MEXTRA=MLIMIT
    IF (IPT.LT.0) NGP=IPT
    IF (IPT.LT.0) IPT=-IPT
    ISYM=2.1-(FLOAT(JSYM)/2.-FLOAT(JSYM/2))*2
    IF (JSYM.EQ.0) ISYM=1
    IF (JSYM.GT.JMAX) ISYM=2
    IF (ISYM.EQ.1) JMAX=((JMAX+1)/2)*2
    RJMX=JMAX
    MSYM=JSYM
    IF ((MSYM.EQ.0).OR.(MSYM.GT.JMAX)) MSYM=1
    FCU=ISYM*JSYM#JMAX
    IF ((JSYM.GT.JMAX).OR.(JSYM.EQ.0)) FCU=JMAX
    QSYM=FCU/RJMX
    IMS=(IMAX+ISYM-1)/ISYM
    JMS=JMAX
    IF (ISYM.EQ.1) JMS=(JMAX/2+1)/2
    IF (JSYM.EQ.0) JMS=JMAX/2
    MODE=ABS(AR(13))
    XO=0.
    YO=0.
    ZD=0.
    PHISYM=0.
    HS=SIZE/2.
    RHOS=1.286
    BOX=RHOFNF*BETA/RHOS/RLAMDA
    RPTS=NPPTS
    XPR=0.
    IF (NPPTS.GT.1) XPR=XPRNG/(RPTS-1.)/2.
    XPM=-XPR
    PIE=3.141592653589793
    MONE=1
    WRITE (6,58) IMAX,JMAX,IMS,JMS,ISYM,JSYM,MSYM,QSYM,FCU
    FORMAT (3X,*, IMAX,JMAX,IMS,JMS,ISYM,JSYM,MSYM,QSYM,FCU,/,
    7I5,2F7.3/)
1  NTWO=2
    IX=IMAX+JMAX

```

58


```

NF=INI
IF ((MODE.EQ.1.) .AND. (NOF.EQ.8) .AND. (DGN.GE.1.)) WRITE (6,69)
IF ((MODE.EQ.1.) .AND. (NOF.EQ.8)) CALL FREAD (NO,RO,NF,ZD)
Z=ZD
IF (DGN.GE.1.) WRITE (6,68)
CALL GARRAY (G,GA,NOF,DGN,MONE,XO,YO,PHISYM)
LM=1
IF ((LPT.EQ.0) .AND. (BND.EQ.0)) LM=0
IIMX=IMAX+1
JIMX=JMAX+1
IJMX=IMAX*JMAX
NBD=1
IF (JSYM.EQ.0) NBD=2
KBD=KLIMIT*NBD
DO 15 IJ=1,IJMX
GA(IJ)=0.
IF (NAF.EQ.0) GO TO 16
NF=IN2
IF ((NAF.EQ.8) .AND. (DGN.GE.1.)) WRITE (6,69)
IF (NAF.EQ.8) CALL FREAD (NA,RA,NF,ZD)
MST=MODE
MODE=1
IF (DGN.GE.1.) WRITE (6,68)
IF (NAF.NE.0) CALL GARRAY (GA,G,NAF,DGN,NTWO,XO,YO,PHISYM)
MODE=MST
DO 6 IJ=1,IJMX
G(IJ)=G(IJ)+GA(IJ)
RLINS=NLINS
IF (NAF.EQ.8) WRITE (6,88) NA,(RA(L),L=1,NA)
IF (NOF.EQ.8) WRITE (6,87) NO,(RO(I),I=1,NO)
IF (LM.EQ.0) GO TO 14
RB(1)=-1.7
DO 1 I=2,7
RB(I)=RB(I-1)+.5
TPIE=2.*PIE
MPIE=-PIE
DYP=0.
DXP=0.
IF (NLINS.GT.1) DYP=YPRNG/(RLINS-1.)
IF (NPTS.GT.1) DXP=XPRNG/(RPTS-1.)
IF (DGN.GE.1.) .AND. (NPN.EQ.2)) WRITE (6,64)
IF (NPN.EQ.2) CALL BDGEN (G,H,SCF,DGN,NBD,BDA,KBD)
DO 5 J=1,NLINS
IF (DGN.GE.1.) WRITE (6,67) J
RJM=J-1
PHI=PHIZ+DELPHI*RJM
YP(J)=YPZERO+DYP*RJM
PSI=(PHI+90.)*PIE/180.

```

15

6
16

1

CAL01180
CAL01190
CAL01200
CAL01210
CAL01220
CAL01230
CAL01240
CAL01250
CAL01260
CAL01270
CAL01280
CAL01290
CAL01300
CAL01310
CAL01320
CAL01330
CAL01340
CAL01350
CAL01360
CAL01370
CAL01380
CAL01390
CAL01400
CAL01410
CAL01420
CAL01430
CAL01440
CAL01450
CAL01460
CAL01470
CAL01480
CAL01490
CAL01500
CAL01510
CAL01520
CAL01530
CAL01540
CAL01550
CAL01560
CAL01570
CAL01580
CAL01590
CAL01600
CAL01610
CAL01620
CAL01630
CAL01640
CAL01650


```

TAU=PSI-PHISYM GO TO 9
IF (LPT.EQ.0) WRITE (6,78) (ST,I=1,124)
IF (LPT.LE.1) WRITE (6,74) (ST,I=1,95)
IF (LPT.GT.1) WRITE (6,74) (ST,I=1,95)
IF ((CMS.EQ.1.) .AND. (LPT.GT.1)) READ (5,79) ZZ
WRITE (6,86)
WRITE (6,85) Z, PHI, YP(J)
WRITE (6,76)
IF (MODE.EQ.1) WRITE (6,83) (RB(I), I=1,7)
IF (MODE.GT.1) WRITE (6,80) (RB(I), I=1,7)
WRITE (6,81) (DH, I=1,54), (PL, I=1,13)
IC=0
DO 3 I=1,NPTS
RIM=I-1
THEO(I,J)=0.
CA(I)=0.
FA(I,J)=0.
ERR(I)=0.
CALC(I,J)=0.
RHO(I)=0.
XP(I)=XPZERO+DXP*RIM
XPI=ABS(XP(I))
IF (XPI.LT.1.E-13) XP(I)=0.
RS=SQRT(XP(I)**2+YP(J)**2)/HS
IF (RS.GT.1.) GO TO 13
THT=ATANM(YP(J),XP(I))
IF (XPI.EQ.0.) THT=0.
SIG=TAU-PIE/2.+THT
IF (SIG.GT.PIE) SIG=SIG-TPIE
IF (SIG.LT.MPIE) SIG=SIG+TPIE
SIGI=SIG
XS=RS*COS(SIG)
IF (DGN.GE.1.) WRITE (6,44) SIG
FORMAT (1, SIG=, E10.3)
YS=RS*SIN(SIG)
IF (DGN.GE.5) WRITE (6,57) PHI, DELPHI, PSI, TAU, THT, SIG, SIGI, XS, YS
I FORMAT (1, ANGLES=, 10E10.3)
RI=I
R=0.
IF (DGN.GE.2.) WRITE (6,66) I
CALL FUNCT (XS, YS, FA(I,J), NAF, DGN, NTWO)
IF (MODE.EQ.1) CALL FUNCT (XS, YS, F, NOF, DGN, MONE)
THEO(I,J)=F
IF (NNN.GE.2) REWIND 3
IF (NNN.GE.2) CALL FIELD (RS, SIGI, SOLN, NBD, BDA, DGN, KBD)
IF (NNN.EQ.1) CALL FIELD2 (RS, SIGI, SOLN, G, H, SCF, DGN)
CA(I)=SOLN/BOX/HS
CALC(I,J)=CA(I)-FA(I,J)

```

9

44

57

CAL01660
 CAL01670
 CAL01680
 CAL01690
 CAL01700
 CAL01710
 CAL01720
 CAL01730
 CAL01740
 CAL01750
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 CAL01770
 CAL01780
 CAL01790
 CAL01800
 CAL01810
 CAL01820
 CAL01830
 CAL01840
 CAL01850
 CAL01860
 CAL01870
 CAL01880
 CAL01890
 CAL01900
 CAL01910
 CAL01920
 CAL01930
 CAL01940
 CAL01950
 CAL01960
 CAL01970
 CAL01980
 CAL01990
 CAL02000
 CAL02010
 CAL02020
 CAL02030
 CAL02040
 CAL02050
 CAL02060
 CAL02070
 CAL02080
 CAL02090
 CAL02100
 CAL02110
 CAL02120
 CAL02130

CAL02140
CAL02150
CAL02160
CAL02170
CAL02180
CAL02190
CAL02200
CAL02210

CAL02260
CAL02270
CAL02280
CAL02290
CAL02300
CAL02310
CAL02320
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CAL02350
CAL02360
CAL02370
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CAL02400
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CAL02470
CAL02480
CAL02490
CAL02500
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CAL02520
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CAL02540
CAL02550
CAL02560
CAL02570
CAL02580
CAL02590
CAL02600
CAL02610

```

13      RHO(I)=RHOINF*(CALC(I,J)+1.)
      ERR(I)=CA(I)
      IF (MODE.EQ.1) ERR(I)=(CALC(I,J)-THEO(I,J))
      IF (MODE.GT.1) THEO(I,J)=FA(I,J)
      IF (LPT.EQ.0) GO TO 3
      LC=0
      TL(I)=BL
      TTL=0
      IF ((XP(I).GT.XPM).AND.(XP(I).LT.XPR)) TTL=1.
      IF (IC.EQ.5) IC=0
      IF (IC.EQ.0) TL(I)=PL
      DO 2 L=2,62
      TL(L)=BL
      IF ((I.EQ.1).OR.(TTL.EQ.1).OR.(I.EQ.NPTS)) TL(L)=PL
      IF (LC.EQ.13) LC=0
      IF ((IC.EQ.0).AND.(LC.EQ.0)) TL(L)=PL
      LC=LC+1
      TL(2)=PL
      TL(22)=PL
      TL(62)=PL
      IC=IC+1
      RLW=(CA(I)+1.)*20.+2.5
      LW=RLW
      IF (LW.GT.62) LW=62
      IF (LW.LT.2) LW=2
      TL(LW)=SC
      RLY=(FA(I,J)+1.)*20.+2.5
      LY=RLY
      IF (LY.GT.62) LY=62
      IF (LY.LT.2) LY=2
      IF (NAF.NE.0) TL(LY)=ST
      RLX=(THEO(I,J)+1.)*20.+2.5
      LX=RLX
      IF (LX.GT.62) LX=62
      IF (LX.LT.2) LX=2
      IF (MODE.EQ.1) TL(LX)=OH
      RLZ=(CALC(I,J)+1.)*20.+2.5
      LZ=RLZ
      IF (LZ.GT.62) LZ=62
      IF (LZ.LT.2) LZ=2
      TL(LZ)=EX
      WRITE (6,82) MOUT,KOUT,INDEX,THEO(I,J),ERR(I),CALC(I,J),RHO(I),
1XP(I),(TL(L),L=1,62)
      IF ((NPTS.LE.20).AND.(I.NE.NPTS)) WRITE (6,79)
      CONTINUE
      IF (LPT.NE.0) WRITE (6,81) (DH,I=1,54),(PL,I=1,13)
      TMAX=0.
      TMIN=0.
3

```


CAL02620
CAL02630
CAL02640
CAL02650
CAL02660
CAL02670
CAL02680
CAL02690
CAL02700
CAL02710
CAL02720
CAL02730
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CAL02750
CAL02760
CAL02770
CAL02780
CAL02790
CAL02800
CAL02810
CAL02820
CAL02830
CAL02840
CAL02850
CAL02860
CAL02870
CAL02880
CAL02890
CAL02900
CAL02910
CAL02920
CAL02930
CAL02940
CAL02950
CAL02960
CAL02970
CAL02980
CAL02990
CAL03000
CAL03010
CAL03020
CAL03030
CAL03040
CAL03050
CAL03060
CAL03070
CAL03080
CAL03090

```

IE=0
BE=0.
EB=0.
DO 4 I=1,NPTS
  TH=THEO(I,J)      TMAX=TH
  IF (TH.GT.TMAX)    TMIN=TH
  IF (TH.LT.TMIN)
  ER=ABS(CALC(I,J)-TH)
  IF (ER.LE.BE) GO TO 4
  BE=ER
  IE=I
  CONTINUE
  TMM=TMAX-TMIN
  EB=RHOINF*(CALC(IE,J)-THEO(IE,J))
  IF (TMM.NE.0.) BE=(CALC(IE,J)-THEO(IE,J))*100./TMM
  IF ((MUDE.EQ.1).AND.(LPT.NE.0)) WRITE (6,75) EB,XP(IE),BE
  IF ((DELPHI.NE.0.) YP(J)=PHI
  CONTINUE
  IF (BND.EQ.0.) GO TO 14
  IF (LPT.EQ.1) WRITE (6,78) (ST,I=1,124)
  IF (LPT.GT.1) WRITE (6,74) (ST,I=1,95)
  IF ((CMS.EQ.1).AND.(LPT.GT.1)) READ (5,79) ZZ
  IF (DGN.GE.1.) WRITE (6,63)
  CALL MAP (NPTS,NLINS,CALC,NOF,Z,BND)
  IF (NAF.EQ.0) GO TO 10
  NAO=10*NOF+NAF
  IF ((DGN.GE.1.).AND.(NGP.EQ.-3)) WRITE (6,62)
  IF (NGP.EQ.-3) CALL GPUNCH (Z,XO,YO,PHISYM,NAO,IMAX,JMAX,G)
  DO 7 IJ=1,IJMX
    G(IJ)=G(IJ)-GA(IJ)
    IF (IPT.LE.0) GO TO 11
    IF ((IPT.EQ.1).OR.(IPT.EQ.3)) WRITE (6,78) (ST,I=1,124)
    IF ((IPT.EQ.2).OR.(IPT.GE.4)) WRITE (6,74) (ST,I=1,95)
    IF ((CMS.EQ.1).AND.((IPT.EQ.2).OR.(IPT.GE.4))) READ (5,79) ZZ
    CALL GPRINT (G,MONE)
    IF (NGP.EQ.-1) CALL GPUNCH (Z,XO,YO,PHISYM,NOF,IMAX,JMAX,G)
    IF (IPT.EQ.3) WRITE (6,78) (ST,I=1,124)
    IF (IPT.GE.4) WRITE (6,74) (ST,I=1,95)
    IF ((CMS.EQ.1).AND.(IPT.GE.4)) READ (5,79) ZZ
    IF (IPT.GE.3) CALL GPRINT (GA,NTWO)
    IF (KPT.LE.0) GO TO 12
    IF ((KPT.EQ.1).OR.(KPT.EQ.3)) WRITE (6,78) (ST,I=1,124)
    IF ((KPT.EQ.2).OR.(KPT.GE.4)) WRITE (6,74) (ST,I=1,95)
    IF ((CMS.EQ.1).AND.((KPT.EQ.2).OR.(KPT.GE.4))) READ (5,79) ZZ
    IF (DGN.GE.1.) WRITE (6,61)
    CALL GPLOT (G,GA,JMS)
    WRITE (6,73) (EX,I=1,124)
    AGAIN=ST
  7 CONTINUE
  10 CONTINUE
  11 CONTINUE
  12 CONTINUE

```


89	IF (CMS.NE.1.) READ(5,60) AGAIN	CAL031100
88	IF (AGAIN.EQ.BL) GO TO 20	CAL031110
	WRITE (6,77)	CAL031120
	FORMAT (6F12.7)	CAL031130
87	1, POINTS) WAS: /7(1F10.3/)	CAL031140
	THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,	CAL031150
86	1, POINTS) WAS: /7(1F10.3/)	CAL031160
85	1, POINTS) WAS: /7(1F10.3/)	CAL031170
	THE INVERTED CROSS SECTION FOR:)	CAL031180
	FORMAT (1H1//, THE INVERTED CROSS SECTION FOR:)	CAL031190
	FORMAT (10X, Z =, F8.3, CM, /10X, PHI=, F8.3, DEGREES, /10X,	CAL031200
	14HY, =, F8.3, CM, .44X, 0 = ORIGINAL FUNCTION:/	CAL031210
84	270X, .* = ADD-ON FUNCTION:)	CAL031220
83	FORMAT (, ADJUST PAGE, HIT SPACE AND RETURN.')	CAL031230
	FORMAT (, LIMIT MAX ORIGINAL ABS. COMPUTED (MG/CC)'/	CAL031240
	1 K TERM FUNCTION ERROR FUNCTION DENSITY', 6H X', F4.1,	CAL031250
82	26F10.1) (2X, I2, 1X, I3, 1X, I4, 1X, F9.4, 2X, F7.3, 1X, F9.4, 1X, F7.3,	CAL031260
	FORMAT (X, 62A1)	CAL031270
81	1 F7.3, 1X, 2X, A1, 12(4X, A1))	CAL031280
80	FORMAT (3X, 54A1, 2X, A1, 12(4X, A1))	CAL031290
	FORMAT (, LIMIT MAX, ADD-ON, 6X, THE, 4X, DENSITY (MG/CC)'/	CAL031300
	1, K TERM, 2X, FUNCTION, 5X, SUM, 3X, FUNCTION DENSITY',	CAL031310
79	2, H X', F4.1, 6F10.1)	CAL031320
78	FORMAT (1X, F10.3)	CAL031330
77	FORMAT (1X, 124A1)	CAL031340
76	FORMAT (1X//)	CAL031350
	1, INVERTED SUM, /70X, X = COMPUTED FUNCTION')	CAL031360
75	1, FORMAT (, LARGEST ERROR: ', F8.6, GMS/CC; AT ', 3HX' =, F6.3	CAL031370
	1, FORMAT (, F10.2, PERCENT, //)	CAL031380
74	1, FORMAT (1X, 47A1, SET PAGE, HIT SPACE, RETURN ', 48A1)	CAL031390
69	FORMAT (, CALL FREAD')	CAL031400
68	FORMAT (, CALL GARRAY')	CAL031410
67	FORMAT (, LINE, I3, DO LOOP')	CAL031420
66	FORMAT (, POINT, I3, CALL FUNCT')	CAL031430
65	FORMAT (, CALL FIELD')	CAL031440
64	FORMAT (, CALL BDGEN')	CAL031450
63	FORMAT (, CALL MAP')	CAL031460
62	FORMAT (, CALL GPUNCH')	CAL031470
61	FORMAT (, CALL GPLOT')	CAL031480
60	FORMAT (80A1)	CAL031490
	STOP	CAL031500
	END	CAL031510
		CAL031520
		SUB000010
		SUB000020

C000001
C


```

C SUBROUTINE BDGEN (G,H,SCF,DGN,NBD,BDA,KBD)
C BDGEN EVALUATES THE B AND D COEFFICIENTS FOR ALL M AND K, AND WRITES
C THE ARRAY LINEARLY ON DISK.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /TAB/ INDEX,KEXTRA,MEXTRA,KLIMIT,MLIMIT,KOUT,MOUT
COMMON /SYM/ ISYM,MSYM,FCU,IMS,JMS,QSYM
DIMENSION G(IJMX),H(IIMX,5),SCF(JJMX,6),BDA(KBD)
C INITIALIZE THE VALUES:
INDEX=0
KL2=NBD*KLIMIT
REWIND 3
JJMX6=JJMX*6
IIMX2=(IIMX+1)/2
PIE=3.141592653589793
RIMAX=IMAX
KLMP=KLIMIT+1
DX=2./RIMAX
RJMAX=JMAX
DXI=2.*PIE/FCU
C INITIALIZE THE MODIFIED HERMITE POLYNOMIAL ARRAY: VECTORS:
C (1)=H1, (2)=H2, (3)=ALPHA*X(1), (4)=HM+2 STORED, (5)=HM+1 STORED
DO 1, II=1, IIMX2
RII=II
IIM=IIMX-II+1
H(II,3)=ALPHA*(RII*DX-DX-1.)
H(IIM,3)=-H(II,3)
H(II,1)=2.*H(II,3)
H(II,2)=(H(II,3)*H(II,1)-1.)/3.
H(IIM,1)=-H(II,1)
H(IIM,2)=H(II,2)
H(II,5)=H(IIM,2)
H(IIM,5)=H(II,5)
H(II,4)=H(IIM,1)
H(IIM,4)=H(II,4)
SIGN=1.
C INITIALIZE THE SIN/COS ARRAY:
DO 2 J=1,JJMX
RJM=J-1
SCF(J,1)=0.
SCF(J,2)=1.
SCF(J,3)=SIN(RJM*DXI-PIE/2.)
SCF(J,4)=COS(RJM*DXI-PIE/2.)
SCF(J,5)=0.
SCF(J,6)=0.
MS=0
C COMMENCE THE M LOOP:
SUB000030
SUB000040
SUB000050
SUB000060
SUB000070
SUB000080
SUB000090
SUB000100
SUB000110
SUB000120
SUB000130
SUB000140
SUB000150
SUB000160
SUB000170
SUB000180
SUB000190
SUB000200
SUB000210
SUB000220
SUB000230
SUB000240
SUB000250
SUB000260
SUB000270
SUB000280
SUB000290
SUB000300
SUB000310
SUB000320
SUB000330
SUB000340
SUB000350
SUB000360
SUB000370
SUB000380
SUB000390
SUB000400
SUB000410
SUB000420
SUB000430
SUB000440
SUB000450
SUB000460
SUB000470
SUB000480
SUB000490
SUB000500

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DO 7 MP=1,MLIMIT
RM=MP-1
RM=M
SIGN=-SIGN
IF (DGN.LE.-4) WRITE (6,88) SCF(1,1),SCF(2,1),SCF(1,2),SCF(2,2)
C TEST FOR SYMMETRY SKIPS:
IF (MS.EQ.MSYM) MS=0
TOTAL=0.
MS=MS+1
IF (MS.NE.1) GO TO 6
C COMMENCE THE K LOOP:
DO 5 KP=1,KLIMIT
K=KP-1
PK=KP
RK=K
INDEX=INDEX+1
C CALL THE B & D COEFFICIENTS AND WRITE THEM ON DISK:
CALL BD (M,K,G,H,SCF,B,D,JJMX6)
IF (DGN.EQ.3.) WRITE (6,89) M,K,B,D
IF (DGN.LE.-2) WRITE (6,89) M,K,B,D
IF (DGN.LE.-4) WRITE (6,88) H(1,1),H(1,2),H(1,4),H(1,5)
KK=K*NBD+1
K2=KP*NBD
BDA(K2)=D
BDA(KK)=B
C GENERATE THE NEXT ORDER OF THE SET OF HERMITE POLYNOMIALS FOR NEW K:
ORDER=M+2*KP+1
HA=SQRT((PK+RM))/(ORDER+1.)/((ORDER+2.))
HB=2.*SQRT((PK+1.))*(RM+PK+1.)/((ORDER+1.)/((ORDER+2.))
DO 5 II=1,IIMX2
IIM=IIMX-II+1
H(II,1)=2.*(H(II,3)*H(II,2)-HA*H(II,1))
H(II,1,1)=SIGN*H(II,1)
H(II,2)=HB*(H(II,3)*H(II,1)-ORDER*H(II,2))
H(II,2,1)=HB*(H(II,3)*H(II,1)-ORDER*H(II,2))
C ADVANCE THE SIN/COS ARRAY FOR THE NEXT M:
DO 3 J=1,JJMX
IF (DGN.LE.-5) WRITE (6,87) (SCF(J,NT),NT=1,6)
FORMAT(1, SIN/COS MXI:8E10.3)
TEMP=SCF(J,1)
SCF(J,1)=SCF(J,4)+SCF(J,2)*SCF(J,3)
SCF(J,2)=SCF(J,1)*SCF(J,4)-STEMP*SCF(J,3)
DO 4 J=1,JMAX
SCF(J,5)=SCF(J+1,1)-SCF(J,1)
SCF(J,6)=SCF(J+1,2)-SCF(J,2)
C WRITE (3) (BDA(I),I=1,KBD)
IF (DGN.LE.-3) WRITE (6,88) (BDA(I),I=1,10)
IF (JSYM.GT.JMAX) RETURN
RM=RM+1

```

SUB00510
SUB00520
SUB00530
SUB00540
SUB00550
SUB00560
SUB00570
SUB00580
SUB00590
SUB00600
SUB00610
SUB00620
SUB00630
SUB00640
SUB00650
SUB00660
SUB00670
SUB00680

SUB00690
SUB00700
SUB00710
SUB00720
SUB00730
SUB00740
SUB00750
SUB00760
SUB00770
SUB00780
SUB00790
SUB00800
SUB00810
SUB00820
SUB00830
SUB00840
SUB00850
SUB00860
SUB00870
SUB00880
SUB00890
SUB00900
SUB00910
SUB00920
SUB00930
SUB00940
SUB00950
SUB00960
SUB00970


```

C REGENERATE THE HERMITE ARRAY FOR NEW M, K=0:
DO 7 II=1,IIMX2
  IIM=IIMX-II+1
  H(II,2)=H(II,4)*SQRT(RM)/(RM+1.)
  H(II,1)=H(II,5)*{(RM+2.)
  H(IIM,1)=-SIGN*H(II,1)
  H(II,2)=2.*SQRT(RM+1)*(H(II,3)*H(II,1)-(RM+1.)*H(II,2))
  H(II,2)=H(II,2)/(RM+3.)
  H(II,4)=H(II,1)
  H(II,5)=H(II,2)
  FORMAT (' M=',I4,' B=',E10.4,' D=',E10.4)
  FORMAT (2X,10E10.3)
  RETURN
END
C000002
C

```

```

SUBROUTINE FIELD (RS,SIG,SOLN,NBD,BDA,DGN,KBD)
C
C FIELD EVALUATES THE VALUE OF THE FIELD FUNCTION AT A PARTICULAR
C POINT DESIGNATED IN CYLINDRICAL COORDINATES, BY USING THE INVERSION
C EQUATION OF MALDONADO. ET.AL. FIELD USES THE ARRAY OF B & D
C COEFFICIENTS GENERATED IN SUBROUTINE BDGEN.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /TAB/ INDEX,KEXTRA,MEXTRA,KLIMIT,MLIMIT,KOUT,MOUT
COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
DIMENSION BDA(KBD),STK(52),STM(52)
C INITIALIZE THE VALUES:
INDEX=0
MTIMER=0
KOUT=0
MOUT=0
MMAX=0
KMAX=0
TOTAL=0.
JJMX6=JJMX*6
REWIND 3
AR=ALPHA*RS
ARG=AR*#2
EXPON=EXP(-ARG)
PIE=3.141592653589793
APP=ALPHA/PIE/PIE
M=0
RM=M
RIMAX=IMAX
DX=2./RIMAX
C
SUB01140
SUB01150
SUB01160
SUB01170
SUB01180
SUB01190
SUB01200
SUB01210
SUB01220
SUB01230
SUB01240
SUB01250
SUB01260
SUB01270
SUB01280
SUB01290
SUB01300
SUB01310
SUB01320
SUB01330
SUB01340
SUB01350
SUB01360
SUB01370
SUB01380
SUB01390
SUB01400
SUB01410
SUB01420
SUB01430

```


SUB01440
SUB01450
SUB01460
SUB01470
SUB01480
SUB01490
SUB01500
SUB01510
SUB01520
SUB01530
SUB01540
SUB01550
SUB01560
SUB01570
SUB01580
SUB01590
SUB01600
SUB01610
SUB01620
SUB01630
SUB01640
SUB01650
SUB01660
SUB01670
SUB01680
SUB01690
SUB01700
SUB01710
SUB01720
SUB01730
SUB01740
SUB01750
SUB01760
SUB01770
SUB01780
SUB01790
SUB01800
SUB01810
SUB01820
SUB01830
SUB01840
SUB01850
SUB01860
SUB01870
SUB01880
SUB01890
SUB01900
SUB01910

```

RJMAX=JMAX
SIGN=1.
STK(1)=0.
STM(1)=0.
SMS=0.
CMI=1.
CMI=SIGN(SIG)
CMI=COS(SIG)
MEP=ME XTRA+1
DO 16 MB=1,MEP
  STM(MB)=0.
  FM=1.
  MS=0
  C COMMENCE THE M LOOP:
  SIGN=-SIGN
  K=0
  RK=K
  RM=M
  ARM=1.
  IF (M.NE.0) ARM=AR**M
  KTIMER=0
  KEP=KEXTRA+1
  DO 15 KB=1,KEP
    STK(KB)=0.
    SIGNK=-1.
  C COMPUTE THE K=0 & K=1 ORDERS OF LAGUERRE POLYNOMIAL FOR GIVEN M:
  PM=0.
  P=SQRT(1./FM)
  PP=(RM+1.-ARG)*SQRT(1./FM/(RM+1.))
  C TEST FOR SYMMETRY SKIPS:
  IF (MS.EQ.MSYM) MS=0
  MS=MS+1
  IF (MS.NE.1) GO TO 7
  C READ A LINE OF B & D COEFFICIENTS FOR GIVEN M:
  READ (3) (BDA(I),I=1,KBD)
  IF (DGN.LE.-6) WRITE (6,88) (BDA(I),I=1,10)
  C COMMENCE THE K LOOP:
  INDEX=INDEX+1
  SIGNK=-SIGNK
  C COMPUTE THE M,K SUMMATION TERM:
  KK=K+NBD+1
  B=BDA(KK)
  D=0.
  IF (NBD.EQ.2) D=BDA(KK+1)
  BRAKET=B
  IF (RM.EQ.0.) GO TO 4
  BRAKET=B*CMS+D*SMS
  ADD=SIGNK*BRAKET*P*ARM

```

16

2

15

3

4


```

TOTAL=TOTAL+ADD
IF (DGN.GT.-5) GO TO 5
STOT=TOTAL*EXPON*APP/BOX/SIZE
WRITE (6,89) M,K,STOT,ADD,BRACKET,P,ARM,B,CMS,D,SMS
ESTABLISH CHECK AS THE RELATIVE SIZE OF THE M,K TERM OF THE SERIES:
5 CHECK=ABS(ADD)
IF (TOTAL.GT.EPS) CHECK=ABS(ADD/TOTAL)
C ADVANCE THE K INDEX:
K=K+1
RK=K
DO 10 KA=1,KEXTRA
KB=KEXTRA-KA+1
STK(KB+1)=STK(KB)
STK(2)=TOTAL
ORDER=M+2*K+1
C GENERATE THE NEXT ORDER OF LAGUERRE POLYNOMIAL FOR NEW K:
PM=PP
P=PP
PP=P*(ORDER-ARG)-PM*SQRT(RK*(RM+RK))
PP=PP/SQRT((RK+1.)*(RM+RK+1.))
C SET K TIMER TO PROVIDE EXTRA K TERMS AFTER CHECK < EPS:
KTIMER=KTIMER+1
IF (K.GE.KLIMIT) GO TO 6
IF (CHECK.GE.EPS) KTIMER=0
IF (KTIMER.LE.KEXTRA) GO TO 3
GO TO 7
6 KOUT=KOUT+1
IF (KEXTRA.EQ.0) GO TO 7
TOTAL=0
DO 11 KA=1,KEXTRA
TOTAL=TOTAL+STK(KA+1)
RKX=KEXTRA
TOTAL=TOTAL/RKX
C END OF K LOOP: ADVANCE M:
M=M+1
RM=M
STP=SMS
SMS=SMS*CMI+CMS*SMI
CMS=CMS*CMI-STP*SMI
IF (K.GT.KMAX) KMAX=K
FM=FM*RM
DO 12 MA=1,MEXTRA
MB=MEXTRA-MA+1
STM(MB+1)=STM(MB)
STM(2)=TOTAL
C SET M TIMER FOR EXTRA M TERMS:
IF (MS.EQ.1) MTIMER=MTIMER+1
IF (JSYM.GT.JMAX) GO TO 9

```

SUB01920
SUB01930
SUB01940
SUB01950
SUB01960
SUB01970
SUB01980
SUB01990
SUB02000
SUB02010
SUB02020
SUB02030
SUB02040
SUB02050
SUB02060
SUB02070
SUB02080
SUB02090
SUB02100
SUB02110
SUB02120
SUB02130
SUB02140
SUB02150
SUB02160
SUB02170
SUB02180
SUB02190
SUB02200
SUB02210
SUB02220
SUB02230
SUB02240
SUB02250
SUB02260
SUB02270
SUB02280
SUB02290
SUB02300
SUB02310
SUB02320
SUB02330
SUB02340
SUB02350
SUB02360
SUB02370
SUB02380
SUB02390


```

13 IF (K.GT.KEXTRA) MTIMER=0
   IF (M.GE.MLIMIT) GO TO 13
   IF (MTIMER.LE.MEXTRA) GO TO 2
   IF (MEXTRA.EQ.0) GO TO 9
   TOTAL=0
   DO 14 MA=1,MEXTRA
   TOTAL=TOTAL+SIM(MA+1)
14 RMX=MEXTRA
   TOTAL=TOTAL/RMX
   C   END OF M LOOP; COMPUTE OUTPUT SOLN.
9   MOUT=M-1
   IF (KOUT.EQ.0) KOUT=KMAX-1
   SOLN=TOTAL*EXPON*APP/2.
89 FORMAT (1,M=1,I4,'',K=1,I4,'', SUBTOTAL='0.9E10.3)
88 FORMAT (2X,10E10.3)
   RETURN
   END
C000003
C

```

```

SUB02400
SUB02410
SUB02420
SUB02430
SUB02440
SUB02450
SUB02460
SUB02470
SUB02480
SUB02490
SUB02500
SUB02510
SUB02520
SUB02530
SUB02540
SUB02550
SUB02560
SUB02570
SUB02580

```

```

C   SUBROUTINE BD (M,K,G,H,SCF,B,D,JJMX6)
C
C   BD EVALUATES THE FIRST (B) AND SECOND (D) COEFFICIENTS IN THE
C   INVERSION EQUATION, FOR A PARTICULAR SET OF INDEXES M & K.
C   BD MAKES USE OF THE HERMITE POLYNOMIAL ARRAY GENERATED BY
C   SUBROUTINE FIELD AS M & K ADVANCE.
C

```

```

SUB02590
SUB02600
SUB02610
SUB02620
SUB02630
SUB02640
SUB02650
SUB02660
SUB02670
SUB02680
SUB02690
SUB02700
SUB02710
SUB02720
SUB02730
SUB02740
SUB02750
SUB02760
SUB02770

```

```

C   COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
C   COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
C   DIMENSION G(IJMX),SCF(JJMX6),H(IIMX)
C   PIE=3.141592653589793
C   B=0.
C   D=0.
C   RM=M
C   RK=K
C   RJMAX=JMAX
C   JJMX4=4*JJMX
C   DXI=2.*PIE/FCU
C   FCRMAT(1X,I10,/)
200 IF (JSYM.LE.0) GO TO 4
   IF (M.NE.0) GO TO 2
   S=DXI
   DO 1 J=1,JMAX
   DO 1 I=1,IMAX
   II=I+1
   IJ=IMAX*(J-1)+I

```

```

SUB02780
SUB02790
SUB02800
SUB02810
SUB02820
SUB02830
SUB02840

```



```

1      DH=H(I,I)-H(I)
      B=B+G(I,I)*S*DH
      B=B*QSYM/2.
      RETURN
2      DO 3 J=1,JMAX
      JS=J+JJMX4
      S=SCF(JS)/RM
      DO 3 I=1,I MAX
      II=I+1
      IJ=I MAX*(J-1)+I
      DH=H(I,I)-H(I)
      B=B+G(I,I)*S*DH
      B=B*QSYM
      RETURN
4      IF (M.NE.O) GO TO 6
      S=DXI
      DO 5 J=1,JMAX
      DO 5 I=1,I MAX
      II=I+1
      IJ=I MAX*(J-1)+I
      DH=H(I,I)-H(I)
      B=B+G(I,I)*S*DH
      B=B/2.
      RETURN
5      DO 7 J=1,JMAX
      JS=J+JJMX4
      J2=JS+JJMX
      S=SCF(JS)/RM
      C=SCF(J2)/RM
      DO 7 I=1,I MAX
      II=I+1
      IJ=I MAX*(J-1)+I
      DH=H(I,I)-H(I)
      B=B+G(I,I)*S*DH
      FORMAT (' BD:',4I5,10F6.2)
      D=D-G(I,I)*C*DH
      RETURN
      END
C000004
C

```

SUBROUTINE FIELD2 (RS,SIG,SOLN,G,H,SCF,DGN)

FIELD2 COMPUTES THE SAME INVERSION AS SUBROUTINE FIELD, EXCEPT THAT THE COEFFICIENTS B AND D ARE COMPUTED INDIVIDUALLY AS USED BY CALLING BD. DISK STORAGE IS NOT REQUIRED, BUT COMPUTING TIME IS MUCH GREATER. FIELD2 IS UTILIZED BY SPECIFYING A NEGATIVE MODE ON

SUB02850
SUB02860
SUB02870
SUB02880
SUB02890
SUB02900
SUB02910
SUB02920
SUB02930
SUB02940
SUB02950
SUB02960
SUB02970
SUB02980
SUB02990
SUB03000
SUB03010
SUB03020
SUB03030
SUB03040
SUB03050
SUB03060
SUB03070
SUB03080
SUB03090
SUB03100
SUB03110
SUB03120
SUB03130
SUB03140

SUB03160
SUB03170
SUB03180
SUB03190
SUB03200
SUB03210
SUB03220
SUB03230
SUB03240

SUB03250
SUB03260
SUB03270
SUB03280
SUB03290
SUB03300

SUB03790
SUB03800
SUB03810
SUB03820
SUB03830
SUB03840
SUB03850
SUB03860
SUB03870
SUB03880
SUB03890
SUB03900
SUB03910
SUB03920
SUB03930
SUB03940
SUB03950
SUB03960
SUB03970
SUB03980
SUB03990
SUB04000
SUB04010
SUB04020
SUB04030
SUB04040
SUB04050
SUB04060
SUB04070
SUB04080
SUB04090
SUB04100
SUB04110
SUB04120
SUB04130
SUB04140
SUB04150
SUB04160
SUB04170
SUB04180
SUB04190
SUB04200
SUB04210
SUB04220
SUB04230
SUB04240
SUB04250

```

11 DO 11 J=1,JMX
    RJM=J-1
    SCF(J,1)=0.
    SCF(J,2)=1.
    SCF(J,3)=SIN(RJM*DXI-PIE)
    SCF(J,4)=COS(RJM*DXI-PIE)
    SCF(J,5)=0.
    SCF(J,6)=0.
    MS=0
    C COMMENCE THE M LOOP:
    2 SIGN=-SIGN
    K=0
    RK=K
    ARM=1.
    IF (M.NE.0) ARM=AR**M
    KTIMER=0
    SIGNK=-1.
    C COMPUTE THE K=0 & K=1 ORDERS OF LAGUERRE POLYNOMIAL FOR GIVEN M:
    PM=0.
    P=SQRT(1./FM)
    PP=(RM+1.-ARG)*SQRT(1./FM/(RM+1.))
    C ADVANCE THE SIN/COS ARRAY FOR NEW M:
    DO 12 J=1,JMX
    12 SCF(J,1)=SCF(J,1)*SCF(J,4)+SCF(J,2)*SCF(J,3)
    SCF(J,2)=SCF(J,2)*SCF(J,4)-SCF(J,1)*SCF(J,3)
    DO 13 J=1,JMAX
    13 SCF(J,5)=SCF(J+1,1)-SCF(J,1)
    SCF(J,6)=SCF(J+1,2)-SCF(J,2)
    C TEST FOR SYMMETRY SKIPS:
    IF (MS.EQ.MSYM) MS=0
    TOTAL=0.
    MS=MS+1
    IF (MS.NE.1) GO TO 7
    RMS=RM*SIG
    CMS=COS(RMS)
    SMS=SIN(RMS)
    C COMMENCE THE K LOOP:
    INDEX=INDEX+1
    SIGNK=-SIGNK
    C CALL THE B & D COEFFICIENTS AND COMPUTE THE M,K SUMMATION TERM:
    CALL BD (M,K,G,H,SCF,B,D,JJMX6)
    IF (DGN.LE.-2.) WRITE (6,89) M,K,B,D
    BRACKET=B
    IF (RM.EQ.0.) GO TO 4
    BRACKET=B*CMS+D*SMS
    ADD=SIGNK*BRACKET*P*ARM
    TOTAL=TOTAL+ADD
    4 C ESTABLISH CHECK AS THE RELATIVE SIZE OF THE M,K TERM OF THE SERIES:
    C

```



```

CHECK=ABS(ADD)
IF (TOTAL.GT.EPS) CHECK=ABS(ADD/TOTAL)
C ADVANCE THE K INDEX:
K=K+1
RK=K
ORDER=M+2*K+1
C GENERATE THE NEXT ORDER OF LAGUERRE POLYNOMIAL FOR NEW K:
PM=PP
P=PP
PP=P*(ORDER-ARG)-PM*SQRT(RK*(RM+RK))
PP=PP/SQRT((RK+1.)*(RM+RK+1.))
C GENERATE THE NEXT ORDER OF THE SET OF HERMITE POLYNOMIALS FOR NEW K:
HA=SQRT((RK+RM))/ORDER
HB=2.*SQRT((RK+1.)*(RM+RK+1.))/((ORDER+1.))/((ORDER+2.))
DO 5 I=1,IIMX2
IIM=IIMX-I+1
H(I,1)=2.*(H(I,3)*H(I,2)-HA*H(I,1))
H(I,1)=SIGN*H(I,1)
H(I,2)=HB*(H(I,3)*H(I,1)-ORDER*H(I,2))
5 C SET K TIMER TO PROVIDE EXTRA K TERMS AFTER CHECK < EPS:
KTIMER=KTIMER+1
IF (K.GE.KLIMIT) GO TO 6
IF (CHECK.GE.EPS) KTIMER=0
IF (KTIMER.LE.KEXTRA) GO TO 3
GO TO 7
C END OF K LOOP: ADVANCE M AND COMPUTE NEW TOTAL:
KOUT=KOUT+1
M=M+1
IF (K.GT.KMAX) KMAX=K
RM=M
FM=FM*RM
C REGENERATE THE HERMITE ARRAY FOR NEW M, K=0:
DO 8 I=1,IIMX2
IIM=IIMX-I+1
H(I,2)=H(I,4)*SQRT(RM)/(RM+1.)
H(I,1)=H(I,5)*(RM+2.)
H(I,1)=-SIGN*H(I,1)
H(I,2)=2.*SQRT(RM+1.)*(H(I,3)*H(I,1)-(RM+1.)*H(I,2))
H(I,2)=H(I,2)/(RM+2.)/(RM+3.)
H(I,4)=H(I,1)
H(I,5)=H(I,2)
C SET M TIMER FOR EXTRA M TERMS:
IF (MS.EQ.1) MTIMER=MTIMER+1
IF ((JSYM.GT.JMAX) GO TO 9
IF (K.GT.KEXTRA) MTIMER=0
IF (M.GE.MLIMIT) GO TO 9
IF (MTIMER.LE.MEXTRA) GO TO 2 SOLN.
C END OF M LOOP: COMPUTE OUTPUT

```

SUB04270
SUB04280
SUB04290
SUB04300
SUB04310
SUB04320
SUB04330
SUB04340
SUB04350
SUB04360
SUB04370
SUB04380
SUB04390
SUB04400
SUB04410
SUB04420
SUB04430
SUB04440
SUB04450
SUB04460
SUB04470
SUB04480
SUB04490
SUB04500
SUB04510
SUB04520
SUB04530
SUB04540
SUB04550
SUB04560
SUB04570
SUB04580
SUB04590
SUB04600
SUB04610
SUB04620
SUB04630
SUB04640
SUB04650
SUB04660
SUB04670
SUB04680
SUB04690
SUB04700
SUB04710
SUB04720
SUB04730
SUB04740

SUB004750
SUB004760
SUB004770
SUB004780
SUB004790
SUB004800
SUB004810
SUB004820
SUB000010
SUB000020

```

9      MOUT=M-1
      IF (KOUT.EQ.0) KOUT=KMAX-1
      SOLN=TOTAL*EXPON#APP/2.
89     FORMAT ('M=',I4,'', K=',I4,'', B=',E10.4,'', D=',E10.4)
      RETURN
      END

```

C0000005
C

SUBROUTINE GARRAY (G,GA,NDF,DGN,NUMB,XO,YO,PHISYM)

GARRAY FILLS THE DATA ARRAY OVER AN ORTHOGONAL AREA WITH
THE REGULAR DATA OBTAINED BY THE METHOD CORRESPONDING TO THE
PARTICULAR MODE:

MODE 1 - DATA OBTAINED BY SAMPLING A KNOWN FUNCTION SUPPLIED
IN SUBROUTINE FUNCT AND SAMPLED IN SUBROUTINE GOLP.

MODE 2 - DATA OBTAINED BY GENERATING A REGULAR ARRAY FROM
IRREGULAR EXPERIMENTAL INPUT DATA READ IN. CALLS
SUBROUTINE SHEET. (EXPERIMENTAL DATA MAY
BE SIMULATED, SEE 'SHEET')

MODE 3 - UTILIZES RAW DATA TAKEN AT THE PROPER INTERVAL, THE
OR PREVIOUSLY GENERATED, AND READ DIRECTLY INTO THE
GARRAY. CALLS SUBROUTINE READ.

```

COMMON IMAX,JMAX,IIMX,IJMX,ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
COMMON /IO/ CMS,IN1,IN2,IN4
DIMENSION G(IMAX,JMAX),GA(IMAX,JMAX)
PIE=3.141592653589793
HS=SIZE/2.
IF (MODE.GT.3) MODE=1
RIMX=IMAX
RJMX=JMAX
DELR=SIZE/RIMX
DELEXI=2.*PIE/FCU
IF (MODE.GT.1) GO TO 2
DO 1 J=1,JMS
  RJ=J
  XI=(RJ-.5)*DELXI-PIE
  J2=J+2*(JMS-J)
  J3=J+JMAX/2
  J4=J2+JMAX/2

```

SUB000030
SUB000040
SUB000050
SUB000060
SUB000070
SUB000080
SUB000090
SUB000100
SUB000110
SUB000120
SUB000130
SUB000140
SUB000150
SUB000160
SUB000170
SUB000180
SUB000190
SUB000200
SUB000210
SUB000220
SUB000230
SUB000240
SUB000250
SUB000260
SUB000270
SUB000280
SUB000290
SUB000300
SUB000310
SUB000320
SUB000330
SUB000340
SUB000350
SUB000360
SUB000370
SUB000380

SUB000390
SUB000400
SUB000410
SUB000420
SUB000430
SUB000440
SUB000450
SUB000460
SUB000470
SUB000480
SUB000490
SUB000500
SUB000510
SUB000520
SUB000530
SUB000540
SUB000550
SUB000560
SUB000570
SUB000580
SUB000590
SUB000600
SUB000610
SUB000620

```

DO 1 I=1,IMS
RI=I
II=I MAX+1-I
R=(RI-.5)*DELR-HS
CALL GOLF (R,XI,GIJ,NOF,DGN,NUMB)
G(I,J)=GIJ
IF (ISYM.EQ.2) G(II,J)=GIJ
IF (ISYM.EQ.2) GO TO 1
G(II,J3)=GIJ
IF (JSYM.EQ.0) GO TO 1
G(I,J2)=GIJ
G(II,J4)=GIJ
CONTINUE
GO TO 4
1 IF (MODE.GT.2) GO TO 3
2 CALL SHEET (G,GA,XO,YO,PHISYM,NOF)
GO TO 4
3 CALL READ (Z,XO,YO,PHISYM,NOF,IMAX,JMAX,G)
4 IF (DGN.GE.2) WRITE (6,39)
39 RETURN
FORMAT (' GARRAY RETURNS')
END
C000006
C

```

SUB000630
SUB000640
SUB000650
SUB000660
SUB000670
SUB000680
SUB000690
SUB000700
SUB000710
SUB000720
SUB000730
SUB000740
SUB000750
SUB000760
SUB000770
SUB000780
SUB000790
SUB000800
SUB000810
SUB000820
SUB000830
SUB000840

```

SUBROUTINE GOLF (R,XI,GIJ,NOF,DGN,NUMB)
C
C GOLF COMPUTES THE FUNCTION G(R,XI) FOR A PARTICULAR LINE OF SIGHT
C FROM A KNOWN FUNCTION CONTAINED IN SUBROUTINE FUNCT.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
ZERO=0.
LMAX=IMAX*3
RLMAX=LMAX
DELXP=SIZE/RLMAX
SXI=SIN(XI)
CXI=COS(XI)
DELYS=DELXP*CXI
DELYS=DELXP*SXI
XP=DELXP*.5-SIZE/2.
XS=XP*CXI-R*SXI
YS=XP*SXI+R*CXI
GIJ=0.
DO 1 L=1,LMAX
RL=L
CALL FUNCT(XS,YS,F,NOF,DGN,NUMB)
GIJ=GIJ+F

```



```

1      XS=XS+DELS
      YS=YS+DELS
      IF (GIJ.NE.0.) GIJ=GIJ*DELP*BOX
      IF ((SD.EQ.0.) OR (NUMB.EQ.1)) GO TO 2
      IF (DGN.GE.3) WRITE (6,28) IX
      CALL GAUSS (IX,SD,ZERO,RV)
      GIJ=GIJ+RV
2      IF (DGN.GE.3) WRITE (6,29) R,XI,GIJ
      RETURN
29     FORMAT (' R=',F8.3,' XI=',F8.3,' GIJ=',F8.3)
28     FORMAT (' GAUSS, IX=',I8)
      END
C000007
C

```

```

C      SUBROUTINE FUNCT (XS,YS,F,NOF,DGN,NUMB)
C

```

```

CP67USERID 1395BOXJ
C

```

```

C      FUNCT EVALUATES AS INPUT FUNCTION AT POSITION (X,Y) IN THE TEST
C      SECTION COORDINATE SYSTEM. NOF IDENTIFIES THE EQUATION USED.
C

```

```

COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /EQPARA/ A,B,C,D,E,P,Q,S,T,U,V,W,RO,RA,NO,NA,N1,N2
DIMENSION RO(101),RA(101)

```

```

AA=A
BB=B
CC=C
DD=D
EE=E
PP=P
IF (NUMB.LE.1) GO TO 50
AA=S
BB=T
CC=U
DD=V
EE=W
PP=0

```

```

50     PIE=3.141592653589793
      HS=SIZE/2.
      R=SQRT(XS**2+YS**2)/HS
      F=0.
      IF (R.GT.1.) GO TO 11
      IF (NOF.LE.0) GO TO 11

```

```

C      1. AXISYMMETRIC GAUSSIAN:
C

```

```

SUB00850
SUB00860
SUB00870
SUB00880
SUB00890
SUB00900
SUB00910
SUB00920
SUB00930
SUB00940
SUB00950
SUB00960
SUB00970
SUB00980

```

```

SUB00990
SUB01000

```

```

SUB01020
SUB01030
SUB01040
SUB01050
SUB01060
SUB01070
SUB01080
SUB01090
SUB01100
SUB01110
SUB01120
SUB01130
SUB01140
SUB01150
SUB01160
SUB01170
SUB01180
SUB01190
SUB01200
SUB01210
SUB01220
SUB01230
SUB01240
SUB01250
SUB01260
SUB01270
SUB01280

```



```

1      IF (NOF.GT.1) GO TO 2
C      F=AA*EXP(-1.*(R*HS/BB)**2)
C      GO TO 11
2      ADJUSTABLE RECTANGULAR STEP FUNCTION:
C      IF (NOF.GT.2) GO TO 3
C      F=PP
C      IF ((ABS(XS-DD).LE.BB).AND.(ABS(YS-EE).LE.CC)) F=AA
C      GO TO 11
3      DISPLACABLE ELLIPTICAL GAUSSIAN:
C      IF (NOF.GT.3) GO TO 4
C      F=AA*EXP(-1.*(((XS-DD)/BB)**2+((YS-EE)/CC)**2))
C      GO TO 11
4      CONSTANT:
C      IF (NOF.GT.4) GO TO 5
C      F=AA
C      GO TO 11
5      ADJUSTABLE AND DISPLACABLE ELLIPTIC RAMP FUNCTION:
C      IF (NOF.GT.5) GO TO 6
C      RBC=SQRT(((XS-DD)/BB)**2+((YS-EE)/CC)**2)
C      F=0
C      IF (RBC.LT.1.) F=AA*((1.-RBC)**PP)
C      GO TO 11
6      DISPLACABLE ELLIPTIC STEP FUNCTION:
C      IF (NOF.GT.6) GO TO 7
C      RBC=SQRT(((XS-DD)/BB)**2+((YS-EE)/CC)**2)
C      F=0
C      IF (RBC.LT.1.) F=AA
C      GO TO 11
7      CIRCULAR COSINE-SQUARED FUNCTION OF BB MAXIMA:
C      IF (NOF.GT.7) GO TO 8
C      F=AA*CCS((2.*BB-1.)*PIE/R/2.)**2
C      GO TO 11
8      NUMERICAL FUNCTION:
C      SUBROUTINE FREAD; N
C      A CONSTANT VALUE AA
C      IF (NOF.GT.8) GO TO 9
C      IF (NUMB.LE.1) N=NO
C      IF (NUMB.GT.1) N=NA
C      NM=N-1
C      RN=N

```

REQUIRES AN INPUT ARRAY READ IN BY
 FOLLOWED BY N POINT VALUES. (101 MAX)
 IS ADDED TO THE FUNCTION.


```

RI=R*(RN-1.)+1.
IR=INT(RI)
RIR=FLOAT(IR)
DI=RIR-RIR
IF (NUMB.LE.1) F=RO(IR)
IF (NUMB.GT.1) F=RA(IR)
IF ((IR.NE.N).AND.(NUMB.LE.1)) F=F+DI*(RO(IR+1)-RO(IR))
IF ((IR.NE.N).AND.(NUMB.GT.1)) F=F+DI*(RA(IR+1)-RA(IR))
F=F*AA+BB
GO TO 11

```

```

C
C
C 9. SPECIAL FUNCTION: MAY BE WRITTEN FOR THE OCCASION AND
C INSERTED IN SUBROUTINE SPFUN
C IF (NOF.GT.9) GO TO 10
C CALL SPFUN (XS,YS,F)
C GO TO 11

```

```

C EQUATIONS NO. 10 AND BEYOND ARE SET TO ZERO.
C F=0.

```

```

C
C
C 11 IF (DGN.GE.4) WRITE (6,99) XS,YS,F
C 99 FORMAT (' F8.3',, YS=' F8.3',, F=' F8.3')
C RETURN
C END
C000008
C

```

```

SUBROUTINE SPFUN (XS,YS,F)

```

```

C
C SPFUN IS A SPECIAL ROUTINE FOR EQ'N NO. 9. ANY FUNCTION MAY BE
C ENTERED.

```

```

C COMMON /EQPAR/ A,B,C,D,E,P,Q,S,T,U,V,W,RO,RA,NO,NA,N1,N2
C DIMENSION RO(101),RA(101)
C F=0.
C IF ((ABS(XS).LE.B).AND.(ABS(YS).LE.C)) F=A
C RETURN
C END
C000009
C

```

SUB01770
SUB01780
SUB01790
SUB01800
SUB01810
SUB01820
SUB01830
SUB01840
SUB01850
SUB01860
SUB01870
SUB01880
SUB01890
SUB01900
SUB01910
SUB01920
SUB01930
SUB01940
SUB01950
SUB01960
SUB01970
SUB01980
SUB01990
SUB02000
SUB02010
SUB02020
SUB02030

SUB02040
SUB02050
SUB02060
SUB02070
SUB02080
SUB02090
SUB02100
SUB02110
SUB02120
SUB02130
SUB02140
SUB02150
SUB02160


```

SUBROUTINE SHEET (G,D,XO,YO,PHISYM,NOF)
SHEET READS IRREGULARLY SPACED VALUES OF THE LINE INTEGRAL, AS
OBTAINED FROM HOLOGRAPHIC INTERFEROGRAMS. THE INTEGRAL LINES MAY BE
DEFINED EITHER BY GRID INTERCEPT POSITIONS, OR BY ANGLE AND RADIUS
ABOUT THE CENTER OF THE LABORATORY COORDINATE SYSTEM CENTER. LINES
MUST BE ENTERED IN CONSECUTIVE ORDER FROM LOWEST (NEG.) TO HIGHEST
(POS.) RADIUS. DATA MAY BE SIMULATED BY SPECIFYING NCODE=1,
FOLLOWED BY APERTURE POSITIONS FOR A FUNCTION NUMBER IN: SUBFUNCT.
SUB02170
SUB02180
SUB02190
SUB02200
SUB02210
SUB02220
SUB02230
SUB02240
SUB02250
SUB02260
SUB02270
SUB02280
SUB02290
SUB02300
SUB02310
SUB02320
SUB02330
SUB02340
SUB02350
SUB02360
SUB02370
SUB02380
SUB02390
SUB02400
SUB02410
SUB02420
SUB02430
SUB02440
SUB02450
SUB02460
SUB02470
SUB02480
SUB02490
SUB02500
SUB02510
SUB02520
SUB02530
SUB02540
SUB02550
SUB02560
SUB02570
SUB02580
SUB02590
SUB02600
SUB02610
SUB02620
SUB02630
SUB02640

COMMON IMAX,JMAX,IMX,JMX,IJMX,JJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /SYM/ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
COMMON /IO/ CMS,IN1,IN2,IN4
DIMENSION G(IMAX,JMAX),D(IMAX,JMAX),XI(303),RR(303)
DIMENSION XG(303),XD(303),YG(303),YD(303),XY(303)
NAR=303
PIE=3.141592653589793
MPIE=-PIE
TPIE=2.*PIE
MPIT=MPIE/2.
PIET=PIE/2.
ZERO THE ARRAYS:
DO 1 J=1,JMAX
DO 1 I=1,IMAX
G(I,J)=0.
D(I,J)=0.
DC 1 I=1,NAR
XG(I)=0.
XD(I)=0.
YG(I)=0.
YD(I)=0.
XY(I)=0.
XI(I)=0.
RR(I)=0.
READ THE BASIC DATA:
IF (CMS.EQ.1.) REWIND 1
READ (IN1,59) NOF,NCODE
READ (IN1,58) Z,XO,YO,PHISYM,XXM,XMN,YMX,YMN
READ (IN1,59) JM
RIMX=IMAX
DR=SIZE/RIMX
R=(-DR-SIZE)/2.*YO*2)
RZO=SQRT(XO*2+YO*2)
GAM=ATANM(YO,XO)
IP=3-ISYM
BT=JSYM
DAN=PIE*TP/BT
HS=SIZE/2.

```



```

C 3  COMMENCE THE READ AND FILL LOOP:
      DO 12 J=1,JM
      READ (IN1,59)IM
      MN=0
      XH=0.
      YH=0.
C 3  READ THE LINES, DETERMINE CODE, CALCULATE RADIUS & ANGLE FOR CODE 1:
      DO 5 I=1,IM
      IF(NCODE.LE.0)READ(IN1,58) XD(I),YD(I),XG(I),YG(I),D(I,J),RR(I),
1XI(I),XY(I)
      IF(NCODE.GE.1)CALL SIM(XD(I),YD(I),XG(I),YG(I),RR(I),XI(I),
1XY(I),XO,YO,PHISYM,XX,XMN,YMX,YMN,NOF,I,IM)
      IF(XY(I).EQ.3.) GO TO 5
      IF(XD(I).NE.0.)OR.(YD(I).NE.0.) XY(I)=1.
      IF(XG(I).NE.0.)OR.(YG(I).NE.0.) XY(I)=1.
      IF((RR(I).NE.0.)OR.(XI(I).NE.0.)).AND.(XY(I).EQ.0.)) XY(I)=2.
      IF(XY(I).EQ.0.)AND.(D(I,J).NE.0.)) XY(I)=2.
      IF(XY(I).NE.1.) GO TO 4
      DEN=SQR((XG(I)-XD(I))**2+(YG(I)-YD(I))**2)
      IF(XY(I).EQ.0.) XY(I)=4.
      IF(XY(I).EQ.4.) GO TO 4
      RR(I)=((XO-XD(I))*((YG(I)-YD(I))-(XG(I)-XD(I))*(YO-YD(I)))/DEN
      XI(I)=ATANM((YG(I)-YD(I)), (XG(I)-XD(I)))
      XIM=XI(I)
      IF(XY(I).EQ.2.) XIM=XI(I)
      XIN=XIM
      IF(XY(I).EQ.2.) RR(I)=RR(I)+RZO*SIN(GAM-XI(I))
      CONTINUE
C 5  COMPUTE MAX AND MIN ANGLE INDEXES FOR APERTURE POSITION LOCATION:
      DO 6 I=1,IM
      IF(XY(I).NE.1.)OR.(XY(I).NE.2.) GO TO 6
      IF(XI(I).GT.XIM) XIM=XI(I)
      IF(XI(I).GT.XIM) IMT=I
      IF(XI(I).LT.XIN) XIN=XI(I)
      IF(XI(I).LE.XIN) INT=I
      CONTINUE
C 6  DETERMINE APERTURE LOCATION:
      LPR=0
      XID=XI(IMT)-XI(INT)
      IF(ABS(XID).LT..00001) LPR=1
      XIH=(XI(IMT)+XI(INT))/2.

```

```

SUB02650
SUB02660
SUB02670
SUB02680
SUB02690
SUB02700
SUB02710
SUB02720
SUB02730
SUB02740
SUB02750
SUB02760
SUB02770
SUB02780
SUB02790
SUB02800
SUB02810
SUB02820
SUB02830
SUB02840
SUB02850
SUB02860
SUB02870
SUB02880
SUB02890
SUB02900
SUB02910
SUB02920
SUB02930
SUB02940
SUB02950
SUB02960
SUB02970
SUB02980
SUB02990
SUB03000
SUB03010
SUB03020
SUB03030
SUB03040
SUB03050
SUB03060
SUB03070
SUB03080
SUB03090
SUB03100
SUB03110
SUB03120

```


SUB03130
SUB03140
SUB03150
SUB03160
SUB03170
SUB03180
SUB03190
SUB03200
SUB03210
SUB03220
SUB03230
SUB03240
SUB03250
SUB03260
SUB03270
SUB03280
SUB03290
SUB03300
SUB03310
SUB03320
SUB03330
SUB03340
SUB03350
SUB03360
SUB03370
SUB03380
SUB03390
SUB03400
SUB03410
SUB03420
SUB03430
SUB03440
SUB03450
SUB03460
SUB03470
SUB03480
SUB03490
SUB03500
SUB03510
SUB03520
SUB03530
SUB03540
SUB03550
SUB03560
SUB03570
SUB03580
SUB03590
SUB03600

```

RRH=10000.
XH=RRH*COS(XIH)
YH=RRH*SIN(XIH)
IF (LPR.EQ.1) GO TO 7
YTX=-RR(IMT)*SIN(XI(IMT))-YO
XTN=-RR(IMT)*SIN(XI(IMT))-YO
XTX=RR(IMT)*COS(XI(IMT))-XO
XTN=RR(IMT)*COS(XI(IMT))-XO
UA=TAN(XI(IMT))
UC=TAN(XI(INT))
UB=YTX-UA*XTX
UC=YTN-UC*XTN
XH=(UD-UB)/(UA-UC)
YH=XH*UA+UB
RRH=SQRT((XH-XO)**2+(YH-YO)**2)
XIH=ATANM((YH-YO),(XH-XO))
CONTINUE
7 C FILL THE ANGLE AND RADIUS FOR ANY CODE 3 OR 4 LINES:
DO 9 I=1,IM
IF (XY(I).NE.3.) GO TO 8
BAS=SQRT(RRH**2-RR(I)**2)
XI(I)=XIH-ATANM(RR(I),BAS)
GO TO 9
XI(I)=ATANM((YH-YD(I)),(XH-XD(I)))
RR(I)=RRH*SIN(XI(I)-XIH)
CONTINUE
8 C ANGLES AND RADII ARE NOW FILLED FOR ALL POINTS IN THIS LINE.
C VACATE THE SET OF VECTORS TO BE USED AS TEMPORARY STORAGE:
DO 10 I=1,IM
XD(I)=0.
YD(I)=0.
XG(I)=0.
YG(I)=RR(I)
XY(I)=D(I,J)
RR(I)=0.
D(I,J)=0.
XI(I)=0.
10 C CONVERT THE LINE TO REGULAR RADII USING INTERPOLATION:
RR(I)=R+DR
CALL SPLINE(YG,XY,IM,RR(1),D(1,J))
DO 11 I=2,IMAX
RI=I
RR(I)=R+DR*RI
11 C CALL SPLINE(YG,XY,IM,RR(I),D(I,J))
C GENERATE THE VECTOR OF ANGLES FOR THIS COLUMN AND STORE IN G ARRAY:
DO 12 I=1,IMAX
BAS=SQRT(RRH**2-RR(I)**2)
G(I,J)=XIH-ATANM(RR(I),BAS)

```


SUB03610
SUB03620
SUB03630
SUB03640
SUB03650
SUB03660
SUB03670
SUB03680
SUB03690
SUB03700
SUB03710
SUB03720
SUB03730
SUB03740
SUB03750
SUB03760
SUB03770
SUB03780
SUB03790
SUB03800
SUB03810
SUB03820
SUB03830
SUB03840
SUB03850
SUB03860
SUB03870
SUB03880
SUB03890
SUB03900
SUB03910
SUB03920
SUB03930
SUB03940
SUB03950
SUB03960
SUB03970
SUB03980
SUB03990
SUB04000
SUB04010
SUB04020
SUB04030
SUB04040
SUB04050
SUB04060
SUB04070
SUB04080

```

12      YG(I)=0.XY(I)
C      D(I,J)=XY(I)
C      XY(I)=0.
C      COLUMNS ARE NOW ALL REGULARLY FILLED.
C      NEXT, INTERPOLATE EACH ROW REGULARLY OVER THE ANGLES.
C      DO 23 I=1,IMAX
C      EXPAND THE DATA TO 2 SETS TO ESTABLISH SMOOTH INTERPOLATION.
      JM3=3*JM
      II=IMAX+1-II
      IF (JSYM.NE.0) GO TO 14
      DO 13 J=1,JMS
      J2=J+JMS
      J3=J2+JMS
      XD(J)=D(I,J)
      XD(J2)=D(II,J)
      XD(J3)=D(II,J)
      XG(J)=G(I,J)-PIE-PHISYM
      XG(J2)=G(II,J)-PIE-PHISYM
      XG(J3)=G(II,J)-PHISYM
      GO TO 16
      DO 15 J=1,JMS
      J1=JMS+1-J
      J2=JMS+J
      J3=JM3+1-J
      XD(J1)=D(I,J)
      XD(J2)=D(II,J)
      XD(J3)=D(II,J)
      XG(J1)=G(I,J)-2.*(G(I,J)-PHISYM)-PIE-PHISYM
      XG(J2)=G(II,J)-PIE-PHISYM
      XG(J3)=G(II,J)+2.*(DAN+PHISYM-G(II,J))-PIE-PHISYM
      CONTINUE
      JM2=2*JMS
      JP=JMS/2
      DO 17 J=1,JM2
      XD(J)=XD(J+JP)
      XG(J)=XG(J+JP)
      JJS=JM2+1
      DO 18 J=JJS,JM3
      XD(J)=0.
      XG(J)=0.
      FIND THE SMALLEST ANGLE
      XY(1)=1.
      SA=XG(1)
      DO 19 J=1,JM2
      IF (XG(J).GE.SA) GO TO 19
      SA=XG(J)
      XY(1)=J
      CONTINUE
13
14
15
16
17
18
19

```



```

C      FIND THA MAX ANGLE IN THE ROW:
      XY(JM2)=JM2
      SB=XG(JM2)
      DO 20 J=1,JM2
      IF (XG(J).LE.SB) GO TO 20
      SB=XG(J)
      XY(JM2)=J
20 C      CONTINUE THE ORDER OF INCREASING ANGLE IN THE ROW
      DETERMINE THE ORDER OF INCREASING ANGLE IN THE ROW
      SB=XG(JM2)
      JJ=2
      JSA=XY(JJ-1)
      SA=XG(JJ-1)
      JTS=0
      DO 22 J=1,JM2
      IF (XG(J).LE.SA) GO TO 22
      IF (XG(J).GT.SB) GO TO 22
      SB=XG(J)
      XY(JJ)=J
      JTS=1
22 C      CONTINUE
      IF (JTS.EQ.0) JM2=JJ
      JJ=JJ+1
      IF (JJ.LE.JM2) GO TO 21
      DO 23 J=1,JM2
      JX=XY(J)
      YD(J)=XD(JX)
23 C      INTERPOLATE:
      DXI=2.*PIE/FCU
      XI(J)=DXI/2.-PIE-PHISYM
      CALL SPLINE (XG,YD,JM2,XI(J),G(I,J))
      DO 24 J=2,JMS
      XI(J)=XI(J-1)+DXI
      CALL SPLINN (XG,YD,JM2,XI(J),G(I,J))
      DO 25 J=1,JMS
      XIJ=XI(J)
      XU=XMX
      IF ((XIJ.GE.O.).AND.(XIJ.LT.PIE)) XU=XMN
      YU=YMN
      IF ((XIJ.GE.MPIT).AND.(XIJ.LT.PIT)) YU=YMX
      XL=XMN
      IF ((XIJ.GE.O.).AND.(XIJ.LT.PIE)) XL=XMX
      YL=YMX
      IF ((XIJ.GE.MPIT).AND.(XIJ.LT.PIT)) YL=YMN
      SXIJ=SIN(XIJ)
      CXIJ=COS(XIJ)
      RMN=(XO-XL)*SXIJ-(YO-YL)*CXIJ
      RMX=(XO-XU)*SXIJ-(YO-YU)*CXIJ

```

```

SUB04090
SUB04100
SUB04110
SUB04120
SUB04130
SUB04140
SUB04150
SUB04160
SUB04170
SUB04180
SUB04190
SUB04200
SUB04210
SUB04220
SUB04230
SUB04240
SUB04250
SUB04260
SUB04270
SUB04280
SUB04290
SUB04300
SUB04310
SUB04320
SUB04330
SUB04340
SUB04350
SUB04360
SUB04370
SUB04380
SUB04390
SUB04400
SUB04410
SUB04420
SUB04430
SUB04440
SUB04450
SUB04460
SUB04470
SUB04480
SUB04490
SUB04500
SUB04510
SUB04520
SUB04530
SUB04540
SUB04550
SUB04560

```



```

25      DO 25 I=1,IMAX
C      IF (RR(I).LT.RMN) G(I,J)=0.
C      IF (RR(I).GT.RMX) G(I,J)=0.
      CONTINUE
      EXPAND SYMMETRY SECTOR INTO AN ORTHOGONAL INTERVAL.
      IF (ISYM.EQ.2) GO TO 27
      DO 26 J=1,JMS
      J2=JMAX/2+1-J
      J3=JMAX/2+J
      J4=JMAX+1-J
      DO 26 I=1,IMAX
      II=IMAX+1-I
      G(I,J2)=G(I,J)
      G(II,J3)=G(I,J)
      G(II,J4)=G(I,J)
      RETURN
      FOR EVEN SYMMETRY, AVERAGE THE GARRY COLUMNS.
      IMS=(2*IMAX+1)/2
      DO 28 J=1,JMAX
      DO 28 I=1,IMS
      II=IMAX+1-I
      GST=(G(I,J)+G(II,J))/2.
      G(I,J)=GST
      G(II,J)=GST
      RETURN
      FORMAT (5I5)
      FORMAT (10F7.3)
      END
C000010
C

      FUNCTION ATANM(Y,X)
C      COMPUTES THE ARCTAN OF Y/X BETWEEN -PI AND +PI.
C
      PIE=3.141592653589793
      PI2=PIE/2.
      ATANM=SIGN(PI2,Y)
      IF(X.NE.0.) ATANM=ATAN(Y/X)
      IF(X.GE.0.) RETURN
      IF(Y.GE.0.) ATANM=PIE+ATANM
      IF(Y.LT.0.) ATANM=-PIE+ATANM
      RETURN
      END
C000011
C
SUB04570
SUB04580
SUB04590
SUB04600
SUB04610
SUB04620
SUB04630
SUB04640
SUB04650
SUB04660
SUB04670
SUB04680
SUB04690
SUB04700
SUB04710
SUB04720
SUB04730
SUB04740
SUB04750
SUB04760

SUB04780
SUB04790
SUB04800
SUB04810
SUB04820
SUB04830
SUB04840
SUB04850
SUB04860

SUB04870
SUB04880
SUB04890
SUB04900
SUB04910
SUB04920
SUB04930
SUB04940
SUB04950
SUB04960
SUB04970
SUB04980
SUB04990
SUB05000
SUB05010

```



```

SUB000100
      READ (NF,89) NO,ZZ
      WRITE(6,90) NO,ZZ
      DO 10 I=1,NO
      READ(NF,88) RO(I)
      WRITE(6,88) RO(I)
      CONTINUE
      FORMAT (I5,F9.3)
10      89      88      90
      FORMAT(F8.5)
      FORMAT(IX,I5,F9.3)
      RETURN
      END
C000013
C

SUB000120
SUB000140
SUB000150
SUB000160
SUB000170

SUB000180
SUB000190
SUB000200
SUB000210
SUB000220
SUB000230
SUB000240
SUB000250
SUB000260
SUB000270
SUB000280
SUB000290
SUB000300
SUB000310
SUB000320
SUB000330
SUB000340
SUB000350
SUB000360
SUB000370
SUB000380
SUB000390
SUB000400
SUB000410
SUB000420
SUB000430
SUB000440
SUB000450
SUB000460
SUB000470
SUB000480
SUB000490
SUB000500

      SUBROUTINE GPRINT (G,NUMB)
      GPRINT PRINTS THE DATA ARRAY 'G' WHICH WAS INPUT TO
      THE PROGRAM IN SUBROUTINE GARRAY.
      COMMON IMAX,JMAX,IIMX,JJMX,IMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
      DIMENSION G(IJMX)
      DIMENSION X(15)
      DATA HYP,VERT/1H-,1H/
      IF (NUMB.EQ.1) WRITE (6,99) MODE,Z
      IF (NUMB.EQ.2) WRITE (6,92) Z
      JM2=JMAX/2
      RIMAX=IMAX
      RJMAX=JMAX
      DX=SIZE/RIMAX
      DXI=360./RJMAX
      INTRVL SETS THE NUMBER OF TERMS PRINTED PER LINE. IF IT IS ALTERED,
      ONE MUST ALSO REDIMENSION X AND ALTER FORMATS 98, 97, AND 95.
      INTRVL=15
      IB=1
      IT=IB+INTRVL-1
      IF (IT.GT.IMAX) IT=IMAX
      IBT=IT-IB+1
      WRITE (6,98) (II,II=IB,IT)
      DO 2 I=1,IBT
      RI=IB-1+I
      X(I)=-SIZE/2.+(RI-.5)*DX
      LM=7*IBT+1
      WRITE (6,97) (X(I),I=1,IBT)
      WRITE (6,96) (HYP,I=1,LM),VERT
      JMH=JM2+1
      DO 3 J=JMH,JMAX
      RJ=J

```


SUB000510
SUB000520
SUB000530
SUB000540
SUB000550
SUB000560
SUB000570
SUB000580
SUB000590
SUB000600
SUB000610
SUB000620
SUB000630
SUB000640
SUB000650
SUB000660
SUB000670
SUB000680
SUB000690
SUB000700
SUB000710
SUB000720
SUB000730

```

3      XI=-180.+DXI*(RJ-.5)
      IGB=(J-1)*IMAX+IB
      IGT=IGB-IB+IT
      WRITE (6,95) J,XI,(G(L),L=IGB,IGT)
      WRITE (6,94) (HYP,L=1,LM),VERT
      IB=IB+INTRVL
      ITOLD=IT
      IT=IT+INTRVL
      IF (ITOLD.LT.IMAX) GO TO 1
      WRITE (6,93)
      FORMAT (1H1//,
99      1 MODE ',11,
      THE ARRAY OF INPUT DATA (G), OBTAINED BY GARRAY,
      CM: ')
      FOR Z='F7.3', '15I7')
98      FORMAT (//11X, 'I= ',15F7.3)
97      FORMAT (//11X, 'X= ',15F7.3)
96      FORMAT (//11X, 'J= ',15F7.3)
95      FORMAT (//11X, 'XI= ',15F7.3)
94      FORMAT (//11X, 'I= ',15F7.3)
93      FORMAT (//11X, 'J= ',15F7.3)
92      FORMAT (//11X, 'I= ',15F7.3)
      THE ADD-ON FUNCTION GARRAY FOR Z='F7.3', CM: ')
      RETURN
      END
C000014
C

```

SUB000740
SUB000750
SUB000760
SUB000770
SUB000780
SUB000790
SUB000800
SUB000810
SUB000820
SUB000830
SUB000840
SUB000850
SUB000860
SUB000870
SUB000880

```

C      SUBROUTINE GPUNCH (Z,XO,YO,PHS,NOF,IMX,JMX,G)
C      GPUNCH PUNCHES OUT THE FIRST NON-SYMMETRIC PORTION OF GARRAY
C      (OR WRITES IT ON FILE 7 IN CMS VERSION)
C
      COMMON /SYM/ ISM,JSM,MSM,FCU,IMS,JMS,QSM
      DIMENSION G(IMX,JMX)
      WRITE (7,39) NOF,IMX,JMX,ISM,JSM,IMS,JMS
      WRITE (7,38) ((G(I,J),I=1,IMS),J=1,JMS)
39      FORMAT(10I5)
38      FORMAT(10F7.3)
      RETURN
      END
C000015
C

```

SUB000890
SUB000900
SUB000910
SUB000920
SUB000930
SUB000940

```

C      SUBROUTINE READ (Z,XO,YO,PHISYM,NOF,IMAX,JMAX,G)
C      READS THE NON-SYMMETRIC PORTION OF THE GARRAY AND EXPANDS IT TO AN
C      ORTHOGONAL SET. NOTE, INSURE SUFFICIENT DIMENSIONS IN MAIN PROGRAM.
C
      COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM

```



```

COMMON /IO/ CMS, IN1, IN2, IN4
DIMENSION G(IMAX, JMAX)
READ (IN1, 39) NOF, IMAX, JMAX, ISYM, JSYM, IMS, JMS
READ (IN1, 38) Z, XO, YO, PHISYM
READ (IN1, 38) ((G(I, J), I=1, IMS), J=1, JMS)
WRITE(6, 37) NOF, Z, XO, YO, PHISYM, IMAX, JMAX, JSYM
RJMX=JMAX
MSYM=JSYM
IF ((MSYM.EQ.0).OR.(MSYM.GT.JMAX)) MSYM=1
FCU=ISYM*JSYM*JMAX
IF (JSYM.GT.JMAX) FCU=JMAX
QSYM=FCU/RJMX
DO 4 J=1, JMS
  IF (ISYM.EQ.1) GO TO 2
  DO 1 I=1, IMS
    II=IMAX+I-1
    G(II, J)=G(I, J)
  GO TO 4
  J2=JMAX/2+1-J
  J3=JMAX/2+J
  J4=JMAX+1-J
  DO 3 I=1, IMAX
    II=IMAX+I-1
    G(I, J2)=G(I, J)
    G(II, J3)=G(I, J)
    G(II, J4)=G(I, J)
  CONTINUE
  FORMAT(10I5)
  FORMAT(10F7.3)
  MODE 3 READS GARRAY DIRECTLY: NOF='I4/'
  FORMAT(//', F7.3,', PHISYM=', F7.3/'
  1,', XO=', F7.3,', JSYM=', I4//')
  2, RETURN
  END
C000016
C

```

```

SUBROUTINE MAP (IM, JM, A, N, Z, BAND)
SUBROUTINE MIMPII AND PLOTS A CONTOUR MAP OF THE ARRAY
DIMENSION A(IM, JM), T(24)
DATA BL/1H /
DO 1 I=1, 24
  T(I)=BL
  ICON=1
  IF(BAND.LT.0.) ICON=0

```


SUB01410
SUB01420
SUB01430
SUB01440
SUB01450
SUB01460
SUB01470
SUB01480
SUB01490
SUB01500
SUB01510
SUB01520

IF(BAND.LT.0.) BAND=-BAND
AMIN=0.
IJT=0
AZ=1.
BZ=0.
WRITE(6,49) N,Z
CALL MTMPII (A, IM, JM, T, BAND, AZ, BZ, AMIN, IJT, ICON)
FORMAT (1H1//, ' THE FUNCTION SURFACE, TEST NO.', I3, ' Z='F5.3//)
49 RETURN
END
C000017
C

SUB01530
SUB01540
SUB01550
SUB01560
SUB01570
SUB01580
SUB01590
SUB01600
SUB01610
SUB01620
SUB01630
SUB01640
SUB01650
SUB01660
SUB01670
SUB01680
SUB01690
SUB01700
SUB01710
SUB01720
SUB01730
SUB01740
SUB01750
SUB01760
SUB01770
SUB01780
SUB01790
SUB01800
SUB01810
SUB01820
SUB01830
SUB01840
SUB01850
SUB01860

SUBROUTINE GPLOT (G, GA, JMS)
C
C GPLOT PRINTS A ROUGH PLOT OF THE LINE INTEGRAL FUNCTIONS IN GARRAY.
C
COMMON IMAX, JMAX, IIMX, JJMX, IJMX, ALPHA, SIZE, EPS, MODE, BOX, SD, IX, Z
COMMON /TAB/ INDEX(7), JSYM, ISYM
COMMON /TAB2/ IPT, KPT, LPT, MPT, REST(5)
DIMENSION G(IMAX, JMAX), GA(IMAX, JMAX), ROW(101)
DIMENSION A(201), B(101), C(201), D(101)
JM=101
DATA BL, PL, ST, DH, EX/1H, 1H+, 1H*, 1H-, 1HX/
JMS2=JMAX/2+1
IF(ISYM.EQ.2) JMS2=1
JMS3=JMS2+JMS-1
DO 8 J=JMS2, JMS3
WRITE(6, 67) (ST, I=1, 120)
DO 1 I=1, IMAX
A(I)=G(I, J)
C(I)=GA(I, J)
AS=.5
BS=.0
CALL INTERP (A, IMAX, AS, B, JM, BS)
CALL INTERP (C, IMAX, AS, D, JM, BS)
WRITE (6, 69) J
BIG=0.
SMALL=0.
DO 2 I=1, IMAX
IF(A(I).GT.BIG) BIG=A(I)
IF(C(I).GT.BIG) BIG=C(I)
IF(A(I).LT.SMALL) SMALL=A(I)
IF(C(I).LT.SMALL) SMALL=C(I)
RANGE=BIG-SMALL
RINK=RANGE/80.
TOP=BIG+RINK
2

SUB01870
SUB01880
SUB01890
SUB01900
SUB01910
SUB01920
SUB01930
SUB01940
SUB01950
SUB01960
SUB01970
SUB01980
SUB01990
SUB02000
SUB02010
SUB02020
SUB02030
SUB02040
SUB02050
SUB02060
SUB02070
SUB02080
SUB02090
SUB02100
SUB02110
SUB02120
SUB02130
SUB02140
SUB02150
SUB02160
SUB02170
SUB02180
SUB02190
SUB02200
SUB02210

```

CEN=BIG
BOT=BIG-RINK
KC=0
DO 7 K=1,41
  IC=0
  DO 6 I=1,101
    ROW(I)=BL
    IF((I.EQ.1).OR.(I.EQ.51).OR.(I.EQ.101)) ROW(I)=PL
    IF((I.EQ.1).OR.(K.EQ.41)) ROW(I)=PL
    IF((K.EQ.1).OR.(BOT.LE.0.)) ROW(I)=DH
    IF((TOP.GE.0.)) AND.(BOT.LE.0.)) ROW(I)=DH
    IF((I.EQ.5) GO TO 3
    GO TO 4
    IC=0
    IF(KC.EQ.10) ROW(I)=PL
    IF(KPT.LE.2) GO TO 5
    IF((D(I).LE.TOP).AND.(D(I).GE.BOT)) ROW(I)=ST
    IF((B(I).LE.TOP).AND.(B(I).GE.BOT)) ROW(I)=EX
    IC=IC+1
    IF(KC.EQ.5) KC=0
    IF(KC.NE.0) WRITE (6,65) (ROW(I),I=1,101)
    IF(KC.EQ.0) WRITE (6,68) CEN,(ROW(I),I=1,101)
    TOP=CEN-2.*RINK
    CEN=CEN-2.*RINK
    BOT=BOT-2.*RINK
    KC=KC+1
    WRITE (6,66) (ST,I=1,120)
    J=I3///
    FORMAT (1X,F8.3,1X,101A1)
    FORMAT (1H1,///121A1//)
    FORMAT (///121A1//)
    FORMAT (10X,101A1)
    RETURN
  END

```

C000018
C

SUB02220
SUB02230
SUB02240
SUB02250
SUB02260
SUB02270
SUB02280
SUB02290
SUB02300
SUB02310
SUB02320

```

SUBROUTINE INTERP (A,IM,AS,B,JM,BS)

INTERP CONVERTS A REGULAR VECTOR A OF IM POINTS TO A REGULAR VECTOR
B OF JM POINTS. OS=.5 FOR A VECTOR WITH POINTS DEFINED IN THE
CENTER OF THE INTERVAL, AS AND BS ARE THE % OF AN INTERVAL FROM THE
EDGE OF THE FIELD TO THE FIRST POINT (.0 OR .5 FOR EDGE OR CENTER
DEFINED POINTS)

```

```

  DIMENSION A(IM),B(JM)
  RIM=IM
  RJM=JM

```

C
C
C
C
C
C

SUB02330
SUB02340
SUB02350
SUB02360
SUB02370
SUB02380
SUB02390
SUB02400
SUB02410
SUB02420
SUB02430
SUB02440
SUB02450
SUB02460
SUB02470

```

RAT=(RIM-1.+2.*AS)/(RJM-1.+2.*BS)
DO 2 I=1,JM
  BI=I
  AI=RAT*(BI+BS)-AS
  IA=AI
  F=AI-FLOAT(IA)
  IF((IA.EQ.0).OR.(IA.EQ.JM)) GO TO 1
  B(I)=A(IA)+F*(A(IA+1)-A(IA))
  GO TO 2
1 IF(IA.EQ.0) B(I)=A(1)*(F-AS)/(1.-AS)
  IF(IA.EQ.JM) B(I)=A(JM)*F/(1.-AS)
  CONTINUE
  RETURN
  END

```

1 2

C000019

SPL00010
SPL00020
SPL00030
SPL00040
SPL00050
SPL00060
SPL00070
SPL00080
SPL00090
SPL00100
SPL00110
SPL00120
SPL00130
SPL00140
SPL00150
SPL00160
SPL00170
SPL00180
SPL00190
SPL00200
SPL00210
SPL00220
SPL00230
SPL00240
SPL00250
SPL00260
SPL00270
SPL00280
SPL00290
SPL00300
SPL00310
SPL00320
SPL00330

SUBROUTINE SPLINE

PURPOSE
PROVIDES INTERPOLATED VALUE USING "CUBIC SPLINE FITTING"

USAGE
FIRST CALL TO SUBROUTINE:
CALL SPLINE(X,Y,M,XINT,YINT)
SUBSEQUENT CALLS:
CALL SPLINN(X,Y,M,XINT,YINT)

DESCRIPTION OF PARAMETERS
X: MONOTONICALLY INCREASING ABSCISSA ARRAY
Y: ONE-FOR-ONE CORRESPONDING ORDINATE ARRAY
M: NUMBER OF X AND Y VALUES SUPPLIED < OR = 300
XINT: VALUE OF ABSCISSA FOR WHICH CORRESPONDING ORDINATE
IS TO BE INTERPOLATED (OR EXTRAPOLATED)
YINT: INTERPOLATED (OR EXTRAPOLATED) ORDINATE VALUE

REMARKS
IF SPECIFIED X FALLS OUTSIDE OF RANGE, AN EXTRAPOLATED
VALUE WILL BE SUPPLIED

SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED
SUBROUTINE SPLICO IS INCLUDED IN SUBROUTINE SPLIN PACKAGE
MATHEMATICAL METHOD
UPON FIRST ENTRY TO SPLIN, A CALL TO SPLICO IS MADE TO

CC

SPL00340
SPL00350
SPL00360
SPL00370
SPL00380
SPL00390
SPL00400
SPL00410
SPL00420
SPL00430

DETERMINE THE COEFFICIENTS TO BE USED IN PERFORMING THE
INTERPOLATIONS. SEARCH FOR BRACKETING ABSCISSA VALUES IS
ALWAYS MADE FROM THE REFERENCE LAST USED IN INTERPOLATING.

REFERENCE
PENNINGTON, RALPH H., "INTRODUCTORY COMPUTER METHODS AND
NUMERICAL ANALYSIS", THE MACMILLAN COMPANY, NEW YORK, 1965

SPL00440
SPL00460
SPL00470
SPL00480
SPL00490
SPL00500
SPL00510
SPL00520
SPL00530
SPL00540
SPL00550
SPL00560
SPL00570
SPL00580
SPL00590
SPL00600
SPL00610
SPL00620
SPL00630
SPL00640
SPL00650
SPL00660
SPL00670
SPL00680
SPL00690
SPL00700
SPL00710
SPL00720

```

SUBROUTINE SPLINE(X,Y,M,XINT,YINT)
DIMENSION X(M),Y(M),C(4,300)
CALL SPLICO(X,Y,M,C)
K=1
ENTRY SPLIN(X,Y,M,XINT,YINT)
IF(XINT-X(1)) 70,1,2
3 70 K=1
GO TO 7
1 YINT=Y(1)
RETURN
2 IF(XINT-X(K+1)) 6,4,5
4 YINT=Y(K+1)
RETURN
5 K=K+1
IF(M-K) 71,71,3
71 K=M-1
GO TO 7
6 IF(XINT-X(K)) 13,12,11
12 YINT=Y(K)
RETURN
13 K=K-1
GO TO 6
7 PRINT 101,XINT
101 FORMAT(8H,XINT = E18.9,32H, OUT OF RANGE FOR INTERPOLATION)
11 YINT=(X(K+1)-XINT)*(C(1,K)-XINT)**2+C(3,K))
YINT=YINT+(XINT-X(K))*(C(2,K)-XINT-X(K))**2+C(4,K))
RETURN
END

```

SPL00730
SPL00750
SPL00760
SPL00770
SPL00780
SPL00790

```

SUBROUTINE SPLICO(X,Y,M,C)
DIMENSION X(M),Y(M),C(4,300),D(300),P(300),E(300),A(300,3),B(300),
1Z(300)
MM=M-1
DO 2 K=1,MM
D(K)=X(K+1)-X(K)

```



```

      P(K)=D(K)/6.-Y(K+1)-Y(K))/D(K)
2  DO 3 K=2,MN
3  B(K)=E(K)-E(K-1)
   A(1,2)=-1.-D(1)/D(2)
   A(1,3)=D(1)/D(2)
   A(2,3)=P(2)-P(1)*A(1,3)
   A(2,2)=2.*(P(1)+P(2))-P(1)*A(1,2)
   B(2,3)=A(2,3)/A(2,2)
   B(2,2)=B(2,3)/A(2,2)
   DO 4 K=3,MN
   A(K,2)=2.*(P(K-1)+P(K))-P(K-1)*A(K-1,3)
   B(K,2)=B(K)-P(K-1)*B(K-1)
   A(K,3)=P(K)/A(K,2)
4  B(K)=B(K)/A(K,2)
   Q=D(M-2)/D(M-1)
   A(M,1)=1.+Q+A(M-2,3)
   A(M,2)=-Q-A(M,1)*A(M-1,3)
   B(M)=B(M-2)-A(M,1)*B(M-1)
   Z(M)=B(M)/A(M,2)
   MN=M-2
   DO 6 I=1,MN
   K=M-I
6  Z(K)=B(K)-A(K,3)*Z(K+1)
   Z(1)=-A(1,2)*Z(2)-A(1,3)*Z(3)
   DO 7 K=1,MN
   Q=1./(6.*D(K))
   C(1,K)=Z(K)*Q
   C(2,K)=Z(K+1)*Q
   C(3,K)=Y(K)/D(K)-Z(K)*P(K)
7  C(4,K)=Y(K+1)/D(K)-Z(K+1)*P(K)
   RETURN
   END

```

C000020

MET000310
 MET000320
 MET000330
 MET000340
 MET000350
 MET000360
 MET000370
 MET000380
 MET000390
 MET000400
 MET000410
 MET000420

SUBROUTINE MTMPII

PURPOSE

MTMPII WILL PRODUCE, ON THE PRINTER, A CONTOUR MAP
OF ANY SINGLE PRECISION TWO DIMENSIONAL ARRAY.

USAGE

CALL MTMPII(Y,N,M,T,BND,AZ,BZ,AMIN,IJT,ICON)

CC

DESCRIPTION OF PARAMETERS		
Y	THE ARRAY TO BE CONTOURED. DIMENSIONED Y(N,M)	MET00130
N	NUMBER OF ROWS IN Y.	MET00140
M	NUMBER OF COLUMNS IN Y.	MET00150
T	ARRAY FOR PLOT TITLE.	MET00160
BND	BANDWIDTH FOR THE CONTOURING. IF BND IS ZERO A BANDWIDTH WILL BE CALCULATED AS FOLLOWS BND=(MAX(Y)-MIN(Y))/15.	MET00170 MET00180 MET00190 MET00200 MET00210
AZ	A LINEAR TRANSFORMATION MAYBE PERFORMED ON THE ARRAY Y OF THE FOLLOWING FORM AZ=Y+BZ. IF AZ=0 THEN AZ WILL BE COMPUTED SUCH THAT MAX(I MAX(Y)), I MIN(Y)) WILL BE LESS THAN 1, AND BZ WILL BE LEFT AS INPUT.	MET00220 MET00230 MET00240 MET00250 MET00260
BZ	SEE UNDER AZ	MET00270
AMIN	THE LEVEL AT WHICH COUTOURING WILL BEGIN. IF AMIN > MIN(Y) THEN AMIN WILL BE CALCULATED TO BE WITH REFERENCE AT ZERO, THE NEXT LOWER CONTOUR LEVEL FROM MIN(Y) AS DETERMINED BY BND.	MET00280 MET00290 MET00300 MET00310 MET00320
IJT	IF IJT=0 AMIN WILL BE CALCULATED AS DESCRIBED ABOVE.	MET00330
ICON	IF ICON=0 NO CONTOURING WILL BE DONE BUT THE ARRAY Y WILL BE PRINTED IN THE PLOT FORMAT.	MET00340 MET00350 MET00360 MET00370 MET00380
REMARKS		MET00390
MTMPII REQUIRES A PRINTER WITH 132 PRINT POSITIONS. IF NECESSARY THE MAP WILL BE SEGMENTED COLUMNWISE. THE ROWS AND COLUMNS ARE NUMBERED ALONG THE EDGES. SO THAT A SEGMENTED MAP MAYBE EASILY JOINED TOGETHER. ONLY THREE SIGNIFICANT FIGURES WILL BE PRINTED AT EACH POINT. THE POSITION OF THE FIRST SIGNIFICANT DIGIT WILL BE DETERMINED BY MAX(I MAX(Y)), I MIN(Y)). THE PLOT WILL BE PRODUCED ON A INCH GRID. IT WILL BE ASSUMED THAT THE SPACING BETWEEN POINTS IN BOTH DIRECTIONS IS THE SAME AND EQUAL FOR ALL POINTS		MET00400 MET00410 MET00420 MET00430 MET00440 MET00450 MET00460 MET00470 MET00480
SUBROUTINES REQUIRED		MET00490
NONE		MET00500 MET00510 MET00520 MET00530 MET00540 MET00550 MET00560 MET00570 MET00580 MET00590 MET00600
METHOD		
THE CONTOUR LEVELS ARE DETERMINED BY SIMPLE LINEAR INTERPOLATION FROM THE FOUR SURROUNDING PIONTS.		
MTMPII SUBROUTINE FOR ONE-INCH GRID SPACING		

C OAKES CODE 5105 15 JAN 69

MET00610
MET00620

SUBROUTINE MTMPJJ(Y,N,M,T,BND,AZ,BZ,AMIN,IJT,ICON)
REAL*4 IH,KG,IJTJZ
DIMENSION A(140),B(140),C(140),D(140),IH(20),Y(N,M),TP(10),TPX(10)
Z ,TPM(10),XMT(10),BTM(10),BTX(10),BT(10),KG(10),T(24)
DIMENSION E(140),F(140),G(140),H(140)

MET00630
MET00640
MET00650
MET00660
MET00670
MET00680
MET00690
MET00700
MET00710
MET00720
MET00730
MET00740
MET00750
MET00760
MET00770
MET00780
MET00790
MET00800
MET00810
MET00820
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MET00870
MET00880
MET00890
MET00900
MET00910
MET00920
MET00930
MET00940
MET00950
MET00960
MET00970
MET00980
MET00990
MET01000
MET01010
MET01020
MET01030
MET01040
MET01050
MET01060

DATA DUE/4H /,EPL/4H+ /,EMI/4H- /,IH/1H0,1H ,1H1,1H ,1H2,
11H ,1H3,1H ,1H4,1H ,1H5,1H ,1H6,1H ,1H7,1H ,1H8,1H /,KG/
21H0,1H1,1H2,1H3,1H4,1H5,1H6,1H7,1H8,1H9/,BLK/4H

MET00700
MET00710
MET00720
MET00730
MET00740
MET00750
MET00760
MET00770
MET00780
MET00790
MET00800
MET00810
MET00820
MET00830
MET00840
MET00850
MET00860
MET00870
MET00880
MET00890
MET00900
MET00910
MET00920
MET00930
MET00940
MET00950
MET00960
MET00970
MET00980
MET00990
MET01000
MET01010
MET01020
MET01030
MET01040
MET01050
MET01060

YMIN=Y(1,1)
YMAX=Y(1,1)
DO 20 I=1,M
DO 10 J=1,N
YMIN=AMINI(YMIN,Y(J,I))
YMAX=AMAXI(YMAX,Y(J,I))
CONTINUE

MET00700
MET00710
MET00720
MET00730
MET00740
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MET00760
MET00770
MET00780
MET00790
MET00800
MET00810
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MET00830
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MET00890
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MET00920
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MET00940
MET00950
MET00960
MET00970
MET00980
MET00990
MET01000
MET01010
MET01020
MET01030
MET01040
MET01050
MET01060

10 DELY=YMAX-YMIN
20 IF(BND) 25,25,30
25 BND=DELY/15.0
30 IF (AMIN-YMIN) 31,31,32
31 IF (IJT) 33,32,33
32 PD=YMIN/BND
PF=ABS(PD-INT(PD))
IF (YMIN) 2,1,1
AMIN=YMIN-PF*BND
GO TO 33
AMIN=YMIN-(1.0-PF)*BND
AHLD=AZ
IF(AZ) 55,35,55
SM=AMAX1(ABS(YMIN),ABS(YMAX))
NS=0
NS=NS+1
SM=10.0*SM
IF(SM-1.0)40,50,45
NS=NS-1
SM=SM/10.0
IF(SM-1.0)50,50,45
AHLD=10.0*NS
HBND=BND/2.0
PRINT 70
PRINT 6,T

MET00700
MET00710
MET00720
MET00730
MET00740
MET00750
MET00760
MET00770
MET00780
MET00790
MET00800
MET00810
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MET00970
MET00980
MET00990
MET01000
MET01010
MET01020
MET01030
MET01040
MET01050
MET01060


```

6  FORMAT(5X,24A4,/)
PRINT 57,AHLD,BZ
57  FORMAT(1H0,65H THE FOLLOWING TRANSFORMATION WAS PERFORMED ON THE IN
      1PUT MATRIX /5X,1H(,E12.5,8H*Y(I,J)+,E12.5,1H) //2X,73HAND THREE
      2 DIGITS TO THE RIGHT OF THE DECIMAL POINT ARE PRINTED IN THE MAP
      1LS
C
PRINT 54,YMAX,YMIN
54  FORMAT(/4X,5HYMAX=,E15.7,5X,5HYMIN=,E15.7)
5  PRINT 11,BND
11  FORMAT(2X,17H THE BAND WIDTH IS, E12.5,6H UNITS //4X,14H CONTOUR LEVE
      1LS
      I=0
      YTOP=AMIN
      IF(ABS(YMIN-YMAX)-100.0*BND) 53,53,58
53  YB=YTOP
      YTOP=YTOP+BND
      I=I+1
      J=MOD(I,20)
      ITJZ=IH(J)
      IF(YB-YMAX) 59,58,58
59  PRINT 61,YB,YTOP,ITJZ
61  FORMAT(/4X,E10.3,4H TO ,E10.3,2H =,1X,A1)
      GO TO 53
58  NCCP=0
      NCP=0
      PRINT 70
70  FORMAT(1H1)
      PRINT 6,T
      NLINE=0
      NCCP=NCP+1
      NCP = NCP + 13
73  IF(NCP-M) 80,80,75
75  NCP=M
80  CONTINUE
      J=-2
      NLINE=NLINE+1
      LLINE=N-NLINE+1
      UP HEADING
      IF(NCCP-1) 85,85,90
85  J = -1
90  DO 100 I = 1,135
      A(I)=BLK
      B(I)=BLK
      H(I)=BLK
100 CONTINUE
110 DO 160 L=NCCP,NCP
      J = J+8

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MET0101070
MET0101080
MET0101090
MET0101100
MET0101110
MET0101120
MET0101130
MET0101140
MET0101150
MET0101160
MET0101170
MET0101180
MET0101190
MET0101200
MET0101210
MET0101220
MET0101230
MET0101240
MET0101250
MET0101260
MET0101270
MET0101280
MET0101290
MET0101300
MET0101310
MET0101320
MET0101330
MET0101340
MET0101350
MET0101360
MET0101370
MET0101380
MET0101390
MET0101400
MET0101410
MET0101420
MET0101430
MET0101440
MET0101450
MET0101460
MET0101470
MET0101480
MET0101490
MET0101500
MET0101510
MET0101520
MET0101530
MET0101540

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MET01550
 MET01560
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 MET01580
 MET01590
 MET01600
 MET01610
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 MET01690
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 MET01940
 MET01950
 MET01960
 MET01970
 MET01980
 MET01990
 MET02000
 MET02010
 MET02020

```

    KI=L
    IF(KI-100) 130,120,120
    LL=KI/100
    A(J)=KG(LL+1)
    KI=KI-100*LL
    GO TO 135
    A(J)=KG(I)
    J=J+1
    IF(KI-10) 150,140,140
    LL=KI/10
    A(J)=KG(LL+1)
    KI=KI-10*LL
    GO TO 155
    A(J)=KG(I)
    J=J+1
    A(J)=KG(KI+1)
    160 CONTINUE
    C SETUP FIRST ROW OF ARRAY
    GO TO 260
    170 NLINE=NLINE+1
    LLINE=N-NLINE+1
    IF(NLINE-N) 180,180,380
    180 DO 190 I=1,135
      A(I)=BLK
      B(I)=BLK
      C(I)=BLK
      D(I)=BLK
      E(I)=BLK
      F(I)=BLK
      G(I)=BLK
      H(I)=BLK
    190 CONTINUE
    IF (ICON)195,260,195
    195 NCY=NCCP-1
    J=4
    IF(NCY)200,200,210
    200 NCY=NCY+1
    210 J=5
    IF(NCY-NCP) 220,220,260
    220 IF(NCY-M) 230,260,260
    230 NLINE = NLINE - 1
    YD1 = Y(NLINE,NCY) - Y(NLINE+1,NCY)
    YD2=Y(NLINE,NCY+1)-Y(NLINE+1,NCY+1)
    TP(1) = Y(NLINE,NCY)-0.125*YD1
    TPX(1)=Y(NLINE,NCY)-0.250*YD1
    TPM(1)=Y(NLINE,NCY)-0.375*YD1
    XMT(1)=Y(NLINE,NCY)-0.500*YD1
    BTM(1)=Y(NLINE,NCY)-0.625*YD1
  
```



```

BTX(1)=Y(NLINE,NCY)-0.750*YD1
BT(1)=Y(NLINE,NCY)-0.875*YD1
TP(10)=Y(NLINE,NCY+1)-0.125*YD2
TPM(10)=Y(NLINE,NCY+1)-0.250*YD2
XMT(10)=Y(NLINE,NCY+1)-0.375*YD2
BTM(10)=Y(NLINE,NCY+1)-0.500*YD2
BTX(10)=Y(NLINE,NCY+1)-0.625*YD2
BT(10)=Y(NLINE,NCY+1)-0.750*YD2
NLINE = NLINE + 1
D1=0.1*(TP(10)-TP(1))
D2=0.1*(TPX(10)-TPX(1))
D3=0.1*(TPM(10)-TPM(1))
D4=0.1*(XMT(10)-XMT(1))
D5=0.1*(BTM(10)-BTM(1))
D6=0.1*(BTX(10)-BTX(1))
D7=0.1*(BT(10)-BT(1))
DO 240 I = 2,9
TP(I)=TP(I-1)+D1
TPX(I)=TPX(I-1)+D2
TPM(I)=TPM(I-1)+D3
XMT(I)=XMT(I-1)+D4
BTM(I)=BTM(I-1)+D5
BTX(I)=BTX(I-1)+D6
BT(I)=BT(I-1)+D7
CONTINUE
DO 250 I = 1,10
J=J+1
I1=MOD(IFIX((TP(I)-AMIN)/BND),20)+1
I2=MOD(IFIX((TPX(I)-AMIN)/BND),20)+1
I3=MOD(IFIX((TPM(I)-AMIN)/BND),20)+1
I4=MOD(IFIX((XMT(I)-AMIN)/BND),20)+1
I5=MOD(IFIX((BTM(I)-AMIN)/BND),20)+1
I6=MOD(IFIX((BTX(I)-AMIN)/BND),20)+1
I7=MOD(IFIX((BT(I)-AMIN)/BND),20)+1
A(J)=IH(I1)
B(J)=IH(I2)
C(J)=IH(I3)
D(J)=IH(I4)
E(J)=IH(I5)
F(J)=IH(I6)
G(J)=IH(I7)
CONTINUE
GO TO 210
250 NCCP=NCCP-1
260 J=-2
IF(NCY) 265,265,270
265 J=-1

```

```

MET02030
MET02040
MET02050
MET02060
MET02070
MET02080
MET02090
MET02100
MET02110
MET02120
MET02130
MET02140
MET02150
MET02160
MET02170
MET02180
MET02190
MET02200
MET02210
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MET02240
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MET02490
MET02500

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MET02510
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 MET02890
 MET02900
 MET02910
 MET02920
 MET02930
 MET02940
 MET02950
 MET02960
 MET02970
 MET02980

```

    GO TO 330
    NCY=NCY+1
    IF(NCY-NCP) 280,280,310
    J=J+7
    THLD=AHLD*Y(NLINE,NCY)+BZ
    IF(THLD) 285,290,290
    H(J)=EMI
    GO TO 295
    H(J)=EPL
    NUM=INT((ABS(THLD-INT(THLD)))*1000.0+0.5)
    NDS=100
    DO 300 KK=1,3
    J=J+1
    KI=NUM/NDS
    H(J)=KG(KI+1)
    NUM=NUM-KI*NDS
    NDS=NDS/10
    CONTINUE
    GO TO 270
    IF(NCP-M) 360,320,320
    IF(J-127)330,330,360
    J=J+3
    KI=NLINE
    IF(KI-100) 340,335,335
    LL=KI/100
    H(J)=KG(LL+1)
    KI=KI-100*LL
    GO TO 343
    H(J)=KG(1)
    J=J+1
    IF(KI-10) 350,345,345
    LL=KI/10
    H(J)=KG(LL+1)
    KI=KI-10*LL
    GO TO 355
    H(J)=KG(1)
    J=J+1
    H(J)=KG(KI+1)
    J=J-5
    IF(NCY-1) 270,270,360
    IF(NLINE-1)362,362,368
    PRINT 370,(A(I),I=1,132),(B(IP1),IP1=1,132),(H(IP2),IP2=1,132)
    GO TO 170
    PRINT 370,(A(I),I=1,132),(B(IP1),IP1=1,132),(C(IP2),IP2=1,132),
    1(D(IP3),IP3=1,132),(E(IP4),IP4=1,132),(F(IP5),IP5=1,132),
    2(G(IP6),IP6=1,132),(H(IP7),IP7=1,132)
    FORMAT(132A1)
    GO TO 170
  
```



```

380 DO 390 I=1,135
    A(I)=BLK
    B(I)=BLK
    C(I)=BLK
    D(I)=BLK
    CONTINUE
390 J=-2
    IF(NCCP-1) 395,395,400
    J=-1
    DO 430 L=NCCP,NCP
        J=J+8
        KI=L
        IF(KI-100) 410,405,405
        LL=KI/100
        C(J)=KG(LL+1)
        KI=KI-100*LL
        GO TO 412
        C(J)=KG(1)
        J=J+1
        IF(KI-10) 420,415,415
        LL=KI/10
        C(J)=KG(LL+1)
        KI=KI-10*LL
        GO TO 422
        C(J)=KG(1)
        J=J+1
        C(J)=KG(KI+1)
        CONTINUE
430 PRINT 370, (B(IP1),IP1=1,132),(C(IP2),IP2=1,132)
    IF(NCP-M)60,500,500
500 RETURN
    END

```

```

MET02990
MET03000
MET03010
MET03020
MET03030
MET03040
MET03050
MET03060
MET03070
MET03080
MET03090
MET03100
MET03110
MET03120
MET03130
MET03140
MET03150
MET03160
MET03170
MET03180
MET03200
MET03210
MET03220
MET03230
MET03240
MET03250
MET03260
MET03270
MET03280
MET03290
MET03300

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C000021

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C .....
C SUBROUTINE GAUSS
C
C PURPOSE
C   COMPUTES A NORMALLY DISTRIBUTED RANDOM NUMBER WITH A GIVEN
C   MEAN AND STANDARD DEVIATION
C
C USAGE
C   CALL GAUSS(IX,S,AM,V)
C
C DESCRIPTION OF PARAMETERS
C
C GAUS 10
C GAUS 20
C GAUS 30
C GAUS 40
C GAUS 50
C GAUS 60
C GAUS 70
C GAUS 80
C GAUS 90
C GAUS 100
C GAUS 110
C GAUS 120
C GAUS 130

```


AND PRODUCES A NEW INTEGER AND REAL RANDOM NUMBER.			
USAGE	CALL RANDU(IX,IY,YFL)		
DESCRIPTION OF PARAMETERS			
IX -	FOR THE FIRST ENTRY THIS MUST CONTAIN ANY ODD INTEGER NUMBER WITH NINE OR LESS DIGITS. AFTER THE FIRST ENTRY, IX SHOULD BE THE PREVIOUS VALUE OF IY COMPUTED BY THIS SUBROUTINE.		
IY -	A RESULTANT INTEGER RANDOM NUMBER REQUIRED FOR THE NEXT ENTRY TO THIS SUBROUTINE. THE RANGE OF THIS NUMBER IS BETWEEN ZERO AND 2**31		
YFL -	THE RESULTANT UNIFORMLY DISTRIBUTED, FLOATING POINT, RANDOM NUMBER IN THE RANGE 0 TO 1.0		
REMARKS	<p>THIS SUBROUTINE IS SPECIFIC TO SYSTEM/360 AND WILL PRODUCE 2**29 TERMS BEFORE REPEATING. THE REFERENCE BELOW DISCUSSES SEEDS (65539 HERE), RUN PROBLEMS, AND PROBLEMS CONCERNING RANDOM DIGITS USING THIS GENERATION SCHEME. MACCLAREN AND MARSAGLIA, JACM 12, P. 83-89, DISCUSS CONGRUENTIAL GENERATION METHODS, AND TESTS. THE USE OF TWO GENERATORS OF THE RANDU TYPE, ONE FILLING A TABLE AND ONE PICKING FROM THE TABLE, IS OF BENEFIT IN SOME CASES. 65549 HAS BEEN SUGGESTED AS A SEED WHICH HAS BETTER STATISTICAL PROPERTIES FOR HIGH ORDER BITS OF THE GENERATED DEVIATE.</p> <p>SEEDS SHOULD BE CHOSEN IN ACCORDANCE WITH THE DISCUSSION THAT GIVEN IN THE REFERENCE BELOW. ALSO, IT SHOULD BE NOTED THAT IF FLOATING POINT RANDOM NUMBERS ARE DESIRED, AS ARE AVAILABLE FROM RANDU, THE RANDOM CHARACTERISTICS OF THE FLOATING POINT DEVIATES ARE MODIFIED AND IN FACT THESE DEVIATES HAVE HIGH PROBABILITY OF HAVING A TRAILING LOW ORDER ZERO BIT IN THEIR FRACTIONAL PART.</p>		
SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED	NONE		
METHOD	POWER RESIDUE METHOD DISCUSSED IN IBM MANUAL C20-8011, RANDOM NUMBER GENERATION AND TESTING		

CC


```

SUBROUTINE RANDU(IX,IY,YFL)
  IY=IX#65539
  IF(IY)5,6,6
  5 IY=IY+2147483647+1
  6 YFL=IY
  YFL=YFL*.4656613E-9
  RETURN
END

```

```

RAND 540
RAND 550
RAND 560
RAND 570
RAND 580
RAND 590
RAND 600

```


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14

KEY WORDS

LINK A

LINK B

LINK C

ROLE

WT

ROLE

WT

ROLE

M/T

Wing-Fuselage Junction

1 MAR 74

20978

Thesis

141295

K823 Kosakoski

c.1

Application of holo-
graphic interferometry to
density field determi-
nation in transonic
corner flow.

1 MAR 74

20978

Thesis

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c.1

Application of holo-
graphic interferometry to
density field determi-
nation in transonic
corner flow.

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